



Lecture 4

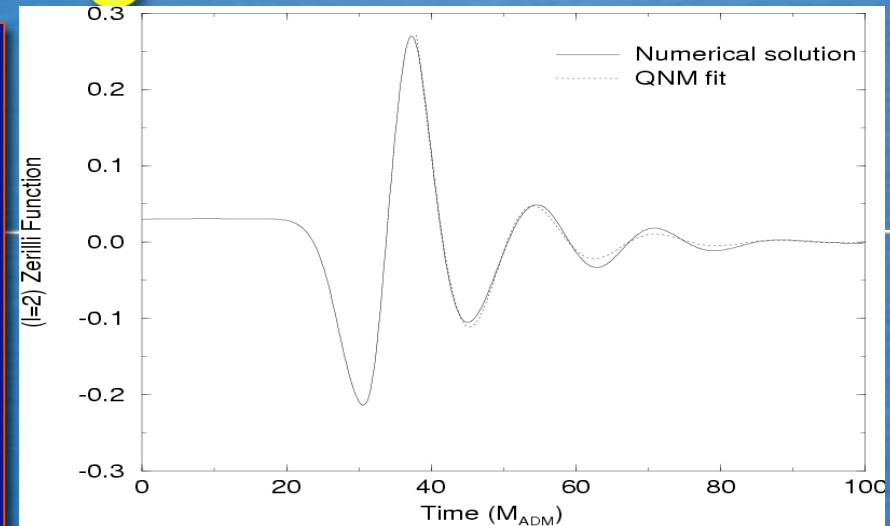
GWs from Binary Systems

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Black-Hole “Ringing”

- The newly formed BH is ringing till settles down to the stationary Kerr state (QNMs).
- The **amplitude** of the ringdown waves and their energy depends on the distortion of the BH.
- Typical frequencies: $10^4\text{-}10^5$ Hz
- **Binary system merging**
- Supermassive BHs absorbing smaller BHs or stars
- The **ringing** due to the excitation by the fallback material (after the collapse) might last for secs
- The **energy** emitted in GWs by the falling material is: $\Delta E \gtrsim 0.01 \mu c^2 (\mu/M)$.



$$\omega \approx \frac{1}{M} (0.37 + 0.19a) \approx 12 \text{kHz} \left(\frac{M_\odot}{M} \right)$$

$$\tau \approx M (1.48 + 2.09a) \approx 0.05 \text{ms} \left(\frac{M}{M_\odot} \right)$$

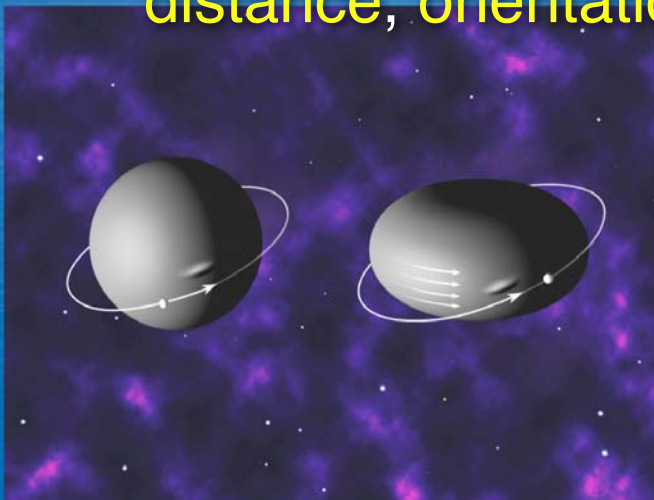
$$h_{\text{eff}} \approx 2 \times 10^{-21} \left(\frac{\varepsilon}{0.01} \right) \left(\frac{d}{10 \text{Mpc}} \right)^{-1} \left(\frac{\mu}{M_\odot} \right)$$

•GWs from BHs: a unique probe of their existence



Compact Binary Inspirals

- Late-time dynamics of compact binaries is **highly relativistic**, dictated by **non-linear GR effects**
- **Post-Newtonian theory**, which is used to model the evolution, is now known to $O(v^7)$
- The shape and strength of the emitted radiation depend on many parameters of binary system: **masses**, **spins**, **distance**, **orientation**, **sky location**, ...



Three archetypal systems

1. **Double Neutron Stars (NS-NS)**
2. **Neutron Star-Black Hole (NS-BH)**
3. **Double Black Holes (BH-BH)**



Coalescence of Compact Binaries

- During the frequency change from 100-200Hz GWs carry away $5 \times 10^{-3} M_{\odot} c^2$.
- In LIGOs band
 - NS/NS (~16000 cycles)
 - NS/BH (~3500 cycles)
 - BH/BH (~600 cycles)
- The GW amplitude is:
- Larger total mass improves detection probability.

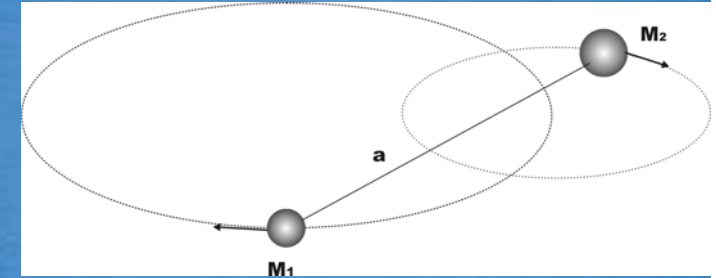
| events/y | LIGO-I | LIGO-II |
|----------|--------|----------------|
| NS/NS | ~0.05 | ~60-500 |
| BH/NS | ~0.02 | ~80 |
| BH/BH | ~0.8 | ~2000 |
| Total | 0.8 | $\gtrsim 2000$ |

- ✓ Phase effects are important, if the signal and the template get out of phase their cross correlation will be reduced.
- ✓ High accuracy templates are needed for accurate detection.

$$h \approx 7.5 \times 10^{-23} \left(\frac{M}{2.8 M_{\odot}} \right)^{2/3} \left(\frac{\mu}{0.7 M_{\odot}} \right) \left(\frac{f}{100 \text{ Hz}} \right)^{2/3} \left(\frac{100 \text{ Mpc}}{r} \right)$$

Rotating Quadrupole (a binary system)

THE BEST SOURCE FOR GWs



- Radiated power
- Energy loss leads to **shrinking of their orbital separation**
- **Period changes** with rate
- ...and the system **will coalesce** after
- The **total energy loss** is
- Typical **amplitude** of GWs

$$-\frac{dE}{dt} = \frac{32 G}{5 c^5} \mu^2 a^4 \omega^6 = \frac{32 G}{5 c^5} \frac{M^3 \mu^2}{a^5}$$

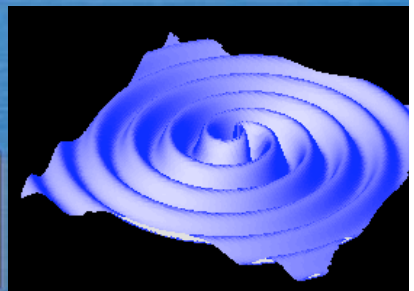
$$\frac{da}{dt} = -\frac{64 G^3 \mu M^2}{5 c^5 a^3}$$

$$\frac{\dot{P}}{P} = -\frac{96 G^3 \mu M^2}{5 c^5 a^4}$$

$$T_{\text{inspiral}} = \frac{5 c^5 a_0^4}{256 G^3 \mu M^2}$$

$$\Delta E_{\text{rad}} = \frac{G}{2} \mu M \left(\frac{1}{a_0} - \frac{1}{a} \right)$$

$$h \approx 5 \times 10^{-22} \left(\frac{M}{2.8 M_{\odot}} \right)^{2/3} \left(\frac{\mu}{0.7 M_{\odot}} \right) \left(\frac{f}{100 \text{ Hz}} \right)^{2/3} \left(\frac{15 \text{ Mpc}}{r} \right)$$





An interesting observation

The **observed frequency change** will be:

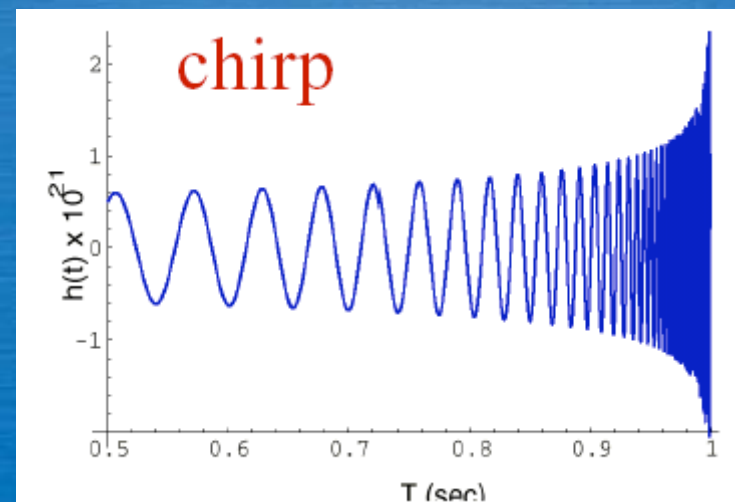
$$\dot{f} \sim f^{11/3} M_{\text{chirp}}^{5/3}$$

$$M_{\text{chirp}}^{5/3} = \mu M^{2/3}$$

The **corresponding amplitude** will be :

$$h \sim \frac{M_{\text{chirp}}^{5/3} f^{2/3}}{r} = \frac{\dot{f}}{f^3 r}$$

- ✓ Since both **frequency** and its **rate of change** are measurable quantities, we can immediately **compute the chirp mass**.
- ✓ The **third relation** provides us with a **direct estimate of the distance of the source**
- ✓ **Post-Newtonian** relations can provide **the individual masses**



Binary systems (examples)

PSR 1913+16

$$M_1 = M_2 \sim 1.4 M_\odot, P = 7\text{h } 45\text{m } 7\text{s}, r = 5\text{kpc},$$

$$h_{\text{earth}} \sim 10^{-20}, f \sim 10^{-4}\text{Hz}, T_{\text{insp}} \sim 3 \times 10^8\text{yr}$$

$$dP_{\text{theo}}/dt = -7.2 \times 10^{-12}\text{s/yr} \quad dP_{\text{obs}}/dt = -(6.9 \pm 0.6) \times 10^{-12}\text{s/yr}$$

The LIGO/VIRGO binary (10-1000Hz)

$$M_1 = M_2 \sim 1.4 M_\odot, f_0 = 10\text{Hz}, f_{\text{final}} = 1000\text{Hz},$$

$$T_{\text{insp}} \sim 15\text{min}, \text{ after } \sim 15000 \text{ cycles (inspiral/merging } 300\text{Mpc)}$$

$$M_1 = 50 M_\odot, M_2 \sim 50 M_\odot, f_0 = 10\text{Hz},$$

$$f_{\text{final}} = 100\text{Hz}, \text{ (inspiral/merging } 400\text{Mpc)}$$

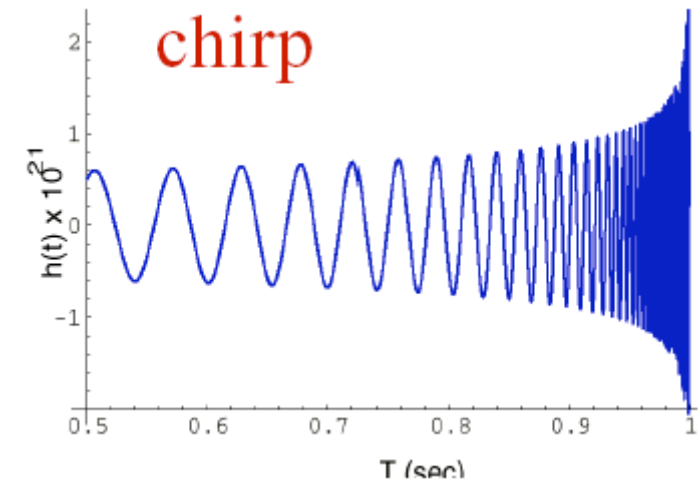
The LISA binary (10^{-5} - 10^{-2} Hz)

$$M_1 = M_2 \sim 10^6 M_\odot, f_0 = 10^{-4}\text{Hz}, f_{\text{final}} = 0.01\text{Hz}, \text{ (inspiral/merging at } r \sim 3\text{Gpc)}$$

$$M_1 = M_2 \sim 10^5 M_\odot, f_0 = 10^{-4}\text{Hz}, f_{\text{final}} = 0.1\text{Hz}, \text{ (inspiral/merging at } r \sim 3\text{Gpc)}$$

$$M_1 = M_2 \sim 10^4 M_\odot, f_0 = 10^{-3}\text{Hz}, f_{\text{final}} = 1\text{Hz}, \text{ (inspiral at } r \sim 3\text{Gpc)}$$

Smaller Stars/BHs plunging into super-massive ones



$$f_{\text{BH}} \sim 12\text{kHz} \left(\frac{1 M_\odot}{M} \right)$$





Coalescing binary signal

- The waveform **2nd PN approximation**:

$$\bar{h}(f) = N \cdot f^{-7/6} \cdot \exp \left[i \sum_{k=1}^6 \psi_k(f) \lambda_k - i \frac{\pi}{4} \right]$$

N : normalization constant

$\psi_k = \psi_k(f, f_a)$: functions of the frequency only

$\lambda_k = \lambda_k(t_c, \Phi_c, \tau_0, \tau_1, \tau_{1.5}, \tau_2)$: signal parameters

t_c, Φ_c : time and phase of coalescence

- “Newtonian case”

$$\bar{h}(f) = N \cdot f^{-7/6} \cdot \exp [2\pi i \lambda_1(t_c, \Phi_c, \tau_0) - i\pi / 4]$$

$$N \sim \frac{M^{5/6}}{R} \left[F_+^2 (1 + \cos^2 i)^2 + 4F_\times^2 \cos^2 i \right]^{1/2},$$

$$\tau_0 = \frac{5}{256} \frac{M^{-3/5}}{(\pi f_a)^{5/3}}$$

- The signal depends on **4 parameters**
 - The amplitude N
 - The phase Φ_C
 - The time parameter τ_0
 - The chirp mass $M = \mu^{3/5} M^{2/5}$

$$\Phi_I = \arctan \left[\frac{\langle h_c | s \rangle}{\langle h_s | s \rangle} \right]$$

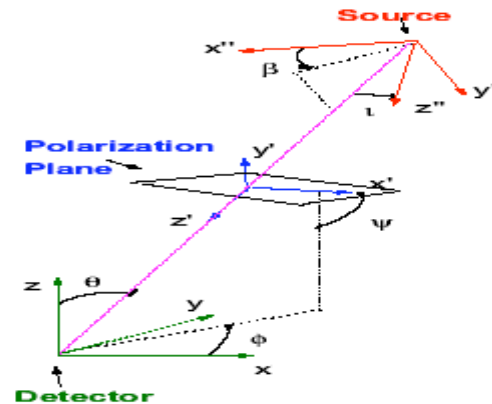
$$N = \frac{\langle h_c | s \rangle \sin \Phi_I + \langle h_s | s \rangle \cos \Phi_I}{\langle h_s | h_s \rangle}$$



Emission, Propagation and Detection

6 parameters determine the position of the binary on the sky

$r, i, \beta, \theta, \phi$ and ψ



- (x'', y'', z'') source's local frame, (i, β) polar angles
- z' -axis along wave propagation, (x', y') polarization plane, ψ angle between the x' -axis and the $\phi = 0$ plane
- (x, y, z) , detectors frame, (θ, ϕ) polar angles

The detector will feel

$$h(t) = F_+(\theta, \phi, \psi)h_+(t; i, \beta) + F_\times(\theta, \phi, \psi)h_\times(t; i, \beta)$$

F_+ & F_\times detector's beam pattern functions ($0 \leq |F| \leq 1$)

For a binary system

$$F_+(\theta, \phi, \psi) = \frac{1}{2} (1 + \cos^2 \theta) \cos 2\phi \cos 2\psi - \cos \theta \sin 2\phi \sin 2\psi$$

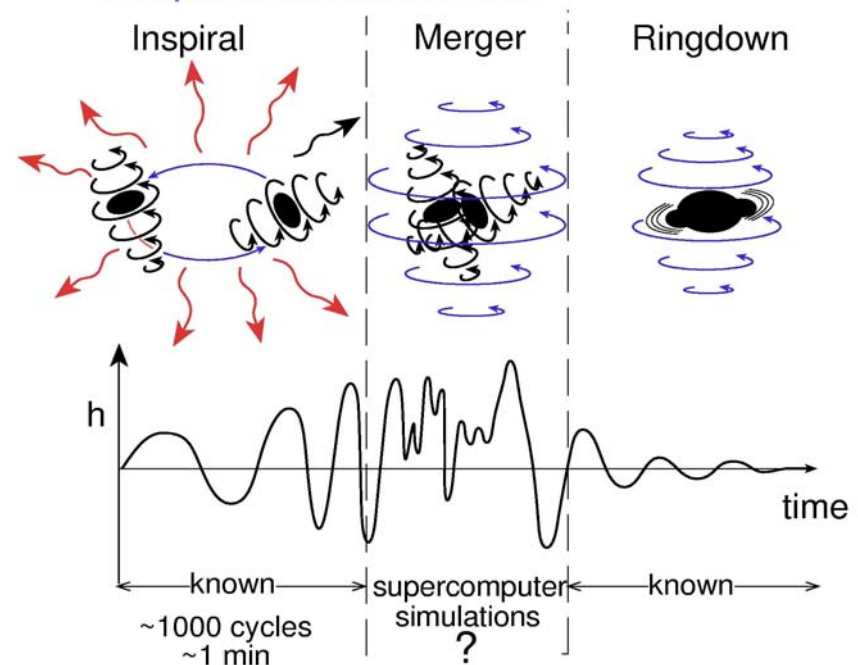
$$F_\times(\theta, \phi, \psi) = \frac{1}{2} (1 + \cos^2 \theta) \cos 2\phi \sin 2\psi + \cos \theta \sin 2\phi \cos 2\psi$$

Gravitational Waves from Binaries

Generically, there are **3 regimes** in which binaries radiate:

- **Orbital in-spiral:** PN- approximations or point-particle orbits.
- **Plunge/merger** after the last stable orbit: numerical simulations or point-particle orbits.
- **Ring-down** of the disturbed black hole as it settles down to a Kerr hole: perturbation theory of black holes.

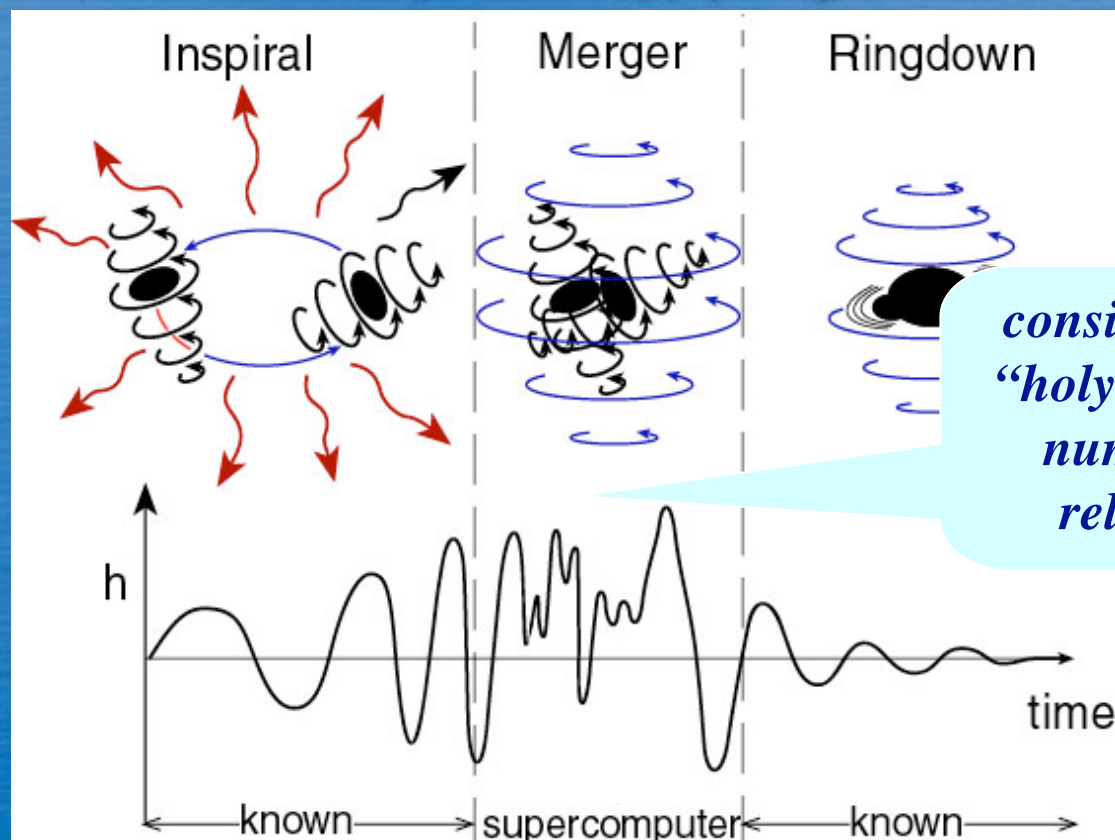
- Merger Science: nonlinear dynamics of spacetime curvature





Final merger of BH binary

- Strong-field merger is brightest GW source, luminosity $\sim 10^{23} L_{\text{SUN}}$
- Requires *numerical relativity* to calculate dynamics & waveforms
- Waveforms scale with **masses**, **spins** \rightarrow apply to ground-based & LISA



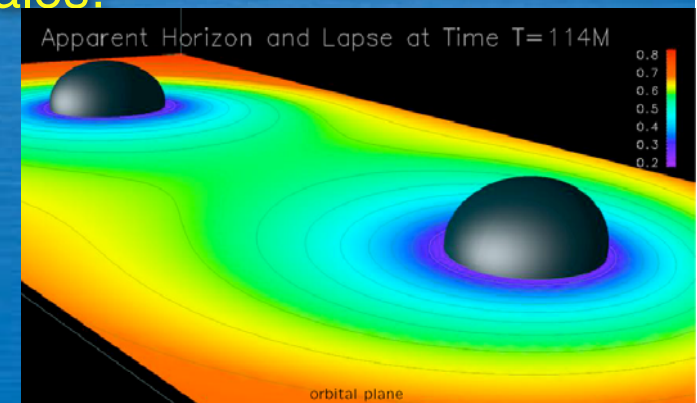
■ ■ ■ A Brief History of binary black hole simulations....

- **1964:** Hahn & Lindquist: try to evolve collision of 2 “wormholes”
- **1970s:** Smarr and Eppley: head-on collision of 2 BHs, extract GWs
 - Pioneering efforts on supercomputers at Livermore Natl Lab
- **1990s:** LIGO moves ahead & work on BBH problem starts again..
 - Work on 2-D head-on collisions at NCSA
 - NSF Grand Challenge: multi-institution, multi-year effort in 3-D
 - ***This is really difficult! Instabilities, issues in formalisms, etc...***
 - Diaspora: multiple efforts (AEI, UT-Austin, PSU, Cornell...)
 - Difficulties proliferate, instabilities arise, codes crash....
 - ***“Numerical relativity is impossible...”***
- **2000 & beyond:** LIGO/GEO/VIRGO and LISA spur more work
 - New groups: Caltech, UT-Brownsville, LSU, Jena, GSFC...
 - ***Since 2004, breakthroughs & rapid progress***
 - ***orbits, at last!***



Recent progress...on a broad front

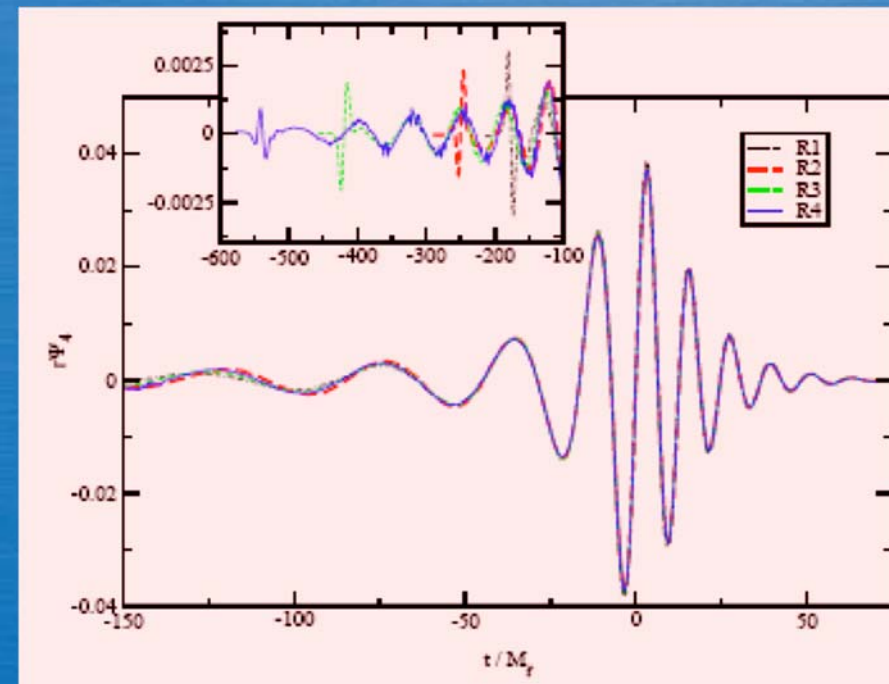
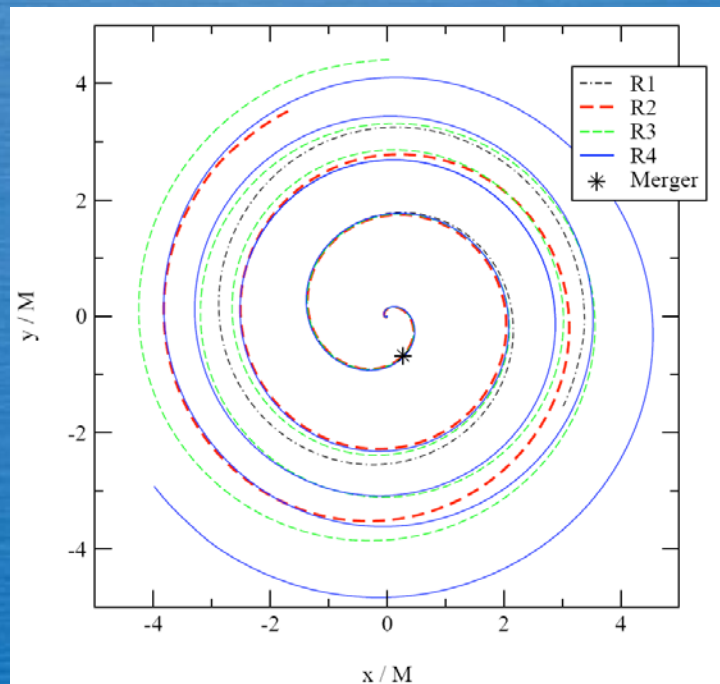
- **Evolutions of BH binary with unequal mass, spinning BHs**
 - start on approx quasi-circular orbits near last stable orbit
 - stable evolution over multiple orbits, plunge, merger, ringdown
- **Independently written codes and different software**
 - Finite differences; spectral methods
- **Different formulations of the Einstein equations**
 - 1ST & 2nd order PDEs; which variables to use; role of constraints
- **How to handle the BHs: excision; “punctures”**
- **Gauge or coordinate conditions: co-moving coords; moving BHs**
- **Variable grid resolution to handle multiple scales:**
 - $\lambda_{\text{GW}} \sim (10 - 100)M$
 - Mesh refinement; spectral decomposition





Universal waveform...

All runs agree to within $< 1\%$ for final orbit, merger & ringdown



• Baker et al 2006, Pretorius 2006, Brugman et.al 2004 (Jena), LSU-AEI 2004,
Penn-State 2005, Brownsville 2005

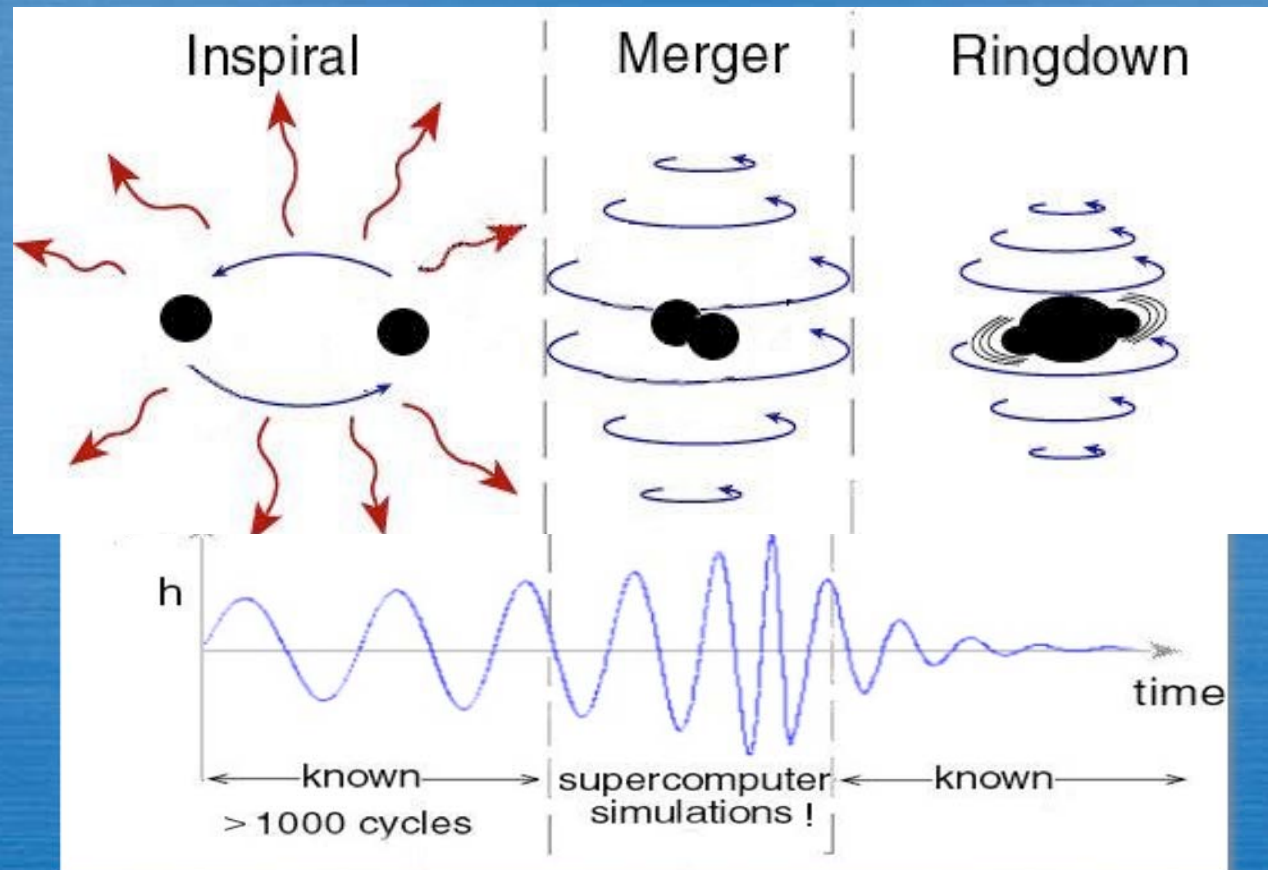
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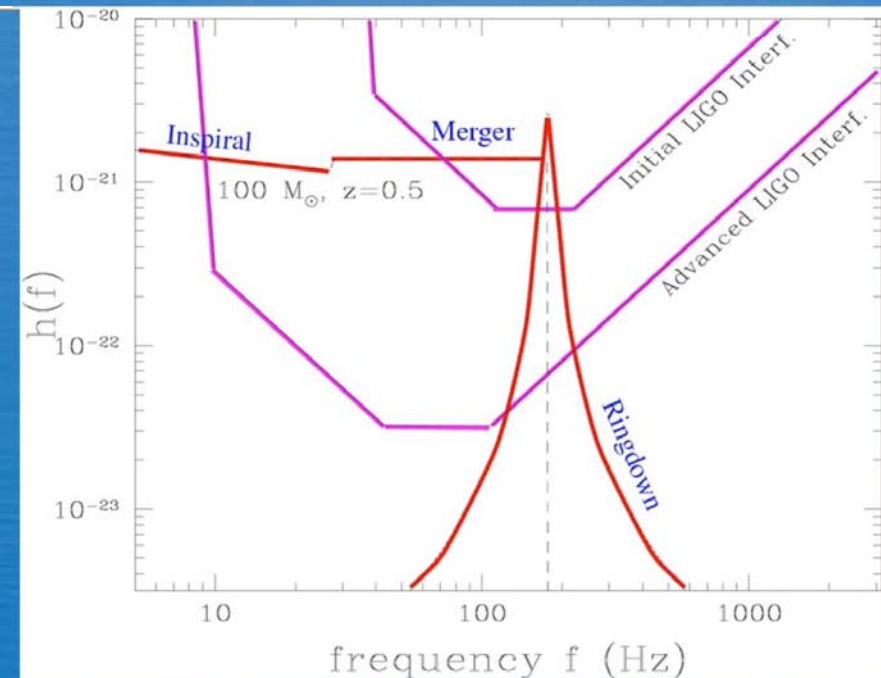
The emerging picture...

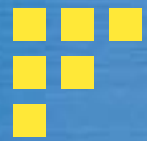




BH/BH coalescence

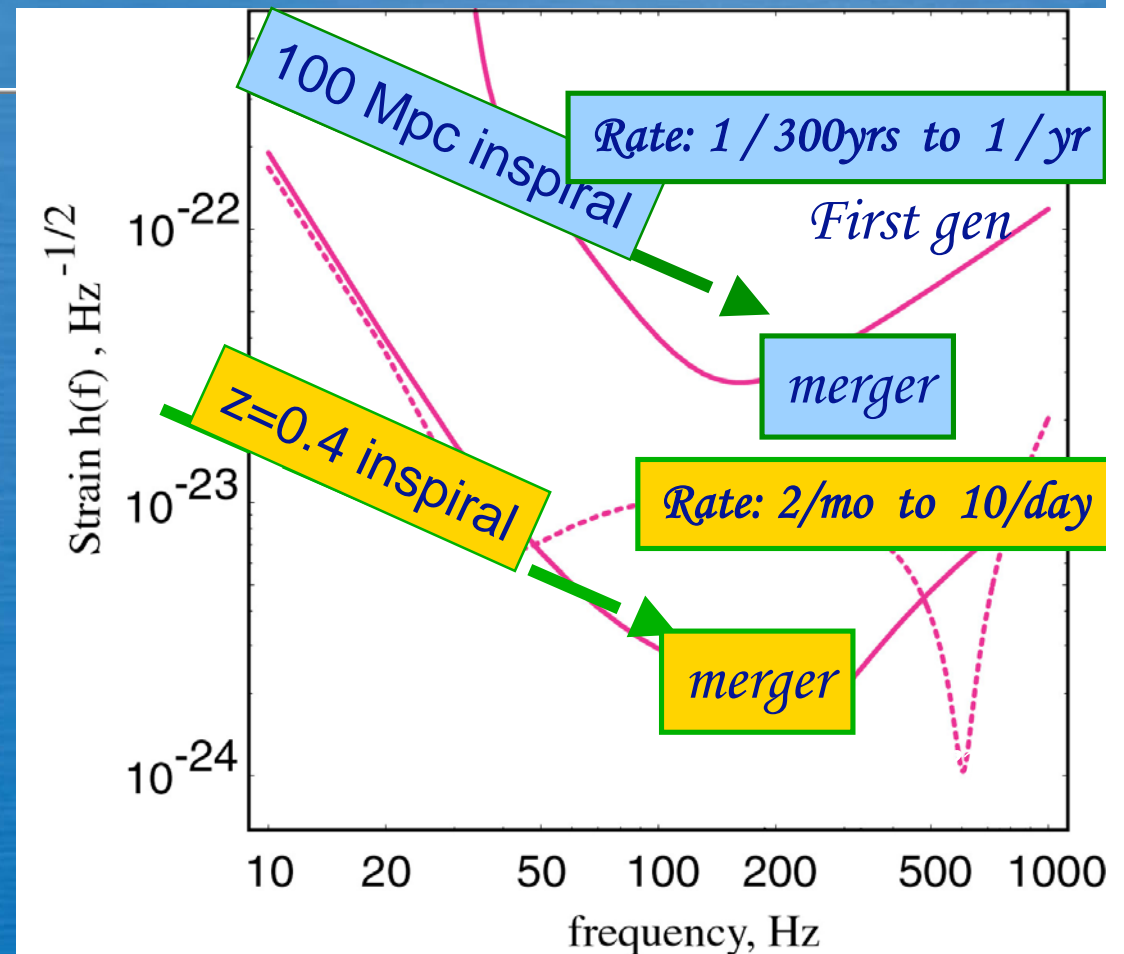
- The **inspiral**, **merger**, and **ringdown** waves from **$50M_{\odot}$ BH binaries** as observed by initial and advanced LIGO.
- The energy spectra are coming from crude estimates (**10%** of the total mass energy is radiated in merger waves and **3%** in ring-down waves).
- We observe that **the inspiral phase is not visible with initial LIGO**, for this case Numerical Relativity is important.





Possible First Source: Binary Black Hole Coalescence

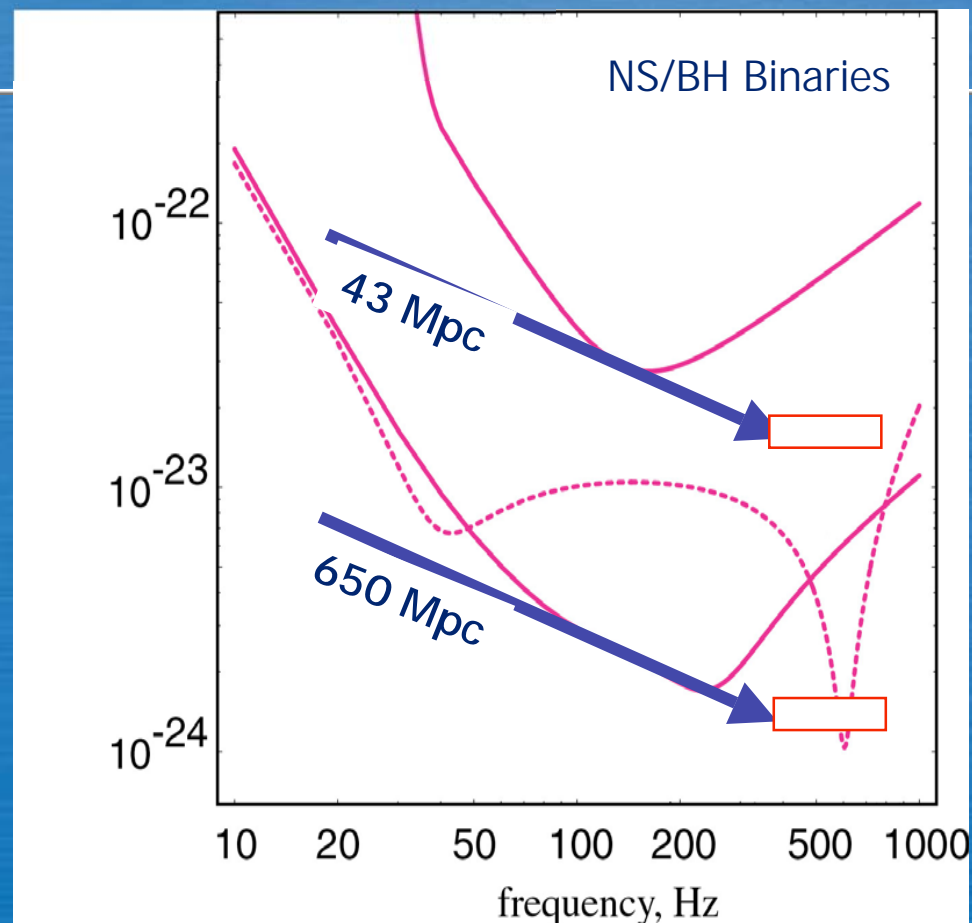
- $10M_{\odot} + 10 M_{\odot}$
BH/BH binary
- Event rates based on population synthesis,
- mostly globular cluster binaries.
- Totally quiet!!





NS-BH inspiral and NS Tidal Disruption

- **NS-BH Event rates**
 - Based on *Population Synthesis*
- Initial interferometers
 - Range: 43 Mpc
 - 1/1000 yrs to 1 per yr
- Advanced interferometers
 - Range: 650 Mpc
 - 2 per yr to several per day





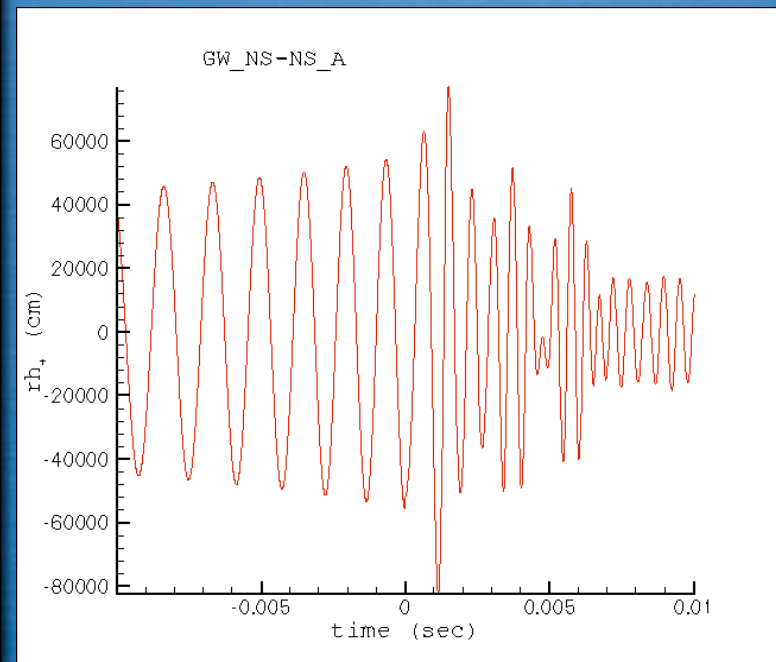
Merging phase: NS/NS & BH/NS

▪ Tidal disruption of a NS by a BH

- GWs could carry information about the EOS of NS eg. estimation of NS radius (15% error).
- The disruption waves lie in the band **300-1000Hz**
- A **few events per year** at **140Mpc** (LIGO-II)

▪ Merging of NS-NS

- Imprint of the NS radii just before merging ($f \lesssim 1\text{kHz}$)
- During the merging we could get important information about the EOS ($f \gtrsim 1\text{kHz}$)

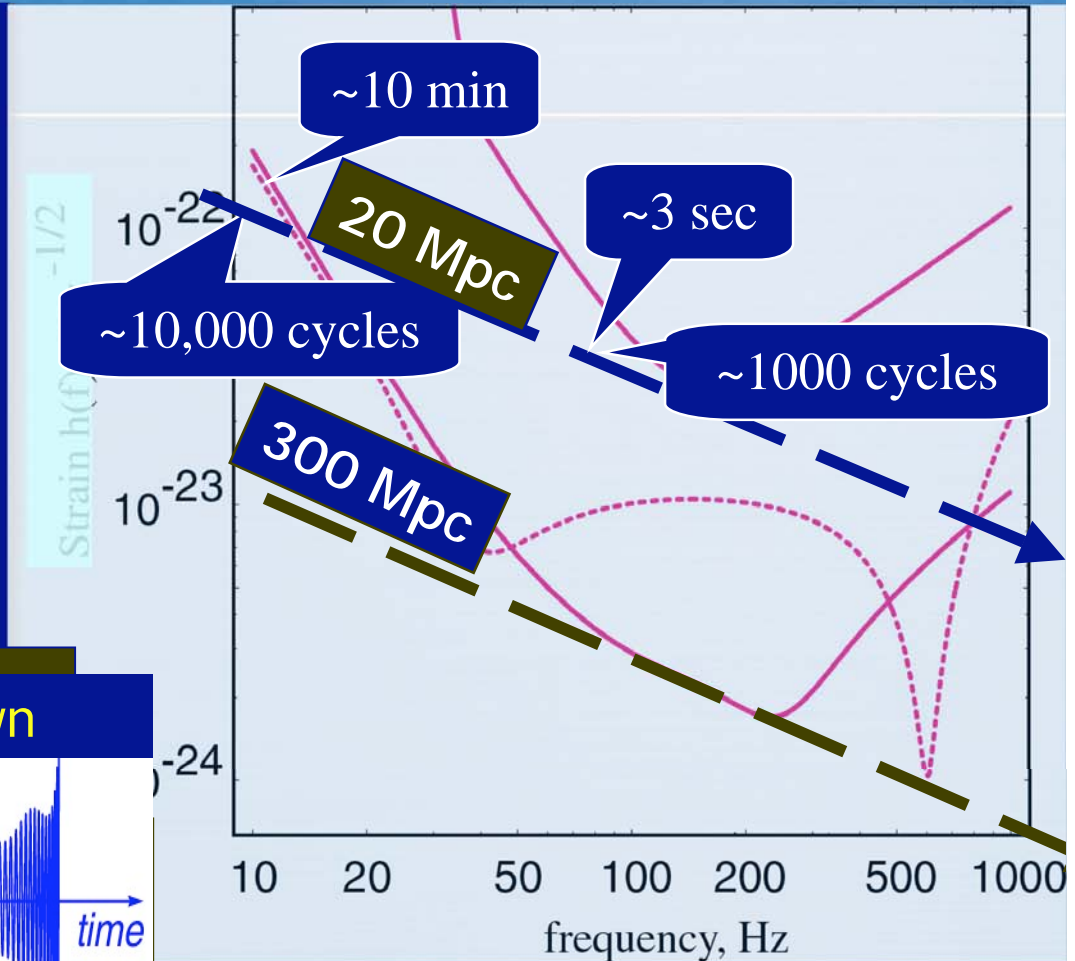




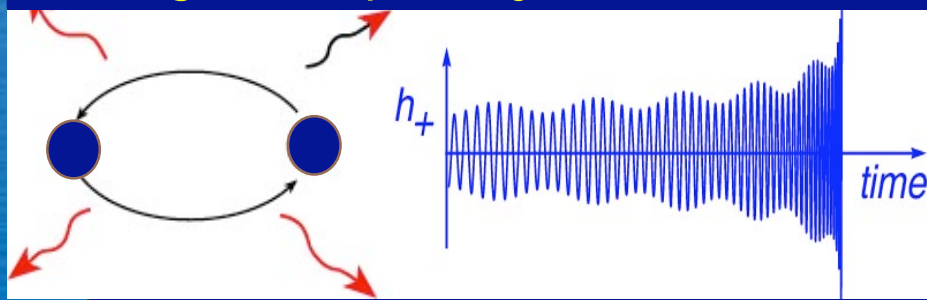
Neutron Star Binary Inspiral

NS-NS coalescence event rates

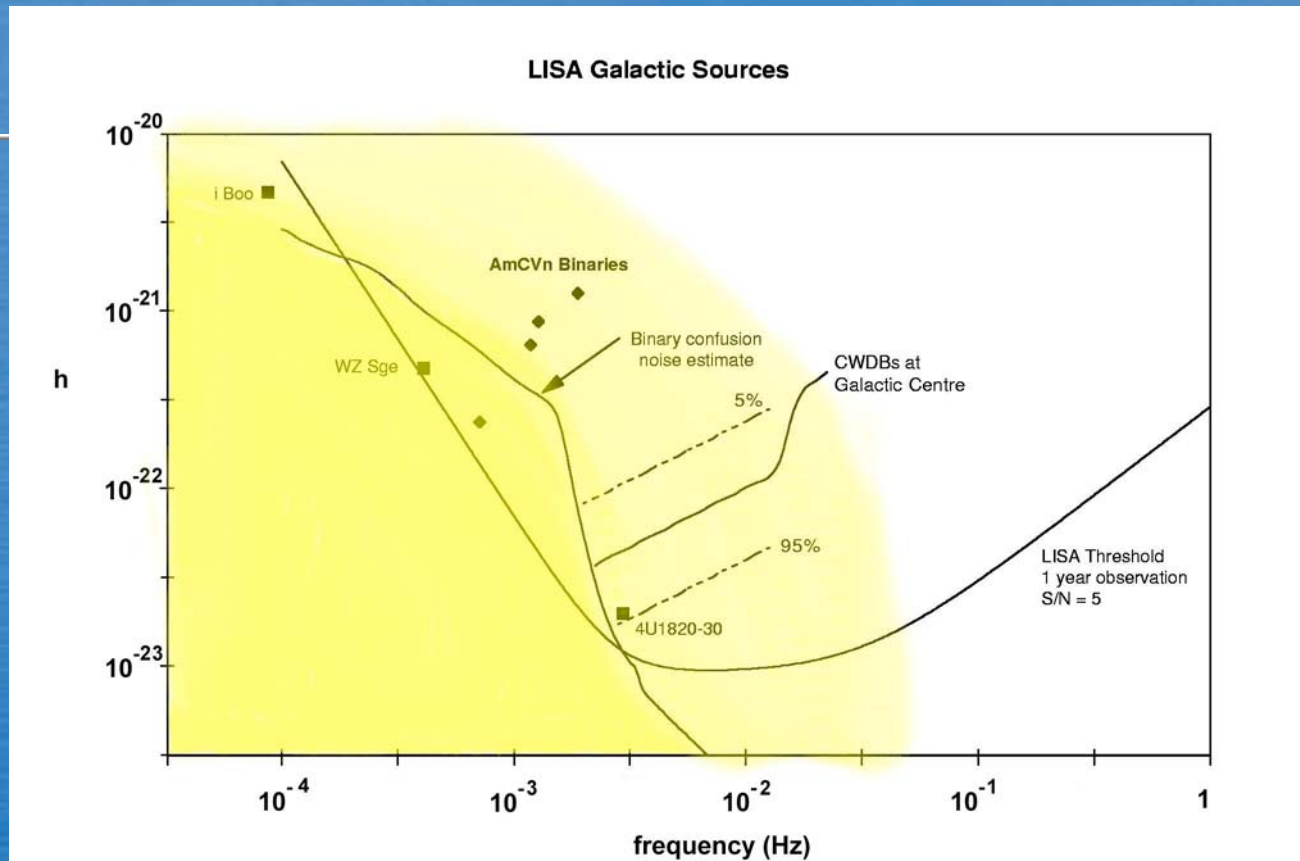
- Initial interferometers
 - Range: 20 Mpc
 - 1 per 40 yrs to 1 per 2 yrs
- Advanced interferometers
 - Range: 300Mpc
 - few per yr to several per day
- The discovery of a new binary pulsar have increased the rate upwards by an order of magnitude



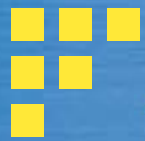
Signal shape very well known



LISA and Compact Galactic Binaries

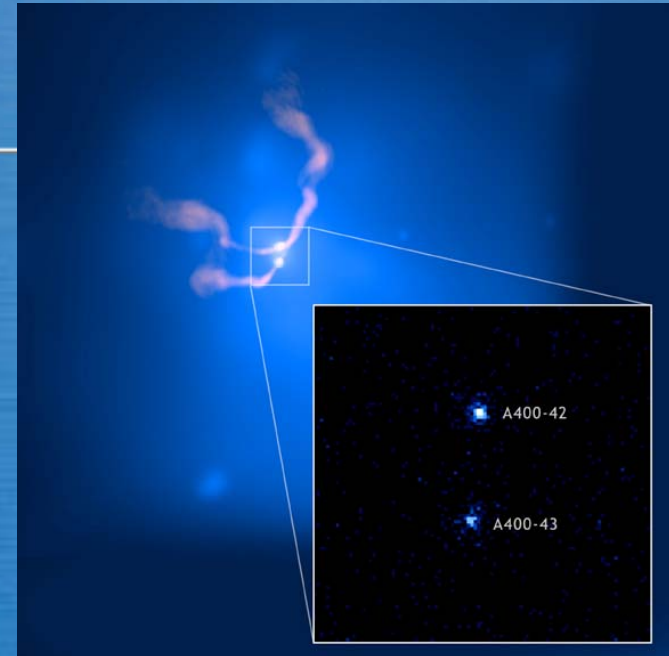


LISA will see all the compact white-dwarf and neutron-star binaries in the Galaxy.



Massive Black Holes in Galaxies

- Detected masses from 10^6 to $10^9 M_{\odot}$. Smaller masses possible.
- Galaxy mergers should produce BH mergers.
 - Rate uncertain: $1/\text{yr}$ for $10^6 M_{\odot}$ at $z=1$?
- Protogalaxy mergers may be richer. Phinney: possibly $10^3/\text{yr}$ for $10^5 M_{\odot}$ at $z = 7$.
- Stellar BHs fall into massive BHs more often, but weaker radiation.



- ***LISA can test GR in the dynamical, strong field regime...if we know the merger waveforms***
- When $m_1 \neq m_2$, GW emission is asymmetric \rightarrow recoil kick



EXAMPLE: Circular, Equatorial orbit; $10 M_{\odot} + 10^6 M_{\odot}$; fast spin

Problems:

1. Seeing them under SMBHs
2. Separating them from each other

1 yr before plunge:

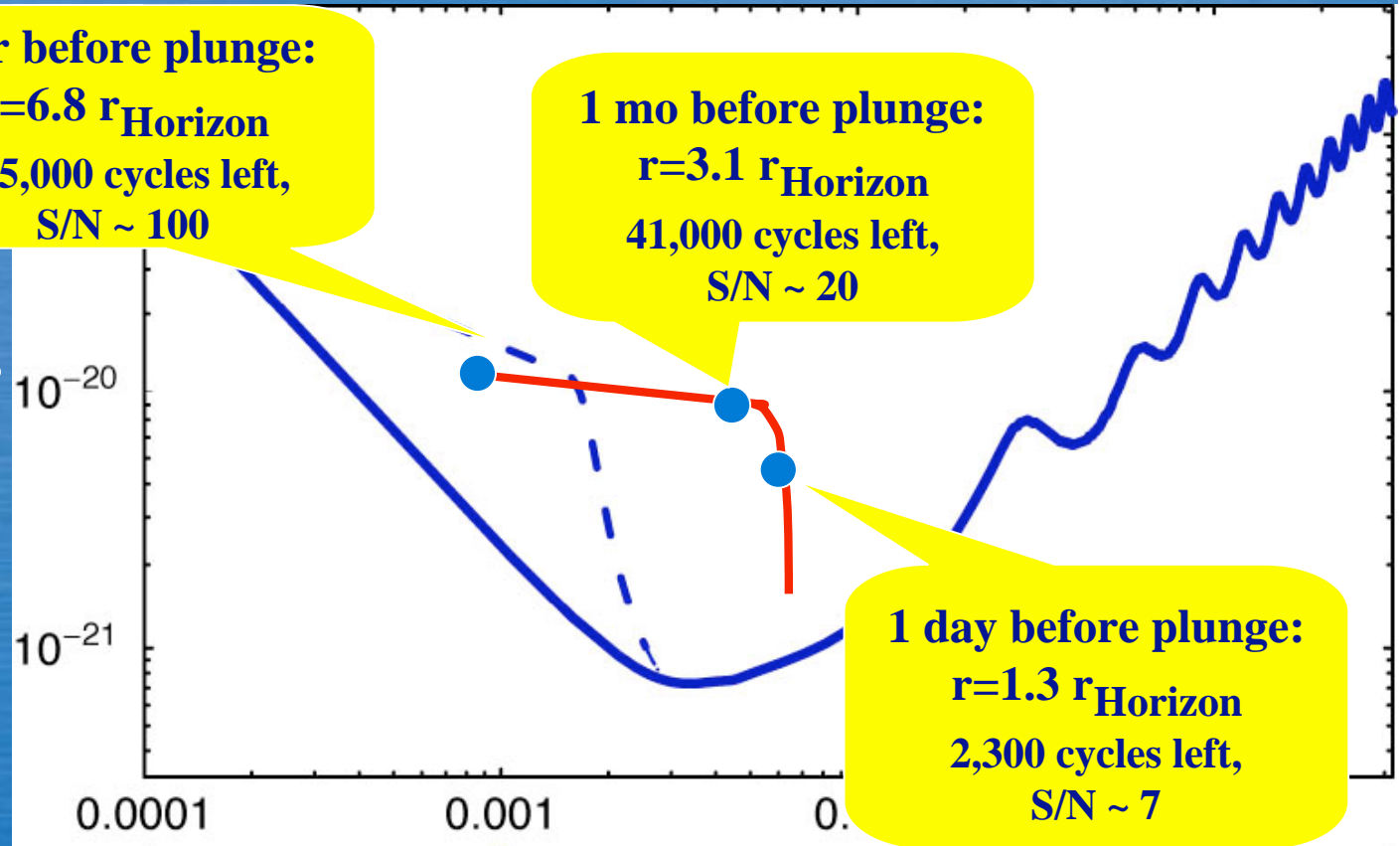
$r=6.8 r_{\text{Horizon}}$
185,000 cycles left,
S/N ~ 100

1 mo before plunge:

$r=3.1 r_{\text{Horizon}}$
41,000 cycles left,
S/N ~ 20

1 day before plunge:

$r=1.3 r_{\text{Horizon}}$
2,300 cycles left,
S/N ~ 7





Small Objects Spiral into Massive Hole

- Scatter into tight orbit

Subsequent evolution:

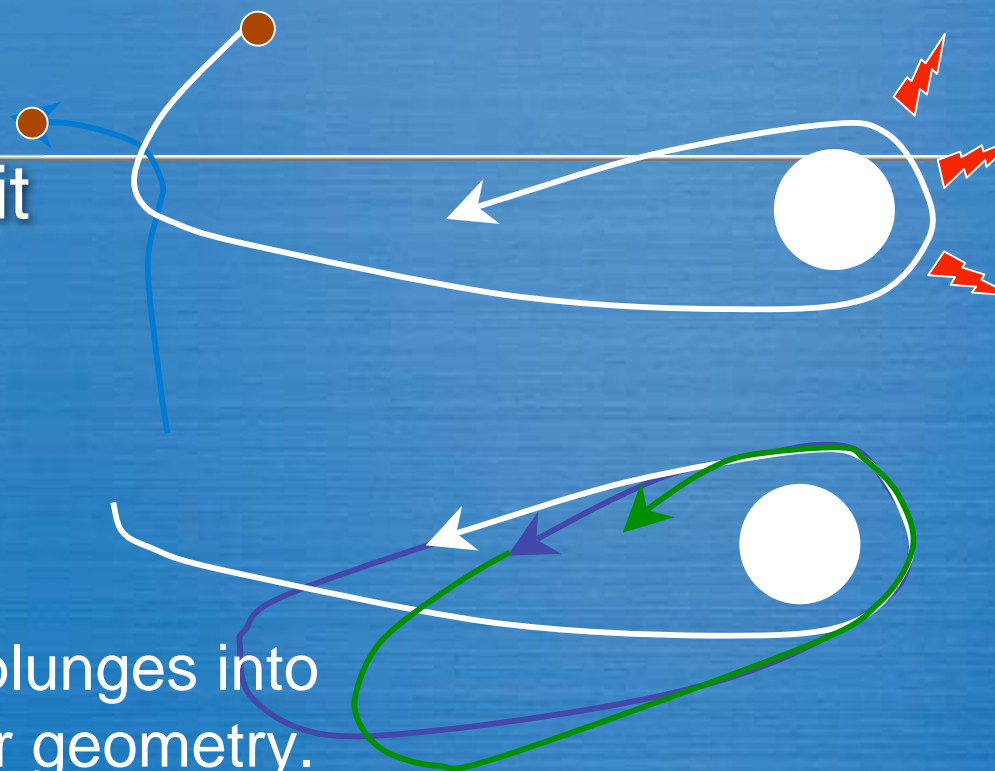
- Gradual decrease of eccentricity
- Little change of periholion
- Relativistic precession

Still quite eccentric when plunges into hole. Waveform maps Kerr geometry.

Test No-Hair Theorems.

A few per year at ~ 1 Gpc [Sigurdsson & Rees]

May be somewhat larger for $\sim 10 M_{\odot}$ holes [Phinney]





Inspirals can be seen to cosmological distances

