



UC SANTA CRUZ

# Stefano Profumo

Santa Cruz Institute for Particle Physics  
University of California, Santa Cruz

## Dark Energy and Dark Matter Lecture 2

XXVII Special Courses at the Observatório Nacional – Rio de Janeiro, Brazil

## Key ideas from last lecture

- **General Relativity** and the **metric** of a homogeneous and isotropic **universe**
- **Thermodynamics** of the early universe – the **thermal decoupling paradigm** ( $\nu, \gamma$ )
- **Cold** Relic Freeze out

# Cold Relics

Try **proton-antiproton** freeze-out:  
what's the **relic** matter **abundance** in a baryon-symmetric Universe?

$$\sigma \sim \Lambda_{\text{QCD}}^{-2}$$

$$n \sigma = H \rightarrow T^3 \Lambda^{-2} = T^2/M_p \rightarrow T = \Lambda^2/M_p$$

doesn't quite work, we're way **outside**  
the regime of validity for **hot relics**, ( $T \gg m_p$ ), since we found  $T \ll \ll \ll \ll m_p \dots$

Need to work out the case of **cold relics**, which looks nastier by eye

$$n \sim (m_\chi T)^{3/2} \exp\left(-\frac{m_\chi}{T}\right)$$

# Cold Relics

$$n \sim (m_\chi T)^{3/2} \exp\left(-\frac{m_\chi}{T}\right)$$

$$n \sigma \sim H \sim T^2/M_P$$

[freeze out condition]

$$n_{\text{f.o.}} \sim \frac{T_{\text{f.o.}}^2}{M_P \cdot \sigma}$$

now define  $m_\chi/T \equiv x$  (cold relic:  $x \gg 1$ )

**Freeze-out** condition (x) now reads

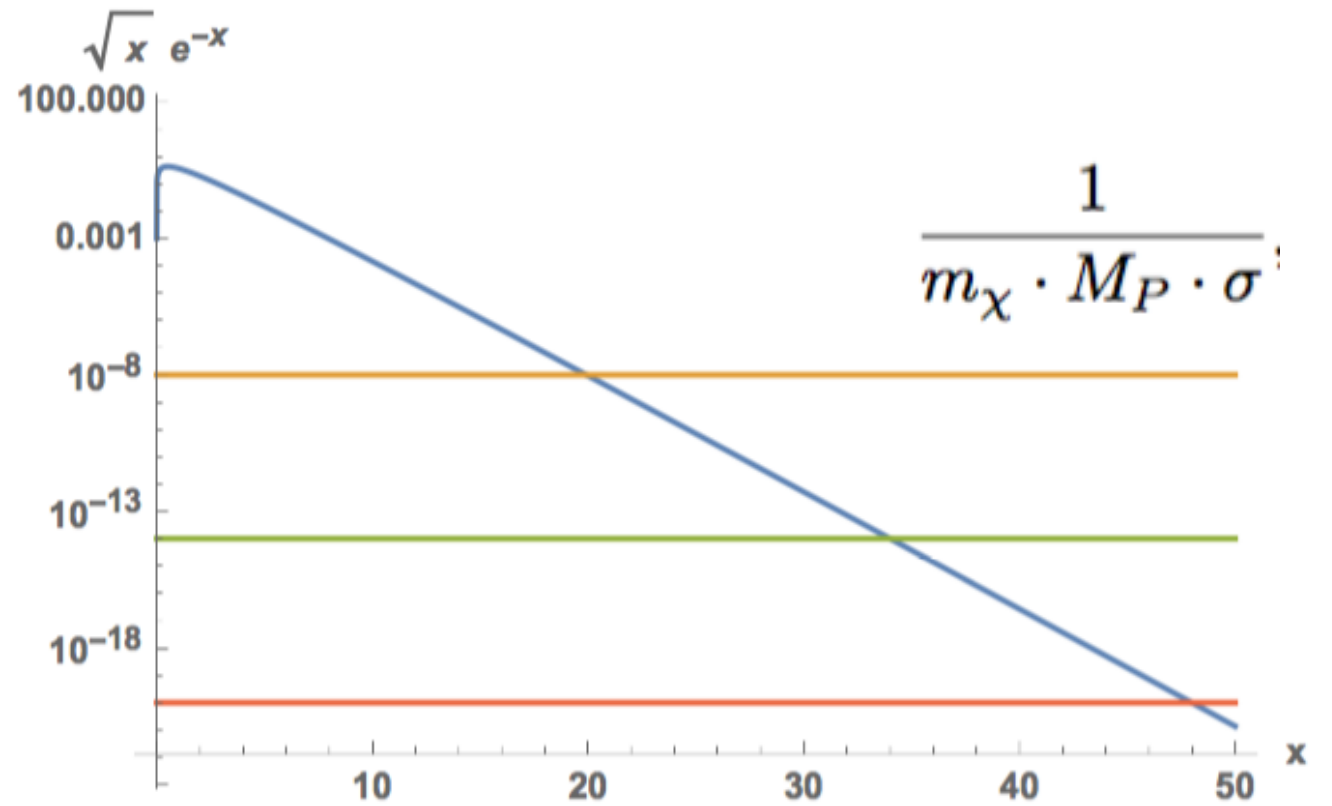
$$\frac{m_\chi^3}{x^{3/2}} e^{-x} = \frac{m_\chi^2}{x^2 \cdot M_P \cdot \sigma}$$

...so we must **solve**

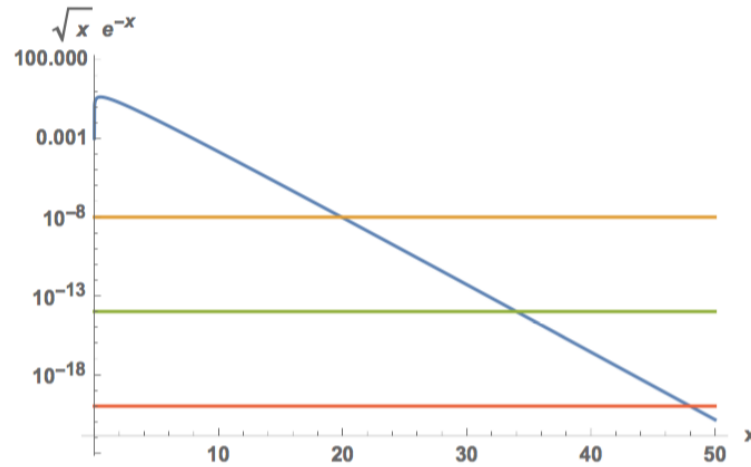
$$\sqrt{x} \cdot e^{-x} = \frac{1}{m_\chi \cdot M_P \cdot \sigma}$$

# Cold Relics

$$\sqrt{x} \cdot e^{-x} = \frac{1}{m_\chi \cdot M_P \cdot \sigma}$$



# Cold Relics



$$\sqrt{x} \cdot e^{-x} = \frac{1}{m_\chi \cdot M_P \cdot \sigma}$$

Take e.g. a "**weakly interacting massive particle**"

$$\sigma \sim G_F^2 m_\chi^2$$

$$m_\chi \sim 10^2 \text{ GeV.}$$

$$\sqrt{x} \cdot e^{-x} = \frac{1}{m_\chi \cdot M_P \cdot \sigma} \sim \frac{1}{10^2 \cdot 10^{18} \cdot 10^{-6}} \sim 10^{-14}$$

$$x = m_\chi / T \sim 35$$

$$\text{proton-antiproton: } 1/(1 \text{ GeV } 10^{18} \text{ GeV } 10^{-2} \text{ GeV}^{-2}) \sim 10^{-16}$$

$$x = m_p / T \sim 37$$

# Cold Relics

Off to calculating the **thermal relic density**  $\Omega_\chi = \frac{m_\chi \cdot n_\chi(T = T_0)}{\rho_c} = \frac{m_\chi T_0^3}{\rho_c} \frac{n_0}{T_0^3}$

iso-entropic universe  $aT \sim \text{const}$   $\frac{n_0}{T_0^3} \simeq \frac{n_{\text{f.o.}}}{T_{\text{f.o.}}^3}$

$$\Omega_\chi = \frac{m_\chi T_0^3}{\rho_c} \frac{n_{\text{f.o.}}}{T_{\text{f.o.}}^3} = \frac{T_0^3}{\rho_c} x_{\text{f.o.}} \left( \frac{n_{\text{f.o.}}}{T_{\text{f.o.}}^2} \right) = \left( \frac{T_0^3}{\rho_c M_P} \right) \frac{x_{\text{f.o.}}}{\sigma}$$

$$\left( \frac{\Omega_\chi}{0.2} \right) \simeq \frac{x_{\text{f.o.}}}{20} \left( \frac{10^{-8} \text{ GeV}^{-2}}{\sigma} \right)$$

# Cold Relics

$$\left(\frac{\Omega_\chi}{0.2}\right) \simeq \frac{x_{\text{f.o.}}}{20} \left(\frac{10^{-8} \text{ GeV}^{-2}}{\sigma}\right)$$

If you have a **WIMP**,  
defined by a cross section

$$\sigma \sim G_F^2 m_\chi^2$$

$$\Omega_\chi h^2 \sim 0.1 \frac{10^{-8} \text{ GeV}^{-2}}{G_F^2 m_\chi^2} \sim 0.1 \left(\frac{10 \text{ GeV}}{m_\chi}\right)^2$$

For **protons**:

$$\sigma \sim 1/(0.1 \text{ GeV})^2$$

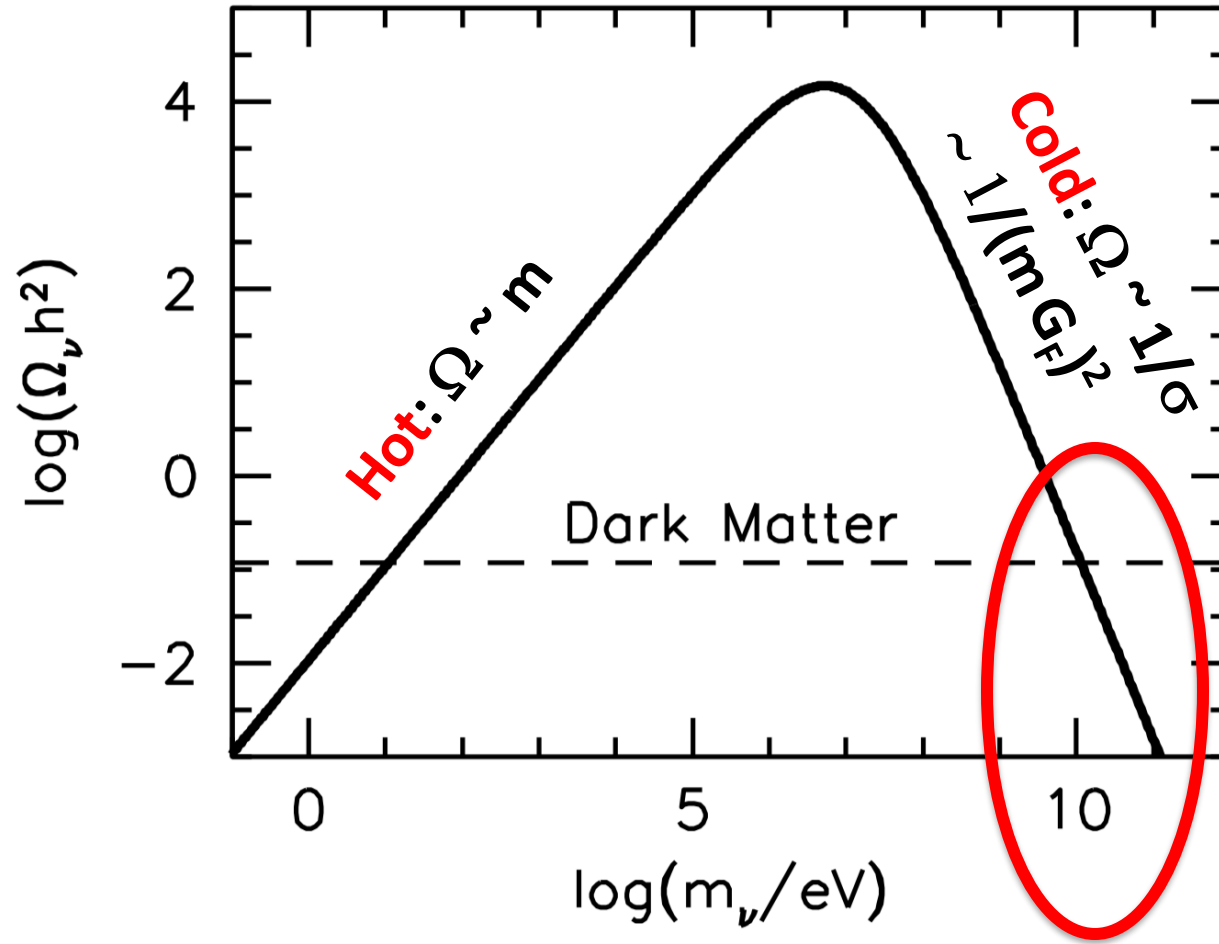
$$\Omega_B \sim 10^{-10}$$

$$\Omega_{B,\text{obs}} \sim 10^{-1}$$

➔ Need **BARYOGENESIS!**



WIMP's **thermal** relic density



"**Lee-Weinberg**" limit

"Classical" **WIMPs**  
**cannot be lighter**  
than  $\sim 1$  GeV

# The CDM in Lambda CDM

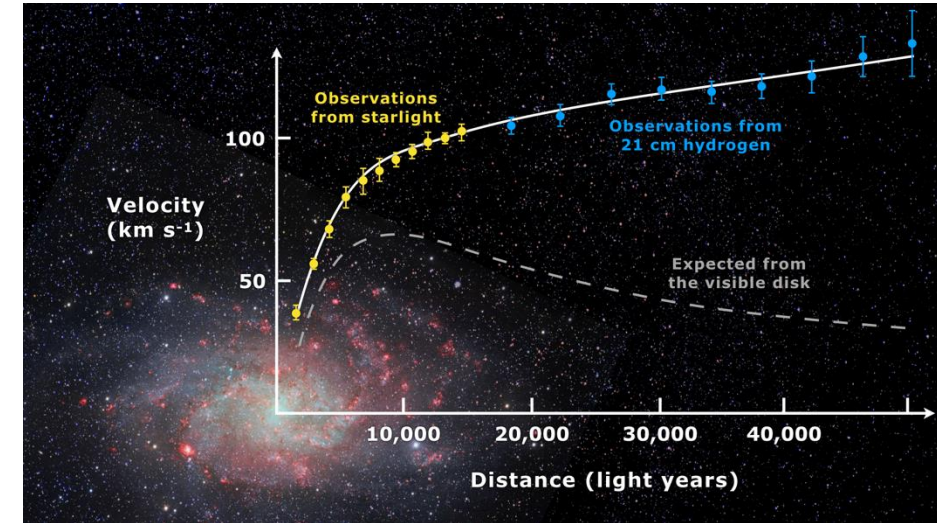


# Astrophysical Evidence

## 1. Galaxy Rotation Curves

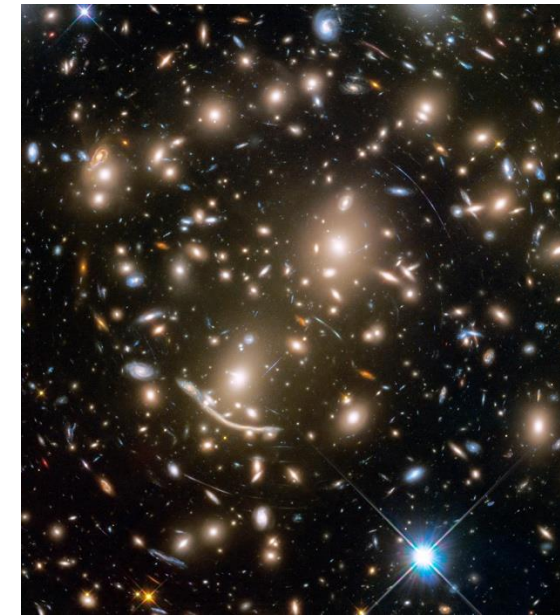
**Observation:** The rotational velocity of stars in spiral galaxies remains approximately constant at large radii, rather than decreasing as expected from the visible mass.

**Implication:** Indicates the presence of an extended, unseen mass component—interpreted as a dark matter halo.



## 2. Galaxy Clusters (Zwicky & the Virial Theorem)

**Observation:** The velocity dispersion of galaxies within clusters suggests much more mass than observed in luminous matter.





# Astrophysical Evidence

## 3. Gravitational Lensing

**Observation:** The deflection of light from background objects by foreground mass often exceeds what is expected from visible matter alone.

**Examples:**

Strong lensing (Einstein rings, arcs)

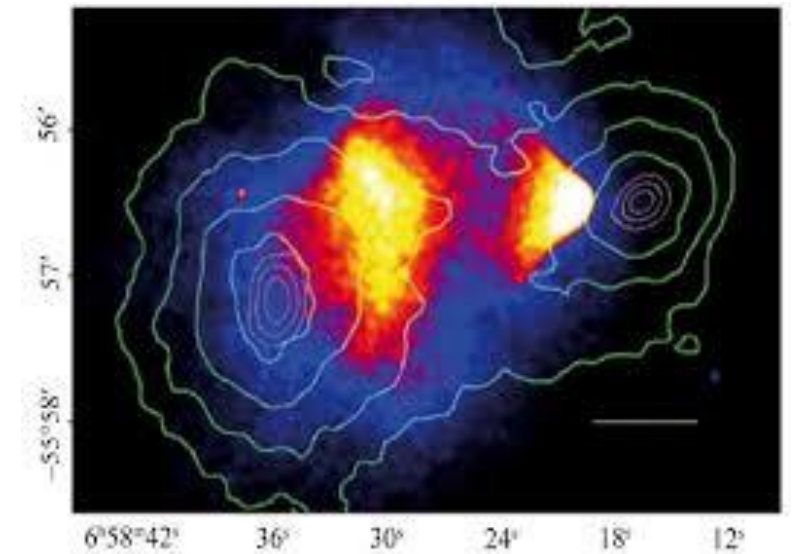
Weak lensing (statistical distortions of background galaxies)



## 4. Bullet Cluster (1E 0657-56)

**Observation:** In a collision of two galaxy clusters, the hot X-ray emitting gas (which contains most of the baryonic mass) is spatially separated from the gravitational potential inferred from lensing.

**Implication:** Demonstrates that most of the mass is collisionless—supporting a non-baryonic dark matter component.

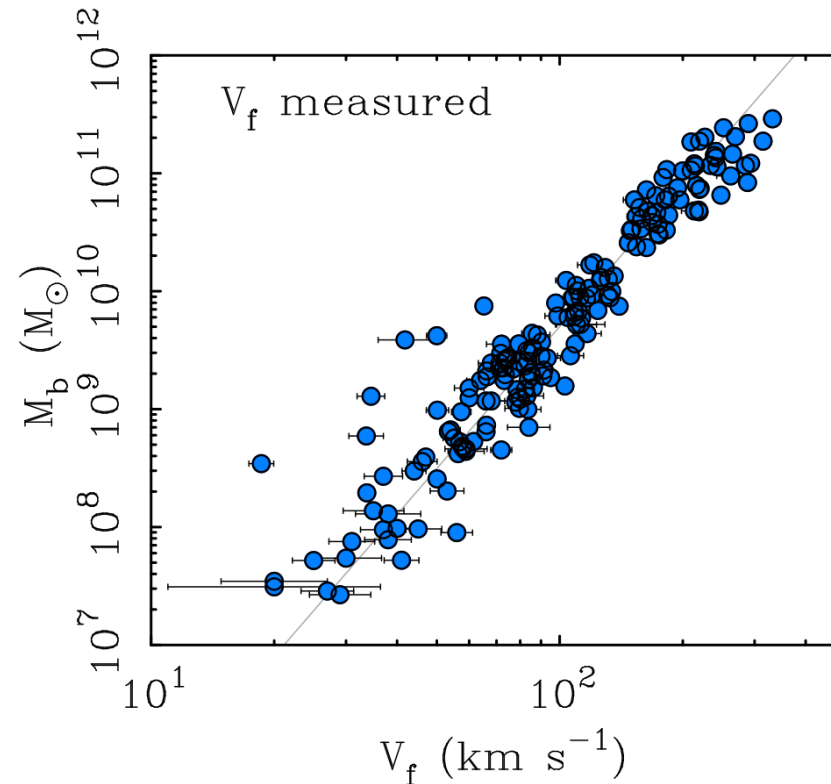




# Astrophysical Evidence

## 5. Galactic and Sub-galactic Dynamics

**Observation:** Dark matter is necessary to explain stability and properties of dwarf spheroidal galaxies, satellite galaxies, and the Tully–Fisher relation.





# Cosmological Evidence

## 6. Cosmic Microwave Background (CMB)

### Anisotropies

**Observation:** Precision measurements (e.g., by Planck and WMAP) of the temperature power spectrum reveal the relative densities of baryons, dark matter, and dark energy.

**Implication:** The height and spacing of acoustic peaks are best explained by a Universe composed of ~5% baryons, ~25% dark matter, and ~70% dark energy.

## 7. Baryon Acoustic Oscillations (BAO)

**Observation:** The imprint of primordial sound waves in the early Universe manifests as a characteristic scale in the distribution of galaxies.

**Implication:** BAO measurements agree with  $\Lambda$ CDM cosmology, requiring a substantial dark matter component.

## 8. Big Bang Nucleosynthesis (BBN)

**Observation:** The abundances of light elements (H, He, Li) match theoretical predictions only if the baryon density is much lower than the total matter density.

**Implication:** Reinforces that most matter is non-baryonic (i.e., dark).

## 9. Large-Scale Structure Formation

**Observation:** The distribution and growth of galaxies over time match simulations only if dark matter is present.

**Key Point:** Cold dark matter (CDM) is essential for early structure formation; baryons alone cannot collapse early enough.

## 9. Large-Scale Structure Formation

**CMB** sky is very **boring** –  $T$  fluctuations very **small**!

$T$  fluctuations proportional to (baryonic) **density** fluctuations,

$$\bar{\delta\rho}/\bar{\rho} \lesssim 10^{-4}$$

Matter **over-densities** in linear regime  
grow **linearly** with scale factor

But the scale factor since CMB decoupling grew by  $z_{rec} \sim 1,100$

Not enough time (since recombination) for structures to go **non-linear**!

$$\delta\rho/\rho * z_{rec} \sim 0.1, \text{ but today } \delta\rho/\rho \gg 1$$

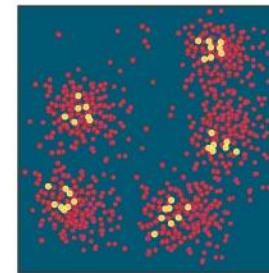
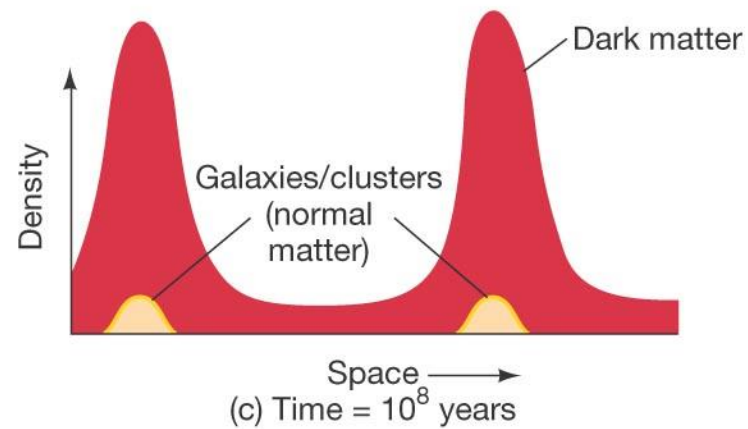
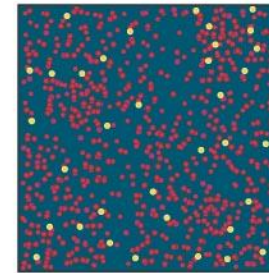
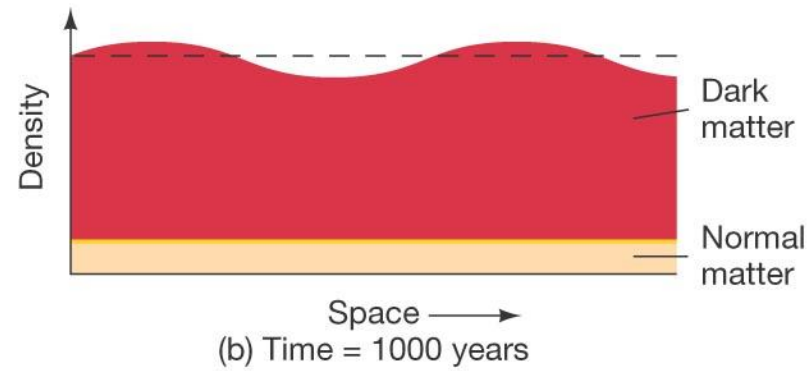
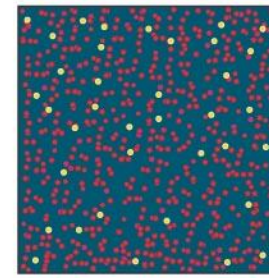
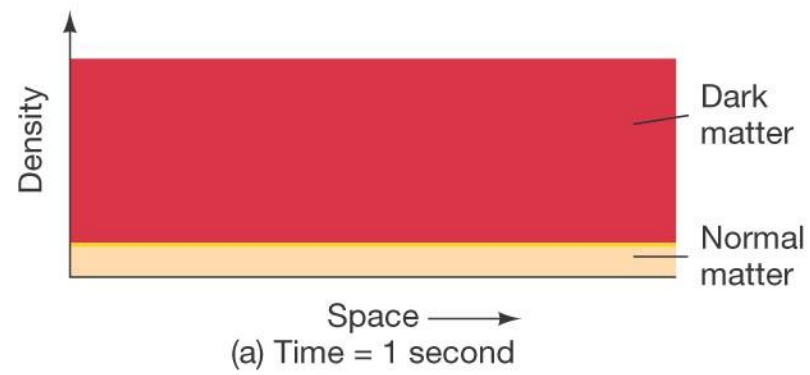
## 9. Large-Scale Structure Formation

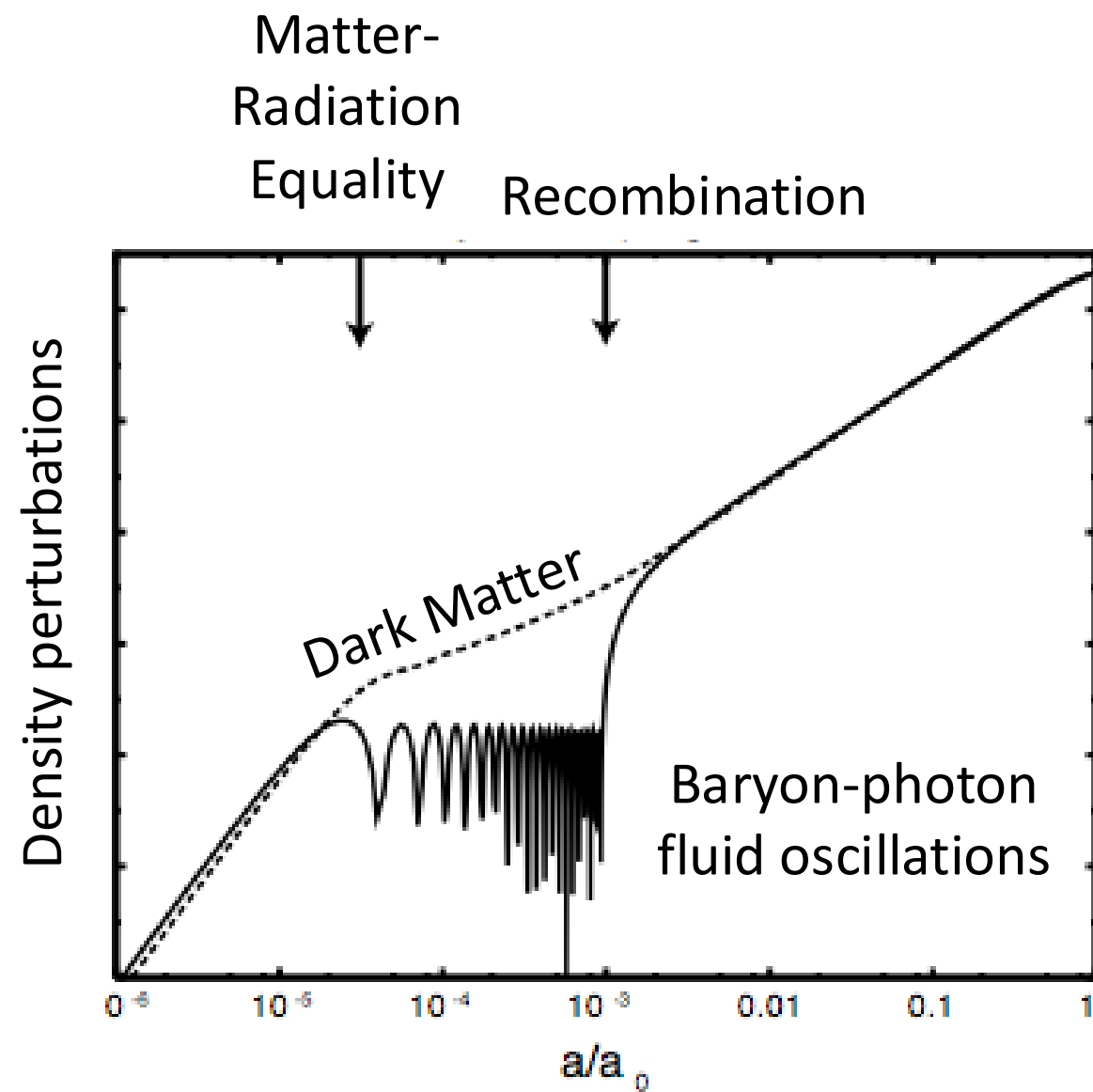
We need a **species** that has **decoupled** from photons much earlier (**Dark Matter**) so that its density **perturbations** are much **larger** at recombination!

$$(\delta\rho/\rho)_{\text{DM}} \gg 10^{-4}$$

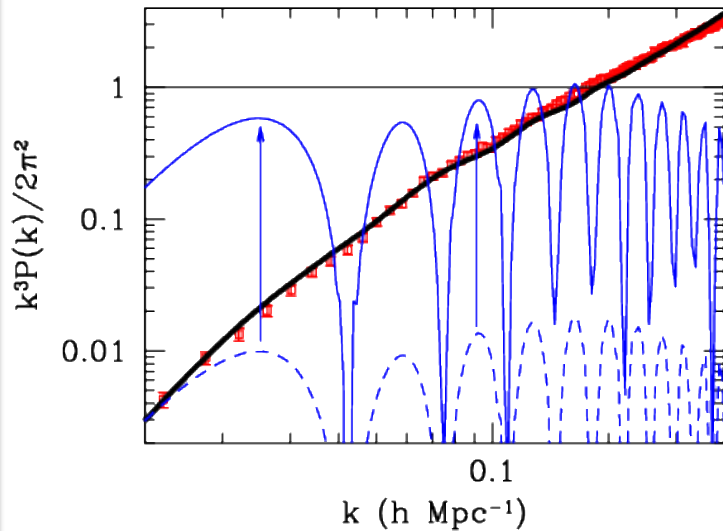
Dark matter **seeds** timely structure formation!



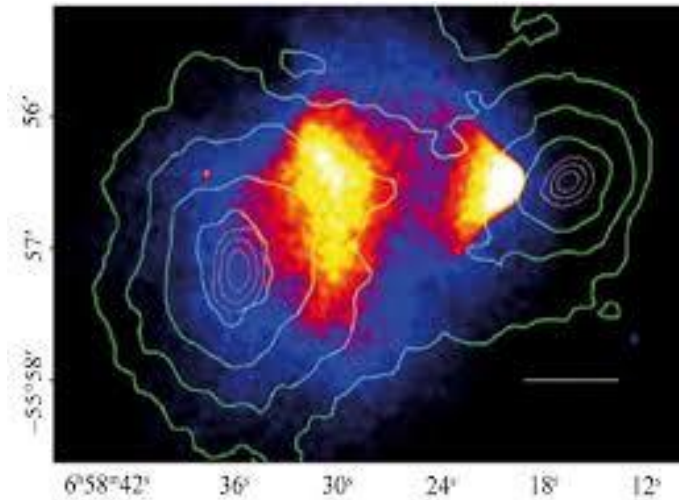




...but what if **gravity works differently** at large scales?



The Fourier Transform of **Matter Perturbations** is completely **wrong**



**Bullet cluster** observations are **impossible** to reproduce



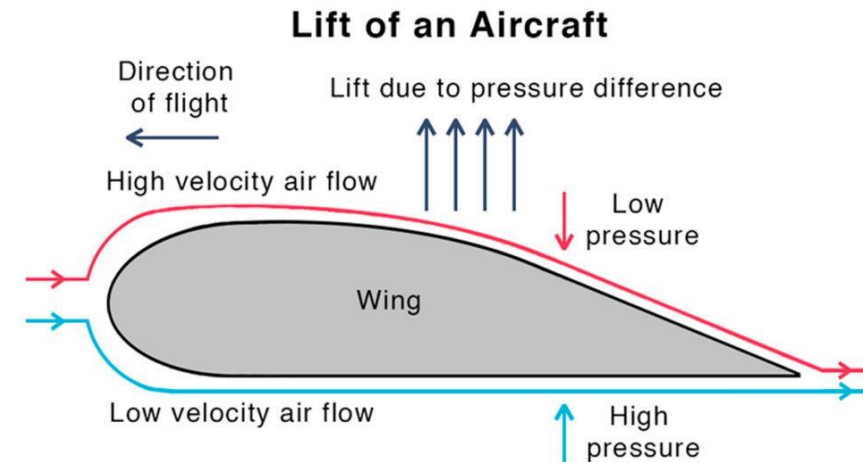
**Gravitational waves** do not propagate at the **speed of light** (they do to 1 part in  $10^{15}$ )

## ...but what if **gravity works differently** at large scales?

*“Maybe we have not formulated yet the right modified theory of gravity”*



*“Maybe we have not formulated yet the right theory of flying gnomes lifting aircrafts”*



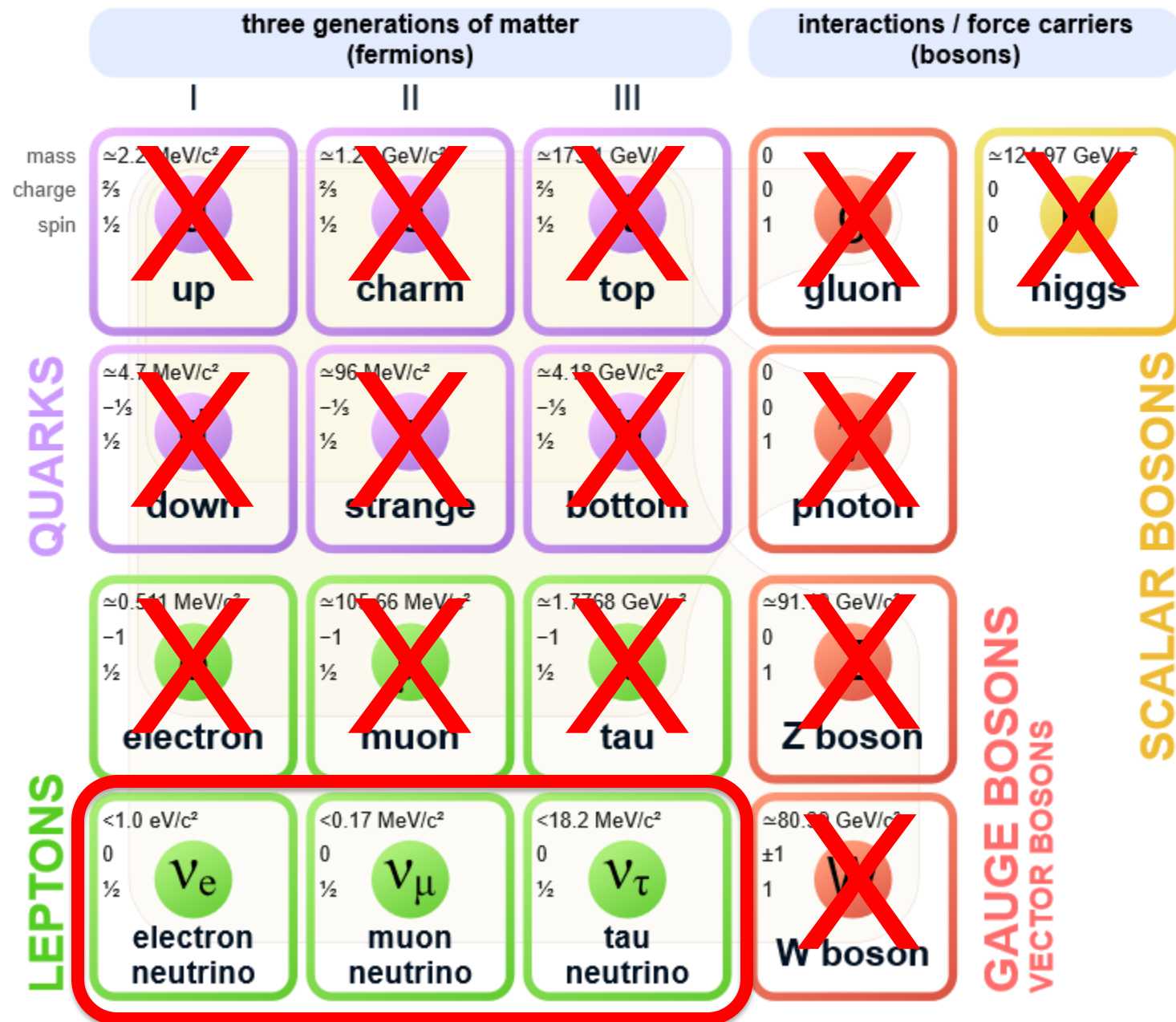
You can't **“prove”** a scientific theory,  
nor can you **rule out all alternatives**  
A scientific theory is useful if it is  
**consistent** with all data and if it makes  
**testable, falsifiable** predictions

# What do we know about the **Dark Matter**?

- **Dark** (no EM charge)
  - **Long-lived**
  - **No color charge**
- (and **B=0**)

**Cosmology + Neutrino Osc.:**  
Neutrinos are 1%  
of the dark matter

## Standard Model of Elementary Particles







What about the **99%**?

# Why Cold and not Hot Relics?

Hot relics decouple when  $T \gg m_\nu$

Structures can only collapse when  $T \sim m_\nu$

(i.e. when things slow down enough for gravitational collapse!)

Structures are cutoff to the **horizon size**  
at that temperature

$$d_\nu \sim H^{-1}(T \sim m_\nu) \quad d_\nu \sim \frac{M_P}{m_\nu^2}$$

$$H = \sqrt{\frac{\pi^2 g_*}{3 \cdot 30} \frac{T^2}{M_P}} \simeq 3.4 \frac{T^2}{M_P}$$

$$\frac{M_P^3}{m_\nu^2} \sim 10^{15} M_\odot \left( \frac{m_\nu}{30 \text{ eV}} \right)^{-2} \sim 10^{12} M_\odot \left( \frac{m_\nu}{1 \text{ keV}} \right)^{-2}$$

# Why Cold and not Hot Relics?

Observational **constraints** give

$$M_{\text{cutoff}} \ll M_{\text{Ly}-\alpha} \simeq 10^{10} M_{\odot}$$

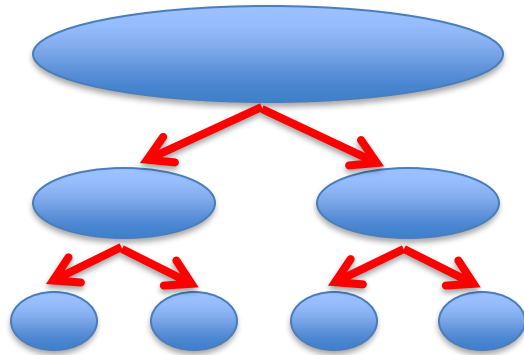
So at best dark matter can be **keV** scale, if produced thermally  
Cannot be **SM neutrinos**!!

$$\frac{M_P^3}{m_\nu^2} \sim 10^{15} M_{\odot} \left( \frac{m_\nu}{30 \text{ eV}} \right)^{-2} \sim 10^{12} M_{\odot} \left( \frac{m_\nu}{1 \text{ keV}} \right)^{-2}$$



# Why Cold and not Hot Relics?

**Structure formation** looks strikingly different  
for **hot** and **cold** dark matter

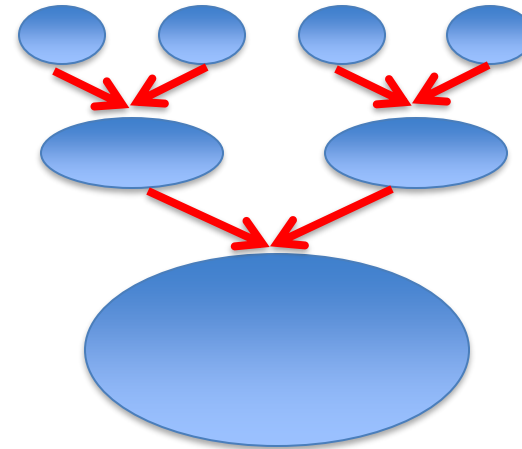


**Hot** Dark Matter

**Top-Down**

[doesn't work!]

Small structure is oldest!]



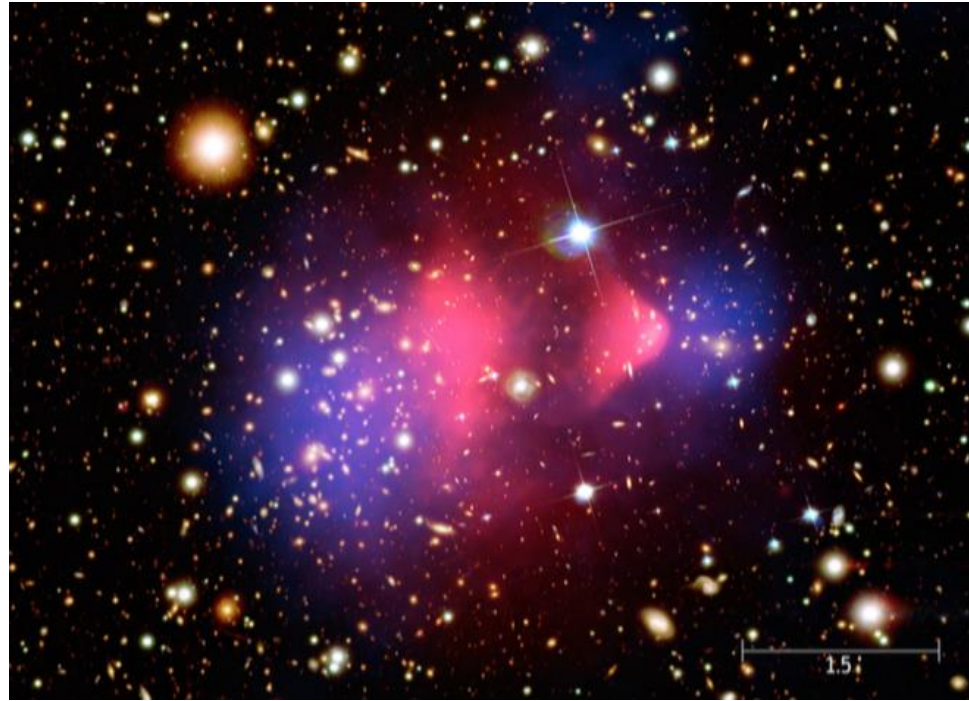
**Cold** Dark Matter

**Bottom-Up**

[Yeah!]

What else do we know about the  
**microscopic** nature of (**cold**) dark matter  
from its **macroscopic** features?

- "**Dark**" (**Milli-charge** allowed, **plasma** effects possible...)
- **~Collisionless**... really? Let's calculate the relevant constraints!

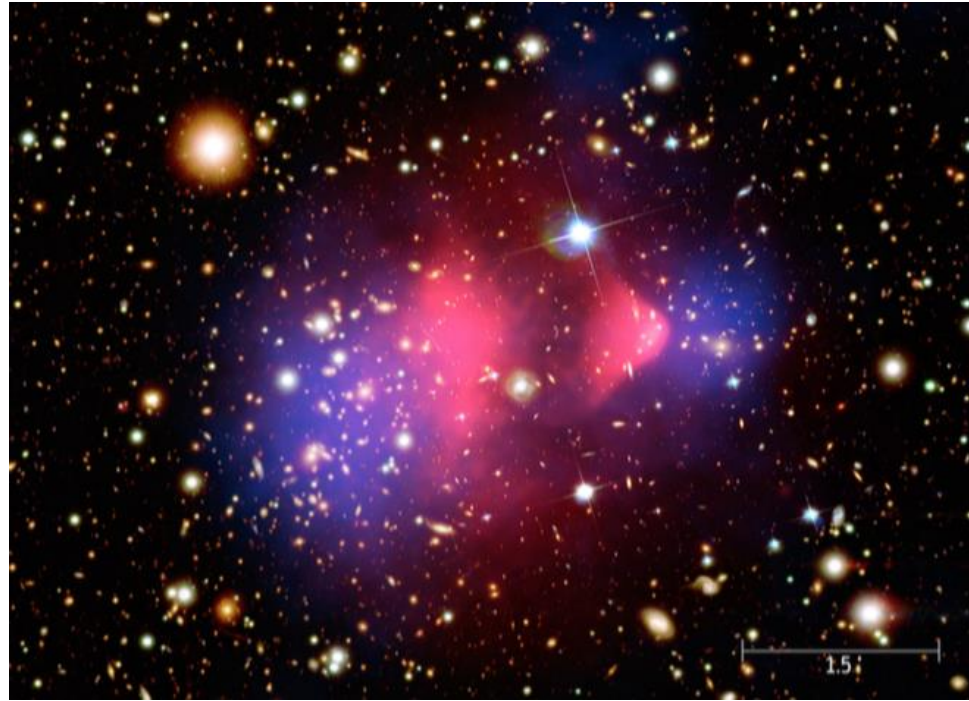


**mean free path  $\lambda$**  larger than cluster size,  $\sim 1$  Mpc

cluster **density**:  $\rho \sim 1 \text{ GeV/cm}^3$ , thus...

$$\lambda = 1/(\sigma (\rho/m)) > 1 \text{ Mpc} \rightarrow \sigma /m < 1 \text{ Mpc} / 1 \text{ GeV/cm}^3$$

$$\rightarrow \sigma /m < 1 \text{ cm}^2/\text{g}, \text{ or } 1 \text{ barn/GeV}$$



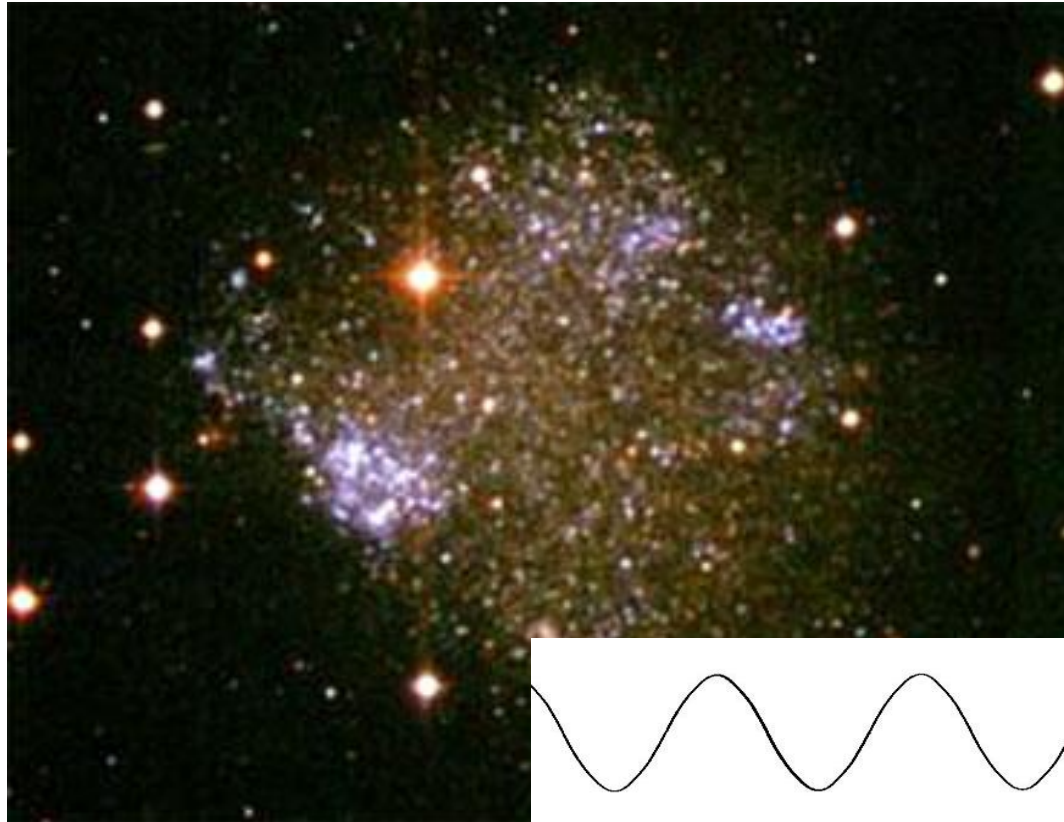
**1 barn/GeV**... which is **strong interaction**-size...

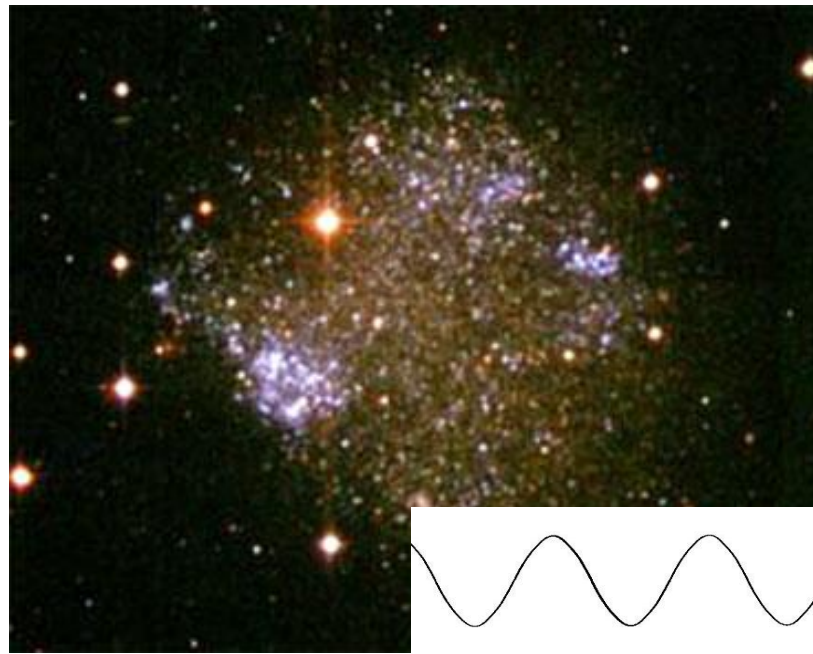
is this **small**?

Also, if cross section is **slightly smaller**, no **visible effect**...  
if cross section **slightly larger**, **disaster**...

Begs the question: is “collisional” **self-interacting** dark matter a “**natural**” possibility??

- **Classical**: quantum effects (deBroglie wavelength) must be **smaller** than **smallest** observed collapsed **structure**





little **exercise**: consider  $v \sim 10 \text{ km/s}$ , show that Compton wavelength:  $\lambda = h/p$

$$\lambda \sim 3 \text{ mm} \left( \frac{1 \text{ eV}}{m} \right)$$

which means that to have  $\lambda \ll \text{kpc} \sim 3 \times 10^{21} \text{ cm}$ ,  **$m > 10^{-22} \text{ eV}$**

Much, much **better constraints** if the DM is a fermion –  
we know that the **phase space** density is bounded  
(Pauli blocking):  $f = gh^{-3}$

Using **observed** density and velocity dispersion of dSph,  
**Tremaine-Gunn** limit (1979): observed phase space  
density cannot exceed upper bound!  
(Liouville theorem) Exercise!

$$\sigma \sim 150 \text{ km/s}$$

$$\rho \gtrsim 1 \text{ GeV/cm}^3$$

$$\left(\frac{\rho}{m}\right) \left(\frac{1}{(2\pi m^2 \sigma^2)^{3/2}}\right) < gh^{-3} \quad \longrightarrow \quad m^4 > \frac{\rho h^3}{g (2\pi \sigma^2)^{3/2}} \sim (25 \text{ eV})^4$$



➤ **Fluid**: don't want to **disrupt** pretty (and old!) **clusters** of stars

State of the art: Koulen, Profumo & Smyth 2403.1901

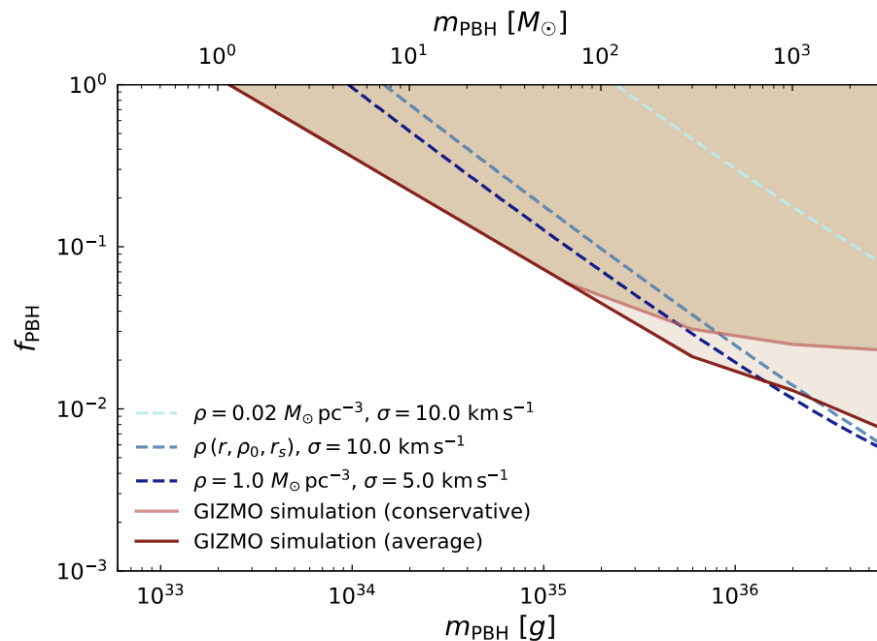


Figure 5. Constraints on  $f_{\text{PBH}}$  as a function of  $m_{\text{PBH}}$ . Limits are derived using the average  $r_h$  over 50 simulations for each mass. Also shown is a conservative case using  $r_h$  one standard deviation below the mean. The dashed lines are the semi-analytically derived constraints.

exercise to estimate the  
**exchanged** by encounters  
C and BH, in the impulse  
mation, demand that that  
y be smaller than binding  
get **maximal mass** for BH

**bility** ("heating")

Bottom line:  **$m < 10^3$  solar masses  $\sim 10^{70} \text{ eV}$**