

Dynamic Ising Model: Reconstruction of Evolutionary Trees

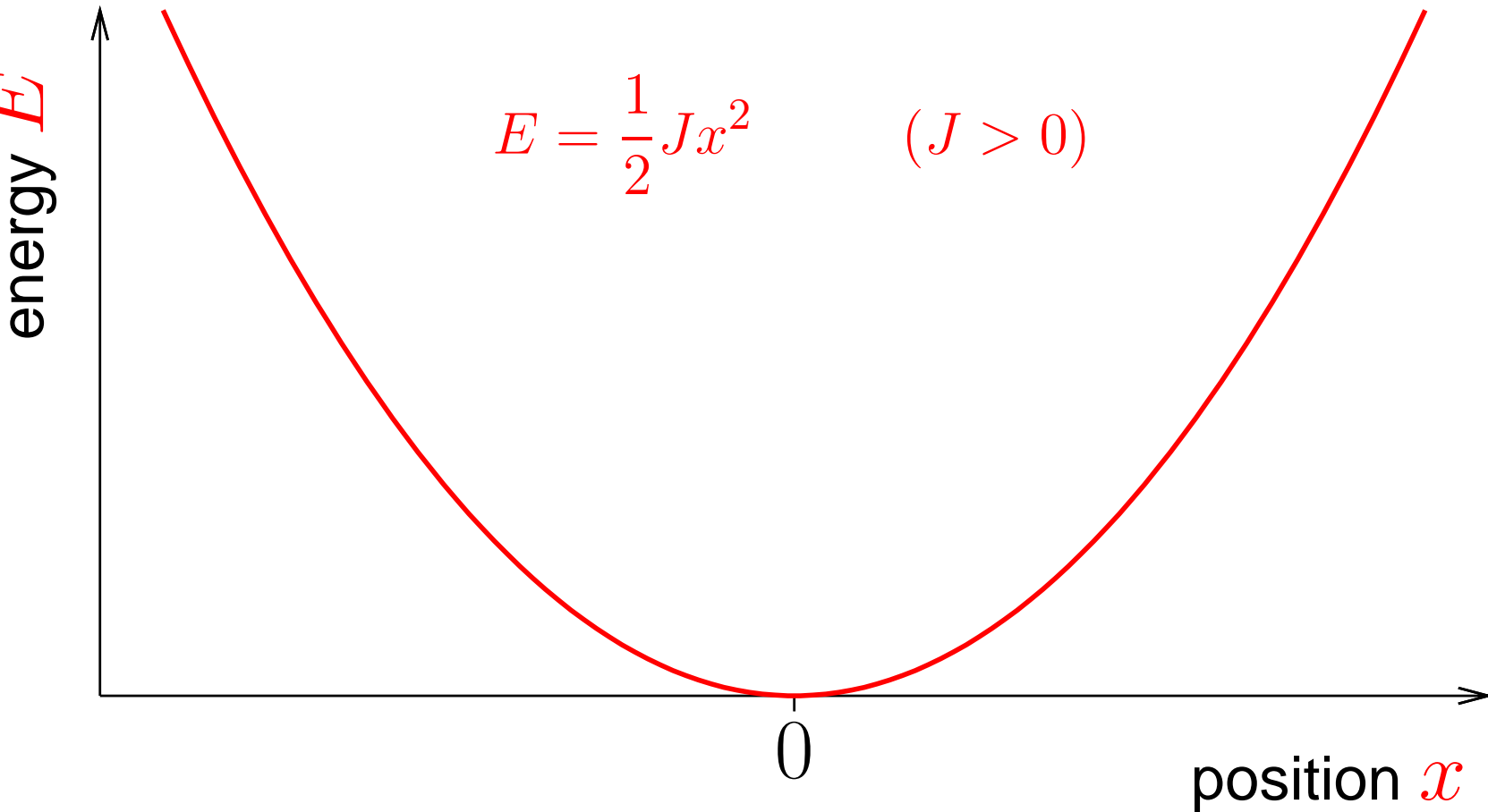
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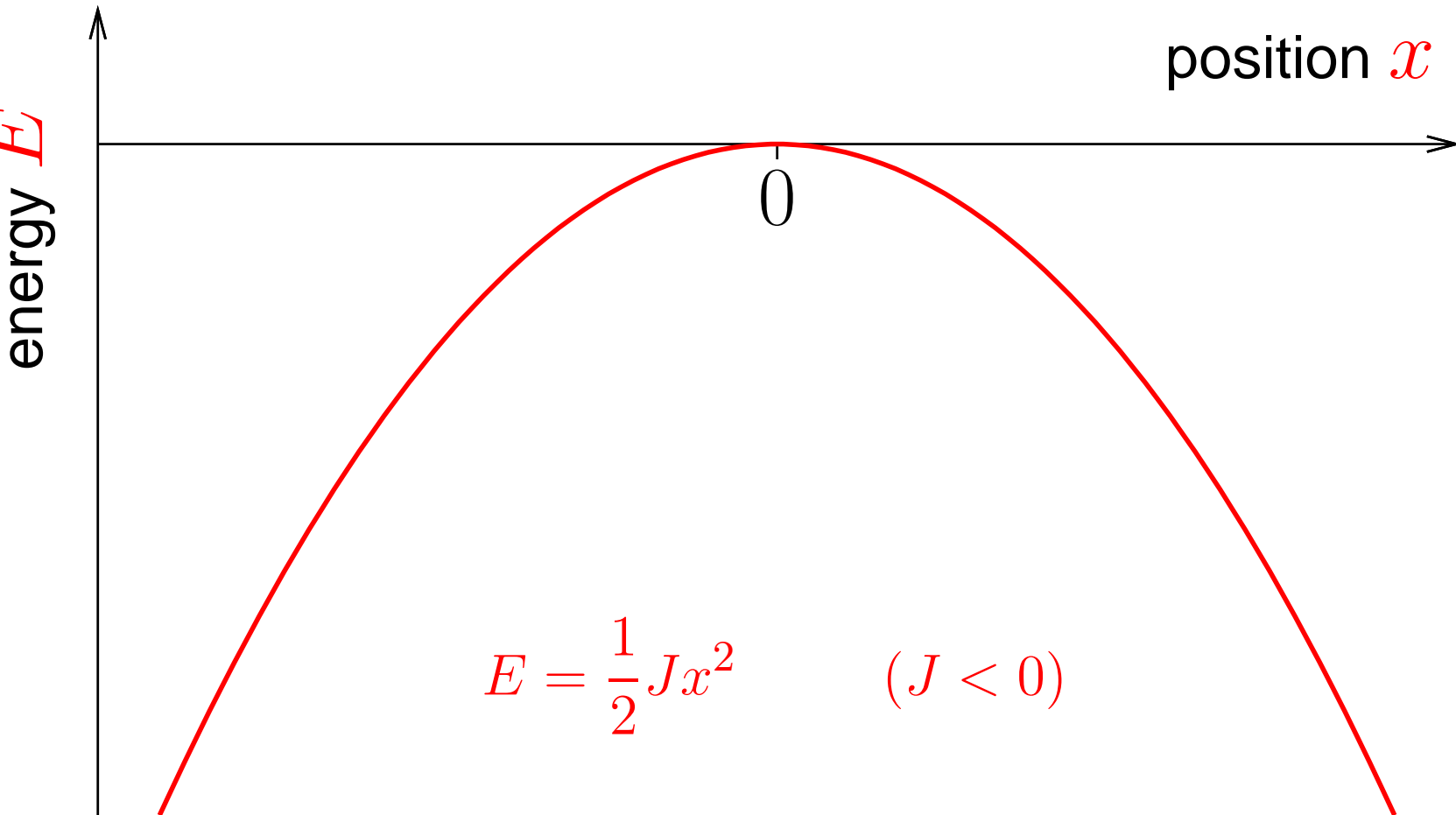
Instituto de Física, Universidade Federal Fluminense

One particle (confinement)



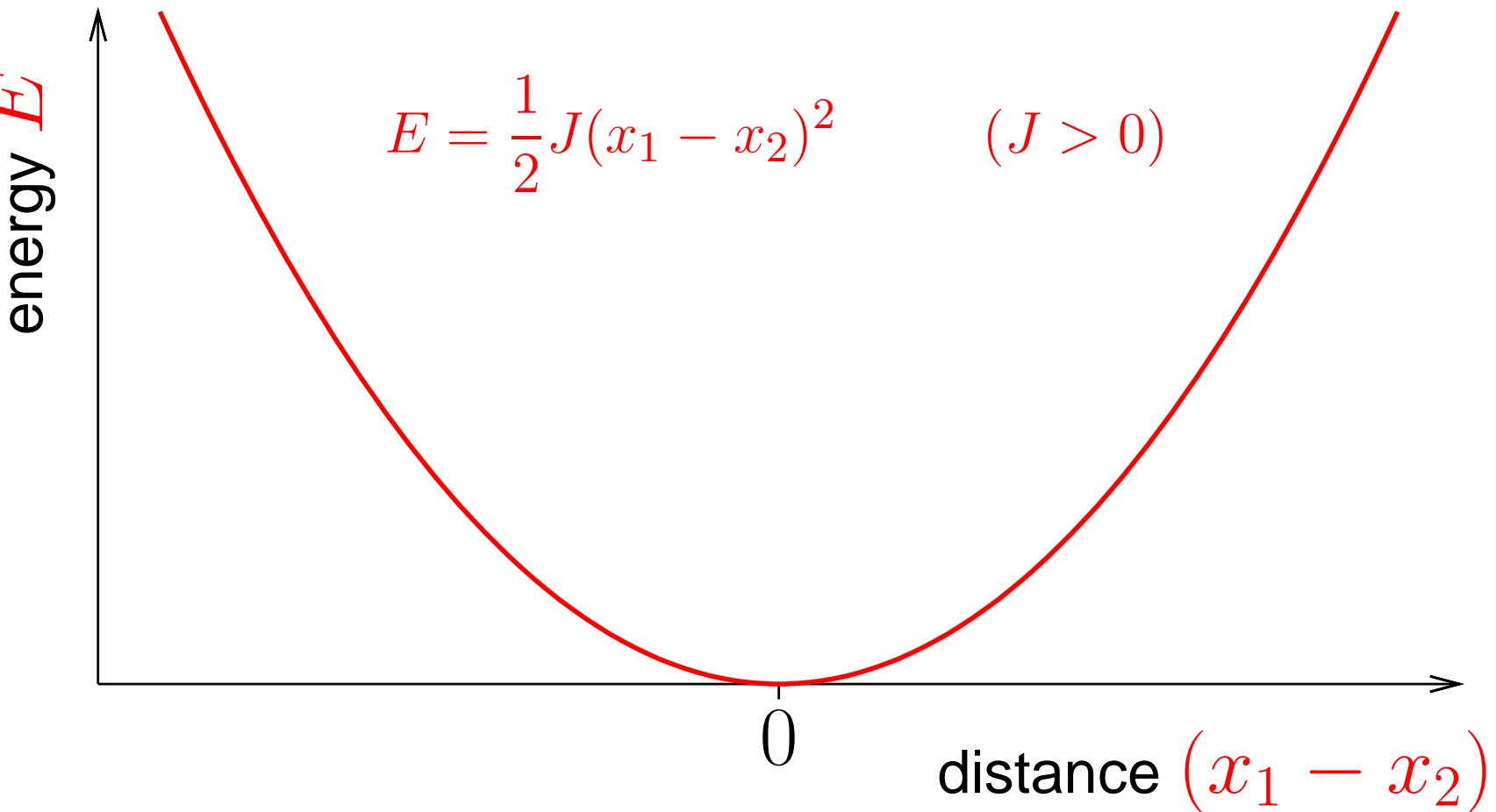
One particle tied to the origin by a spring.

One particle (runaway)



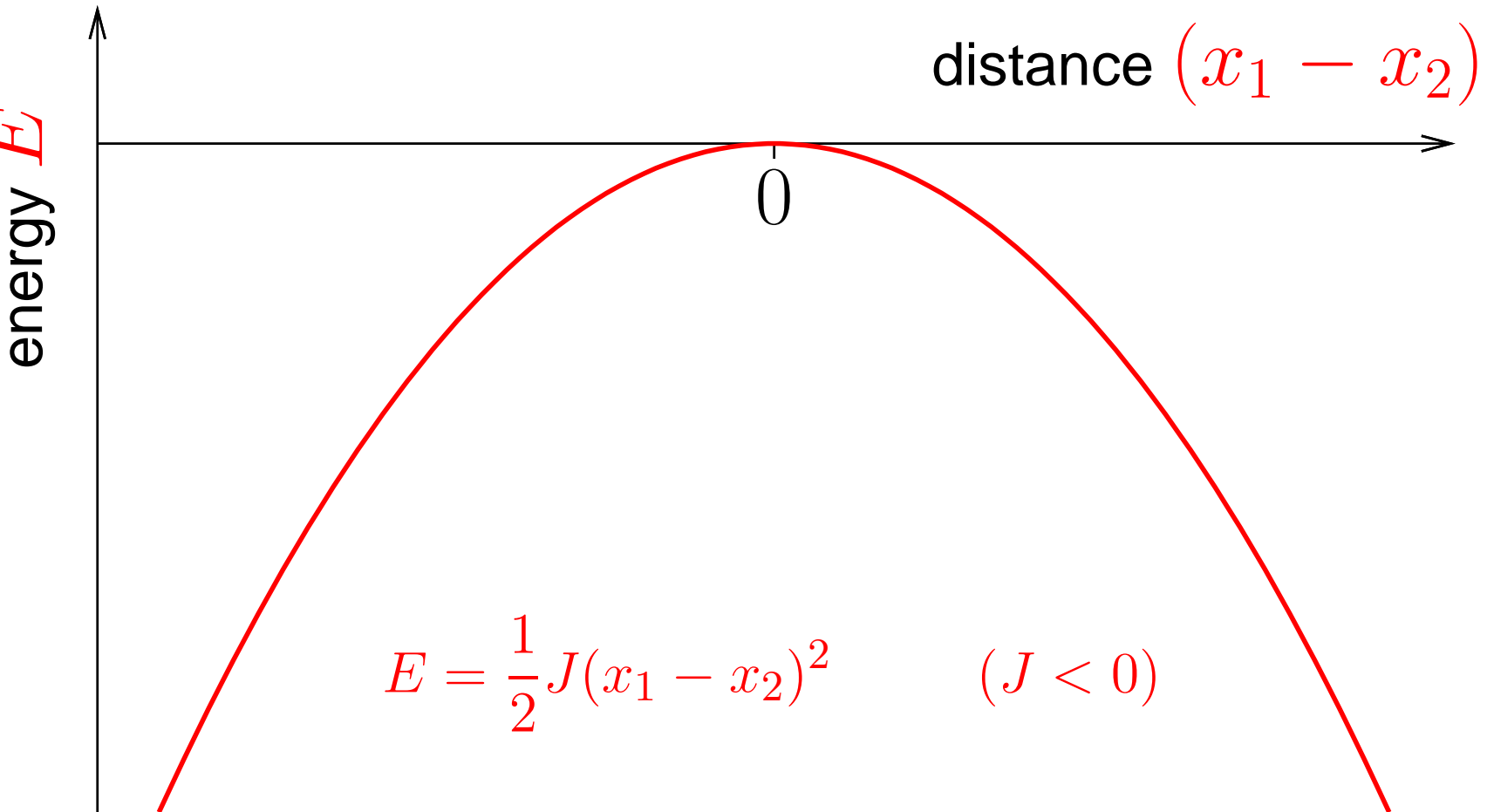
One particle repelled away from the origin by a anti-spring.

Two attracting particles



Two particles tied to each other by a spring.

Two repelling particles



Two particles repelling each other by a anti-spring.

Many (N) particles



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update :
$$\begin{cases} x_i(t + \Delta t) = x_i(t) + v_i(t)\Delta t + \frac{1}{2} \frac{F(x_i) + F(\bar{x}_i)}{2} \Delta t^2 \\ v_i(t + \Delta t) = v_i(t) + \frac{F(x_i) + F(\bar{x}_i)}{2} \Delta t \end{cases} \quad \text{Verlet}$$

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Ising (1925), most widespread model in Statistical Physics:

$$E = - \sum_{i,j} J_{ij} S_i S_j$$

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Due to universality, many different systems can be described by this model: ferromagnetism or liquid gas transition for $J_{ij} > 0$; antiferromagnetism for $J_{ij} < 0$ in bipartite lattices; spin glasses for both positive and negative, random J_{ij} ; and Boolean systems in general.

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Onsager’s exact solution (1944) for the thermodynamic behaviour of the uniform ferromagnet in two dimensions is a paradigmatic scientific achievement.

Back to the current model

$$\frac{1}{2}J_{ij}(x_i - y_j)^2 = \frac{1}{2}J_{ij}x_i^2 + \frac{1}{2}J_{ij}x_j^2 - J_{ij}x_ix_j$$

Back to the current model

“external” potential

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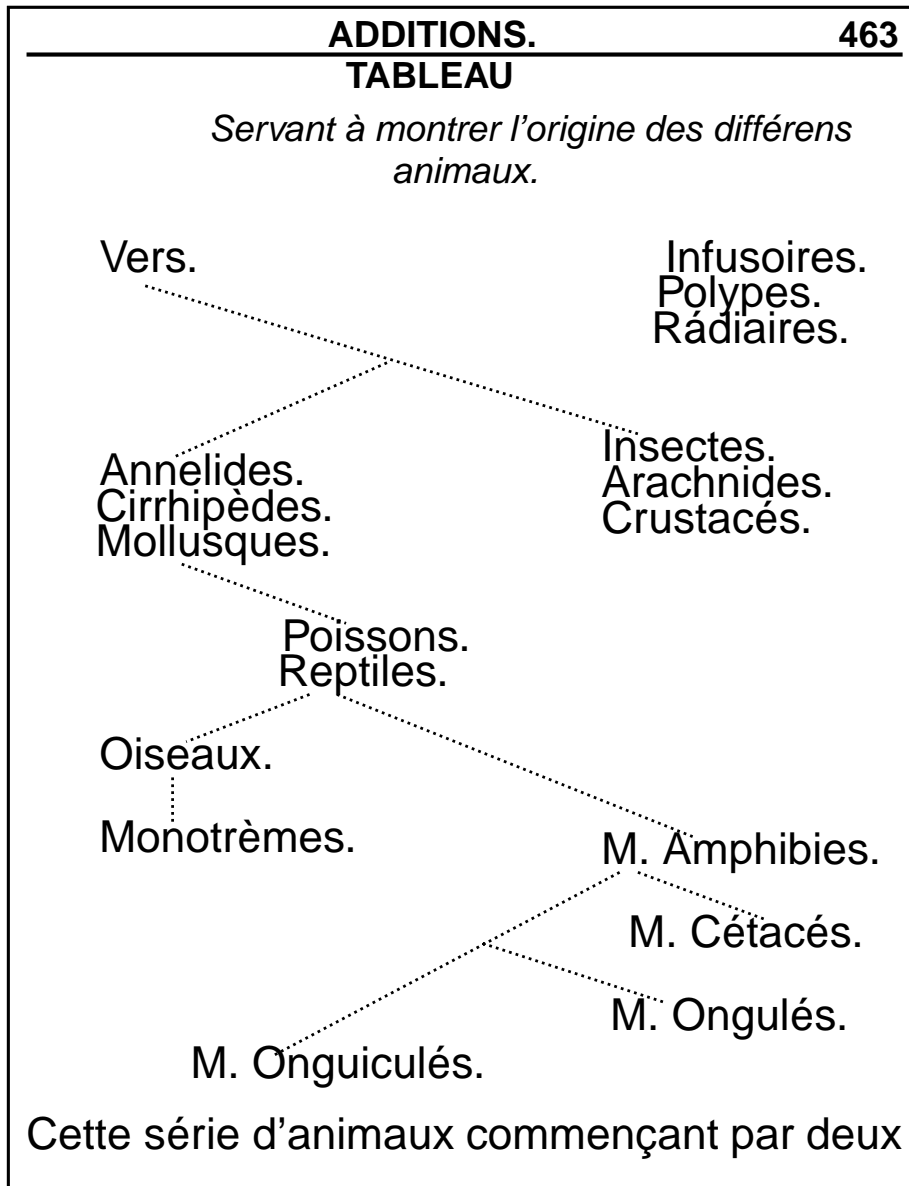
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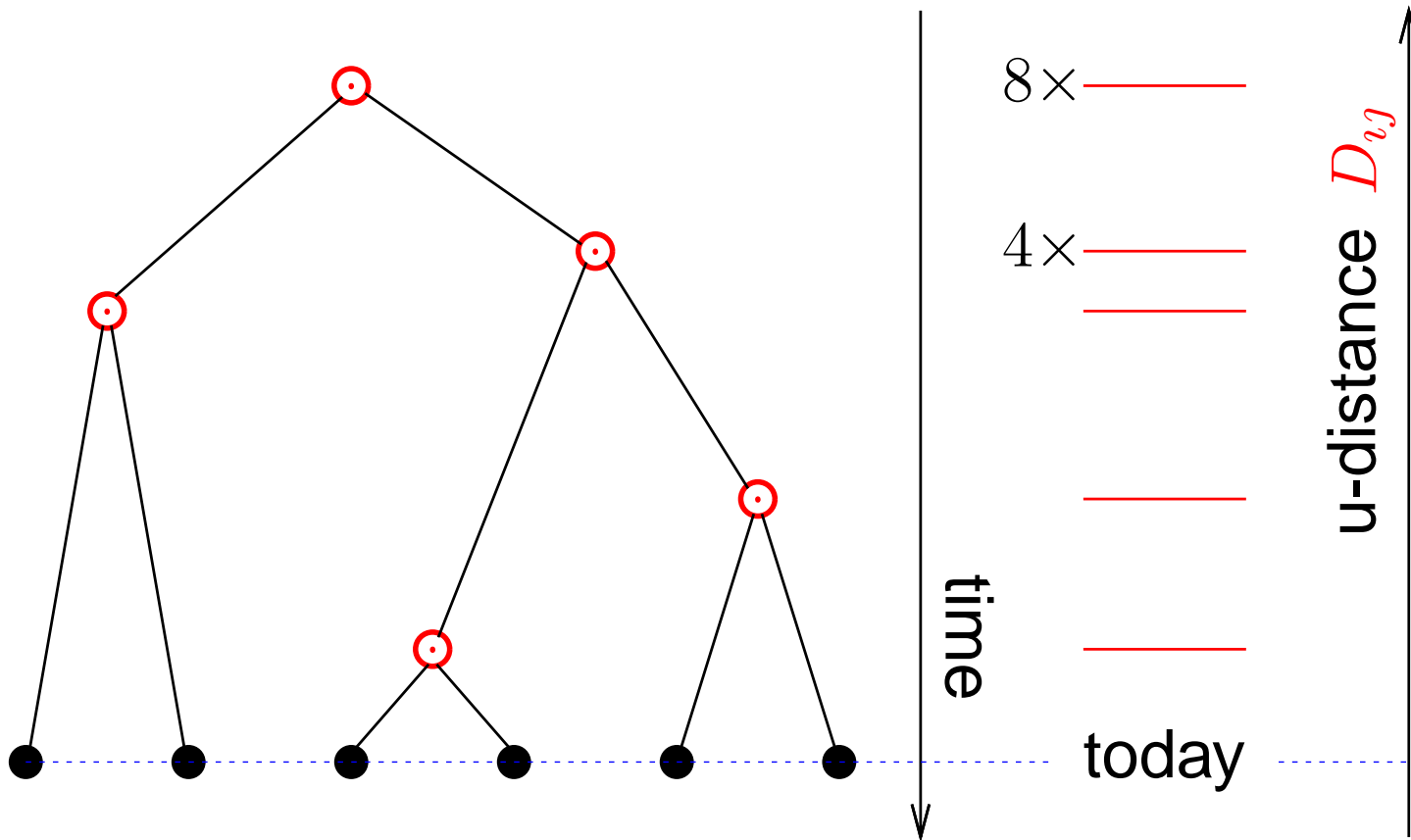
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That is why I called it **Dynamic Ising Model**.

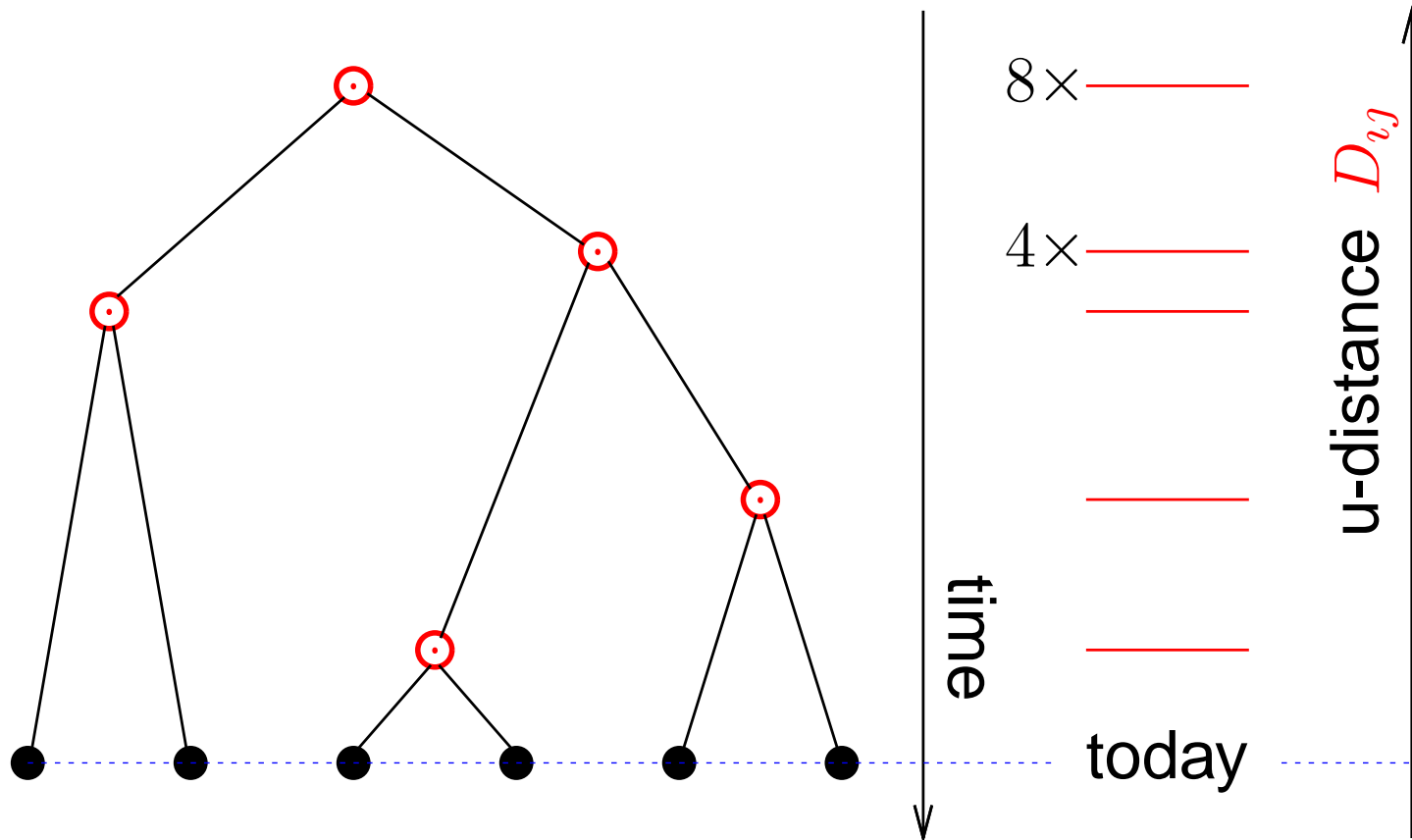
The first evolutionary tree



Another evolutionary tree



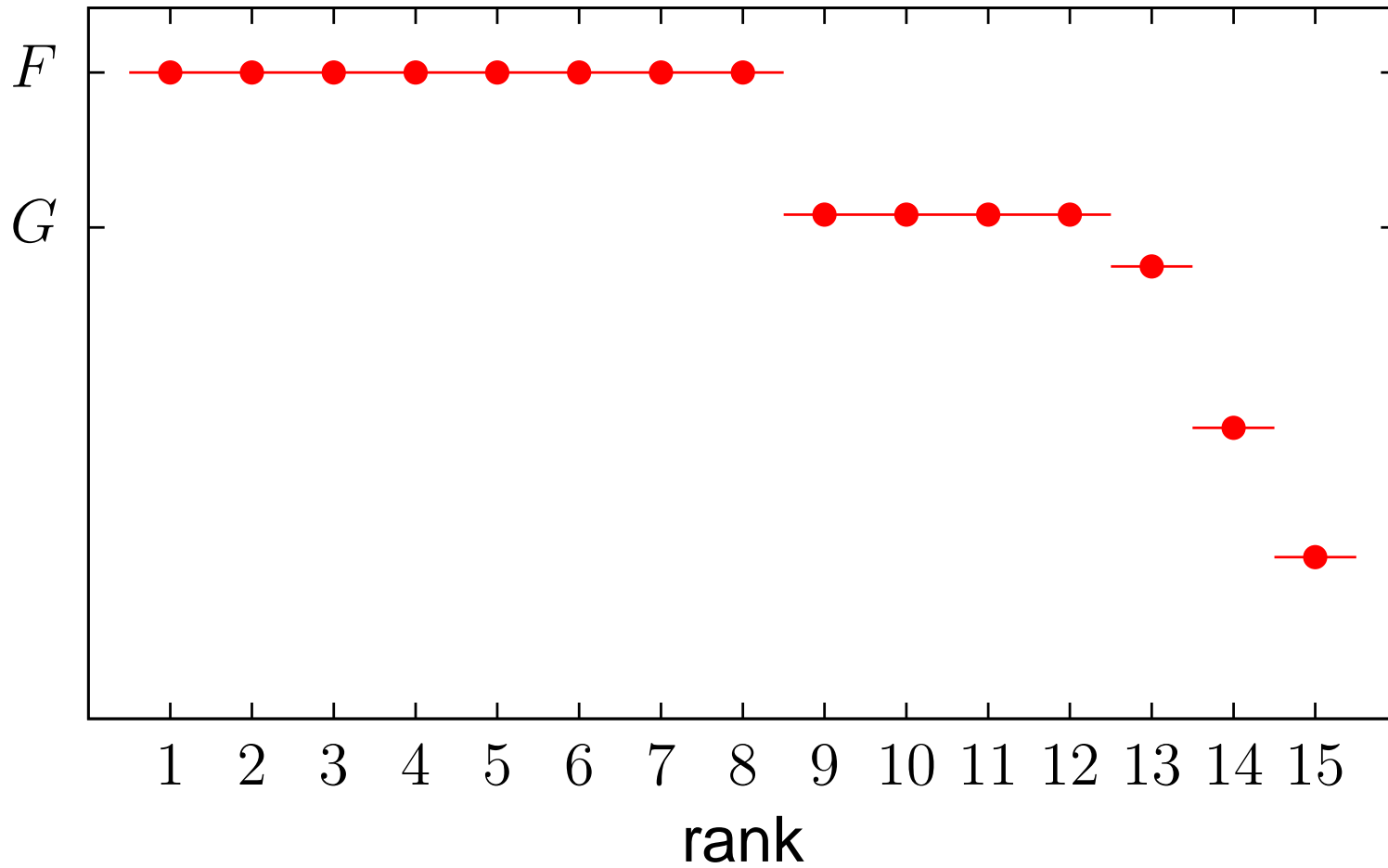
Another evolutionary tree



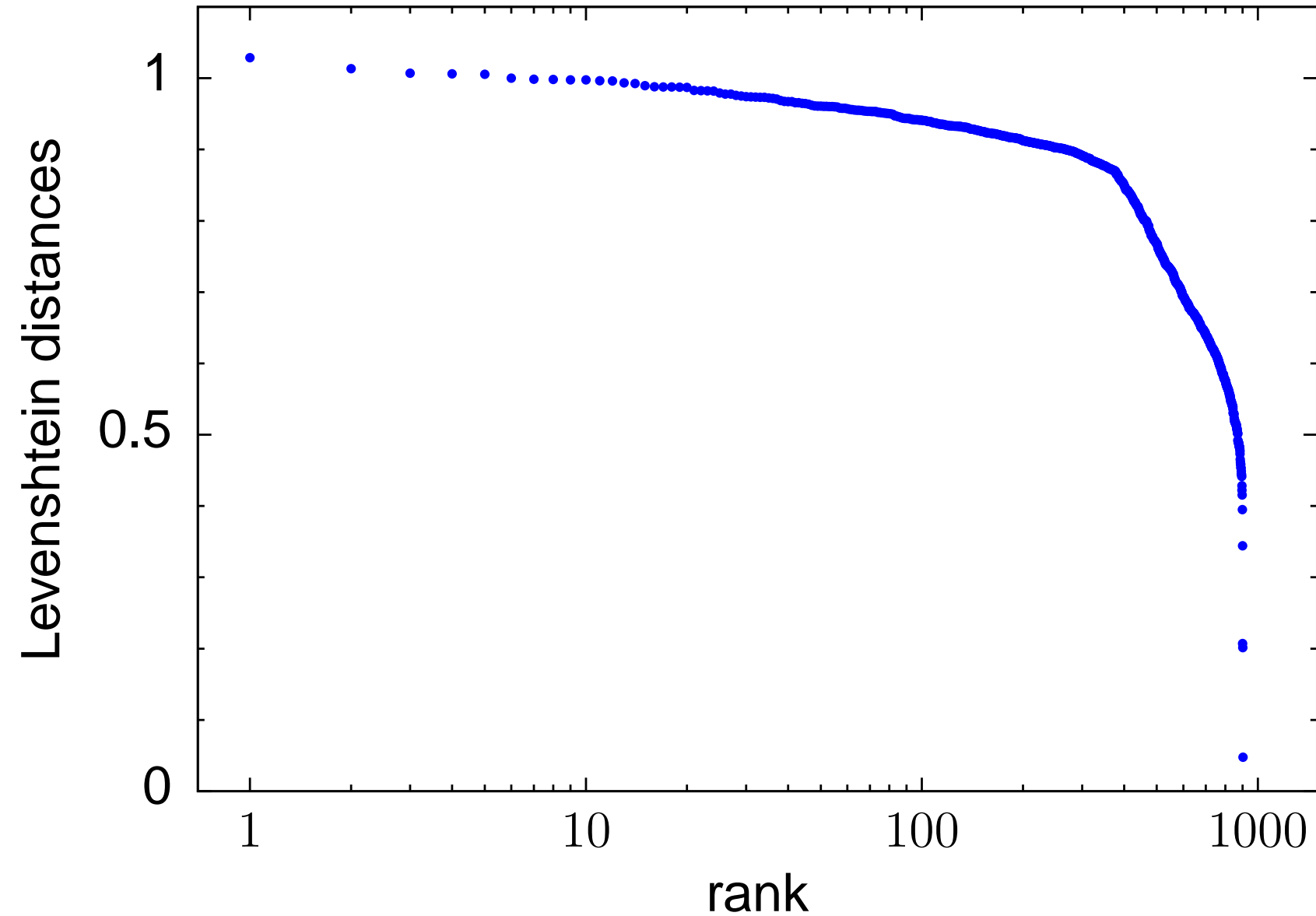
$$J_{ij} = D - D_{ij}$$

Parameter D is adjusted in between the two higher levels.

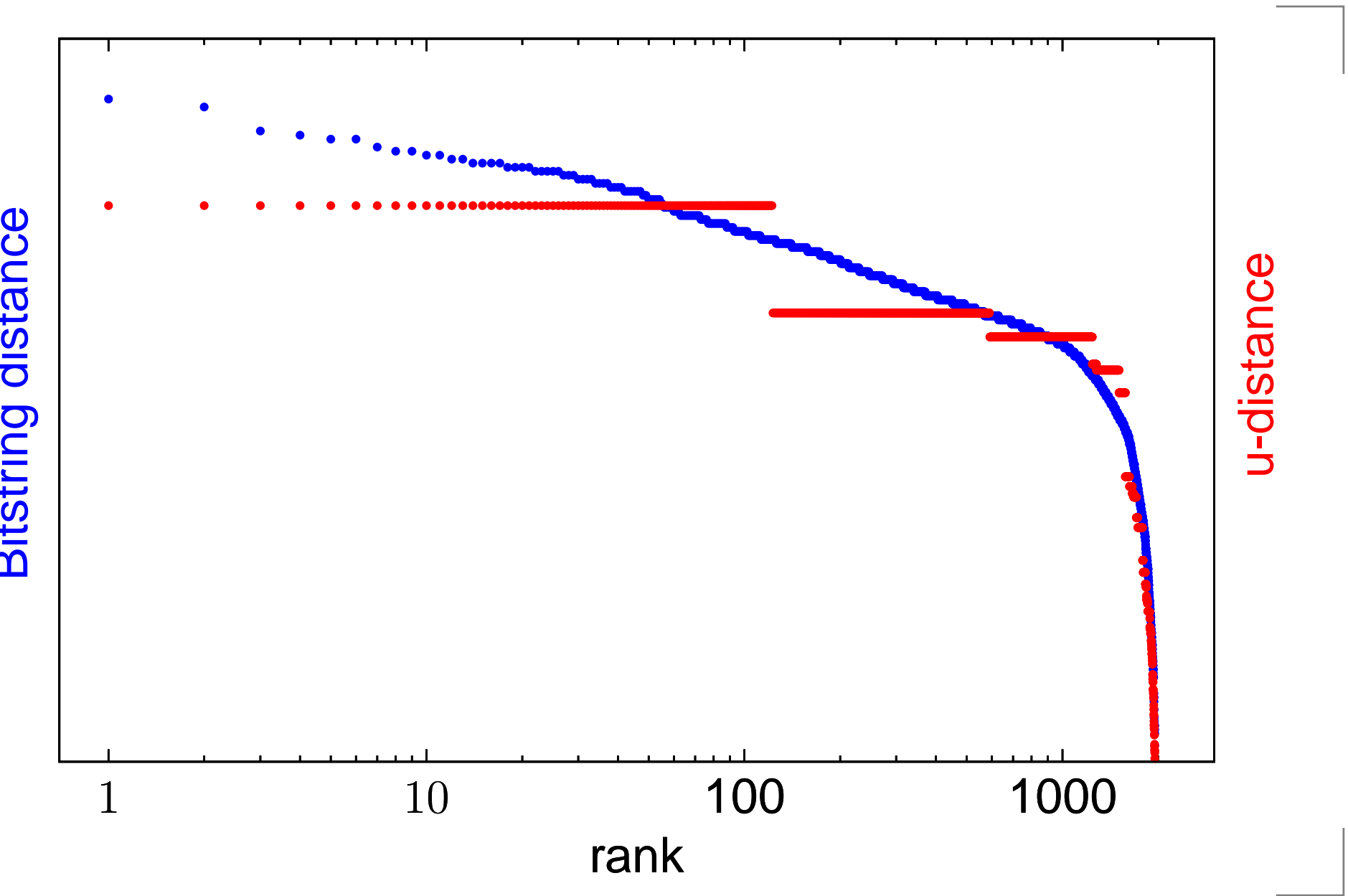
Rank plot, same tree



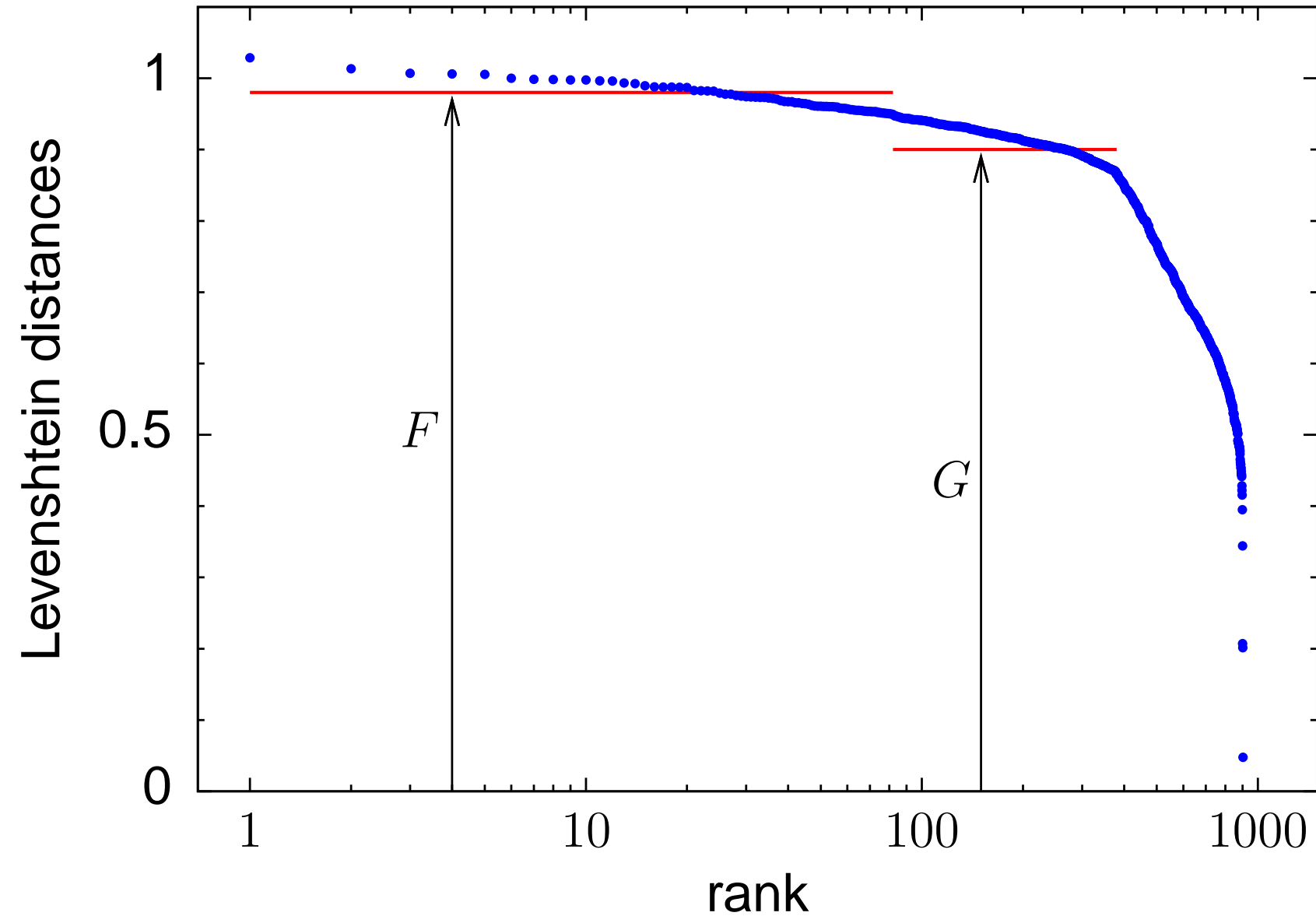
Tupian family, 43 languages



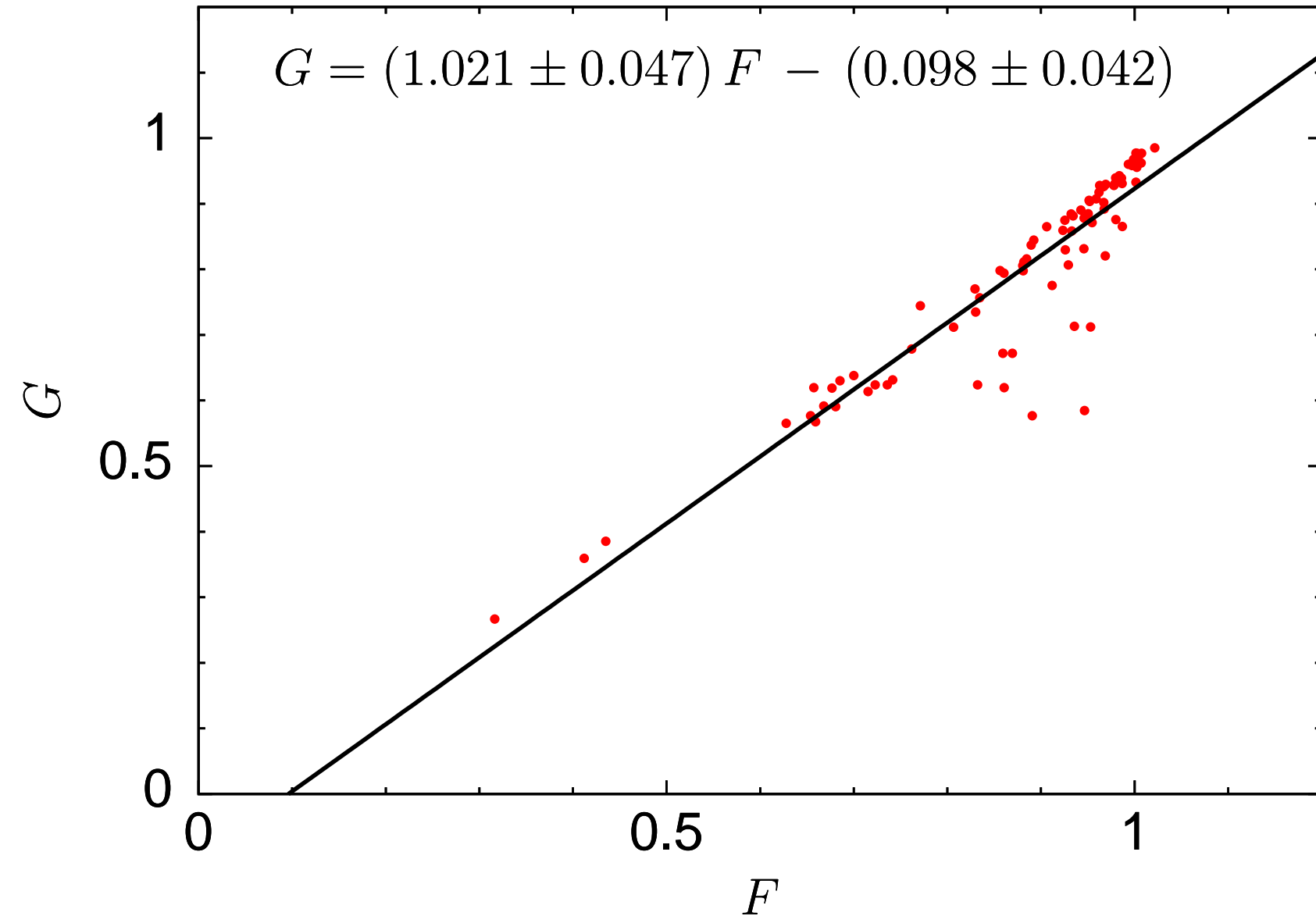
Simulated tree, 63 languages



Measures for each real family



89 real families



Our group at UFF



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Marcio Argollo de Menezes
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Edgardo Brigatti (UFF-VR)
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