

DISTRIBUIÇÕES DE
VELOCIDADES NO
FERROMAGNETO DE
HEISENBERG INERCIAL

FERNANDO D. NOBRE

CBPF - BRASIL

- Fundamental Hypothesis: Ergodicity

→ Time Averages \equiv Ensemble Averages

- Maximum Lyapunov Exponent:

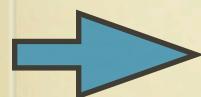
$$\lambda_{\max} = \lim_{t \rightarrow \infty} \frac{1}{t} \ln \left[\frac{d(t)}{d(0)} \right]$$

$\lambda_{\max} > 0 \Rightarrow$ Chaos

- Dynamical Approaches:
(e.g., Numerical Simulations)
- Order of limits $t \rightarrow \infty$, $N \rightarrow \infty$
is irrelevant for ergodic systems
- Non-Ergodic Systems: ??

Good Systems to Investigate:

- Long-Range Interactions
- Should present their own dynamics
- Examples: Systems of Rotors



Present investigation:

Infinite-Range Heisenberg
Rotors

Motivation:

Anomalies in the dynamics of the inertial infinite-range-interaction XY ferromagnet

- Metastable / Quasi-Stationary State (QSS)

$$u \approx u_c :$$

$$t_{QSS} \sim N ; \quad \lambda_{\max} \sim N^{-\kappa} ; \quad (\kappa > 0)$$

- The Model:

$$H = K + V$$

$$H = \frac{1}{2} \sum_{i=1}^N \sum_{\mu=1}^n L_{i\mu}^2 + \frac{1}{2N} \sum_{i,j=1}^N (1 - \vec{S}_i \cdot \vec{S}_j)$$

- Normalization: $\sum_{\mu=1}^n S_{i\mu}^2 = 1 \quad (\forall i)$

- n=2: XY rotors
- n=3: Heisenberg rotors

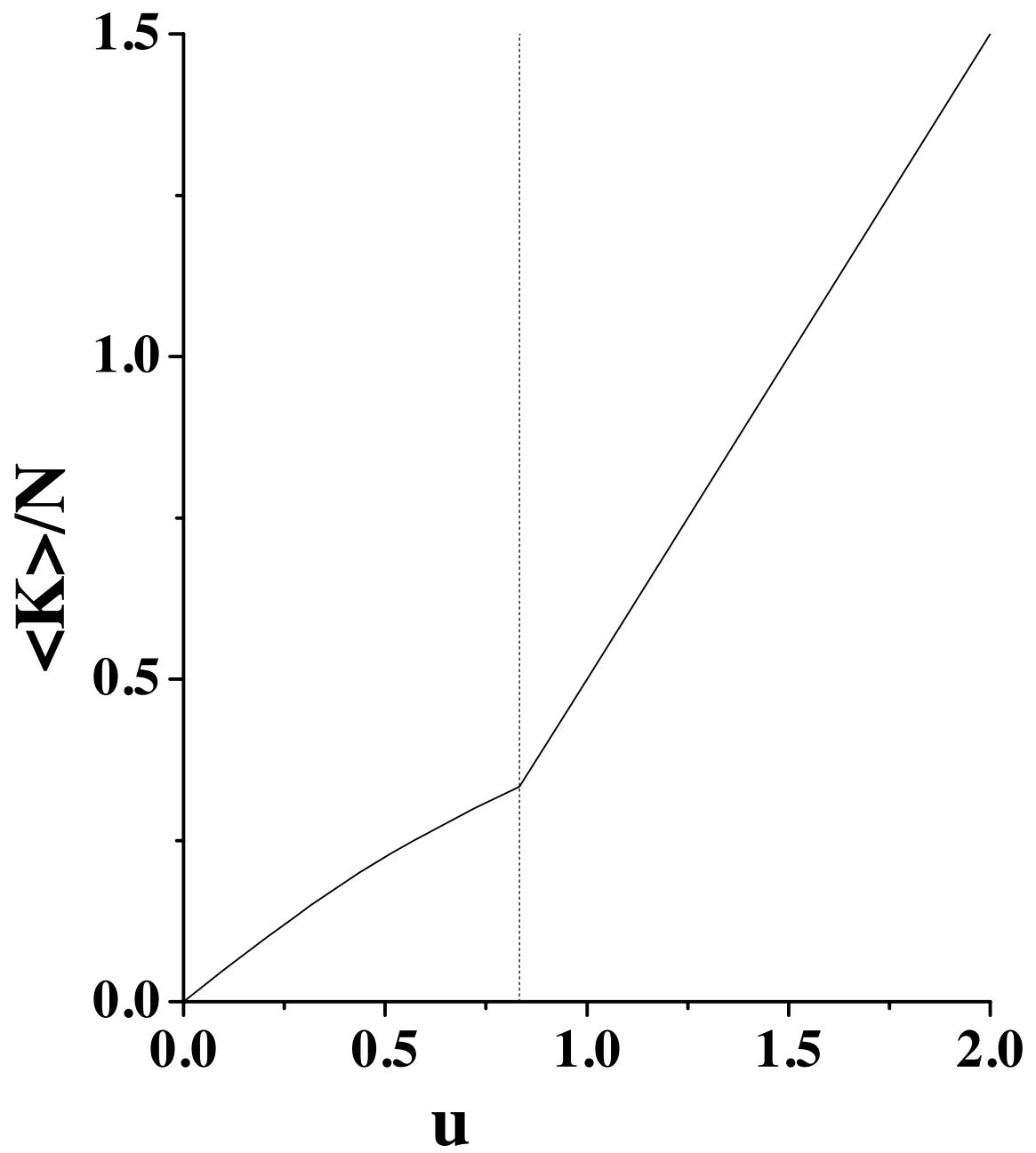
● Canonical Ensemble Solution:

$$u = \frac{n-1}{2\beta} + \frac{1}{2}(1 - \vec{m}^2)$$

$$m \equiv |\vec{m}| = \frac{I_{n/2}(\beta m)}{I_{(n-2)/2}(\beta m)}$$

$$T_c = \frac{1}{n} ; \quad u_c = 1 - \frac{1}{2n} \quad (k_B = 1)$$

$$n = 3 \rightarrow u_c = \frac{5}{6}$$



- Dynamics of the Inertial Heisenberg

$$\dot{\vec{L}}_i = \vec{S}_i \times \left(\frac{1}{N} \sum_{j=1}^N \vec{S}_j \right) \quad (i = 1, 2, \dots, N)$$

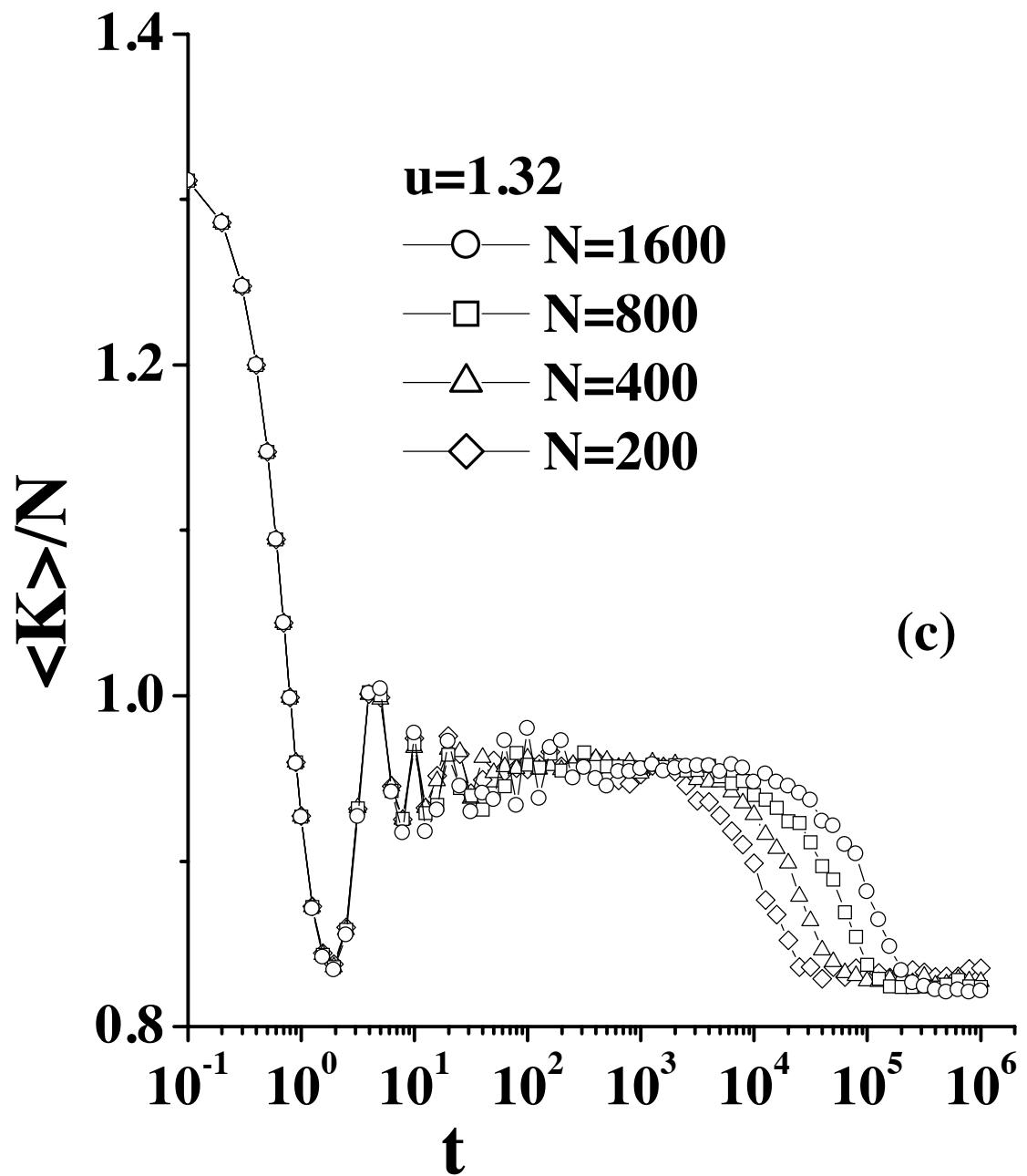
$$\dot{\vec{S}}_i = \vec{L}_i \times \vec{S}_i$$

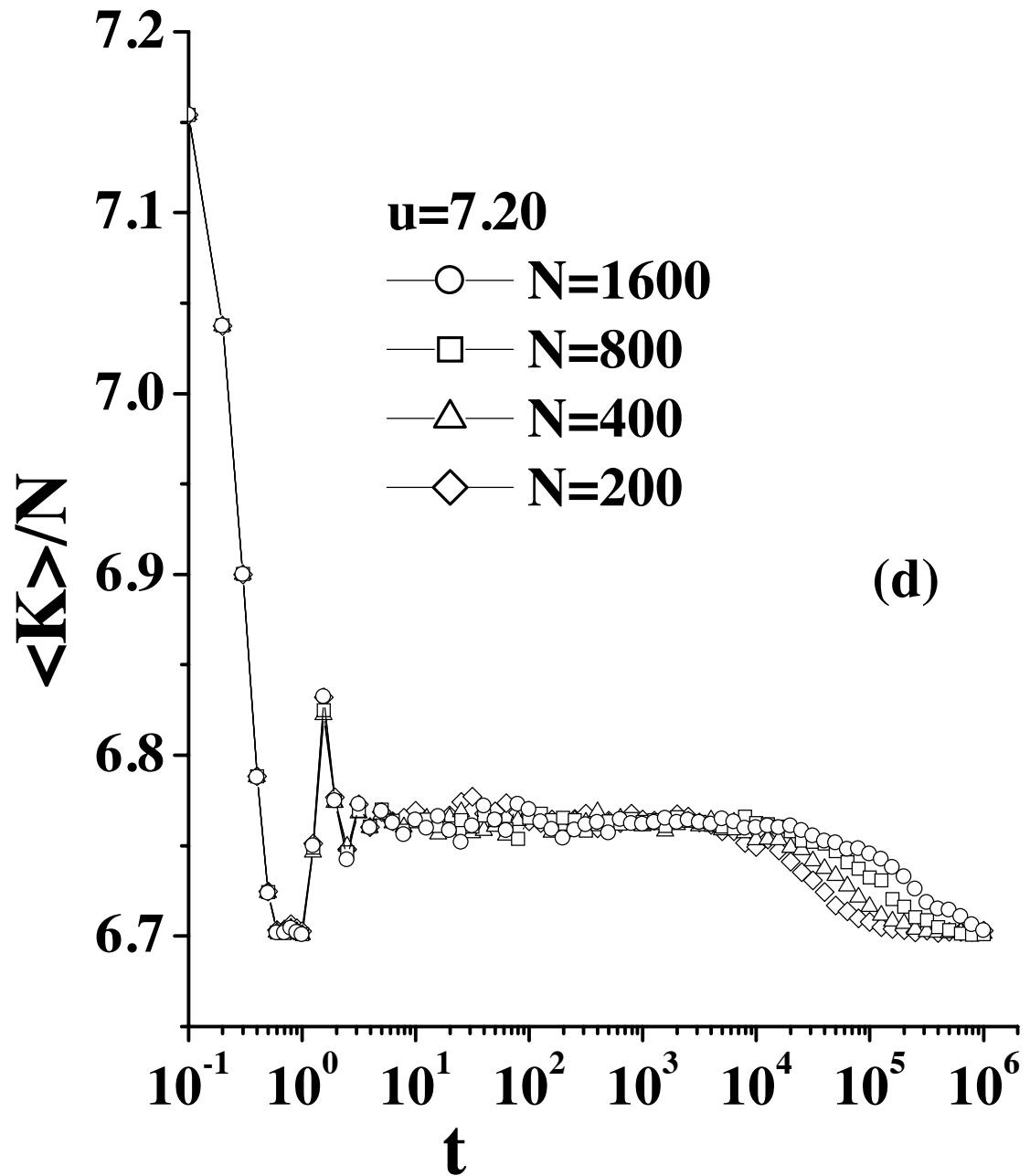
→ 6N Equations

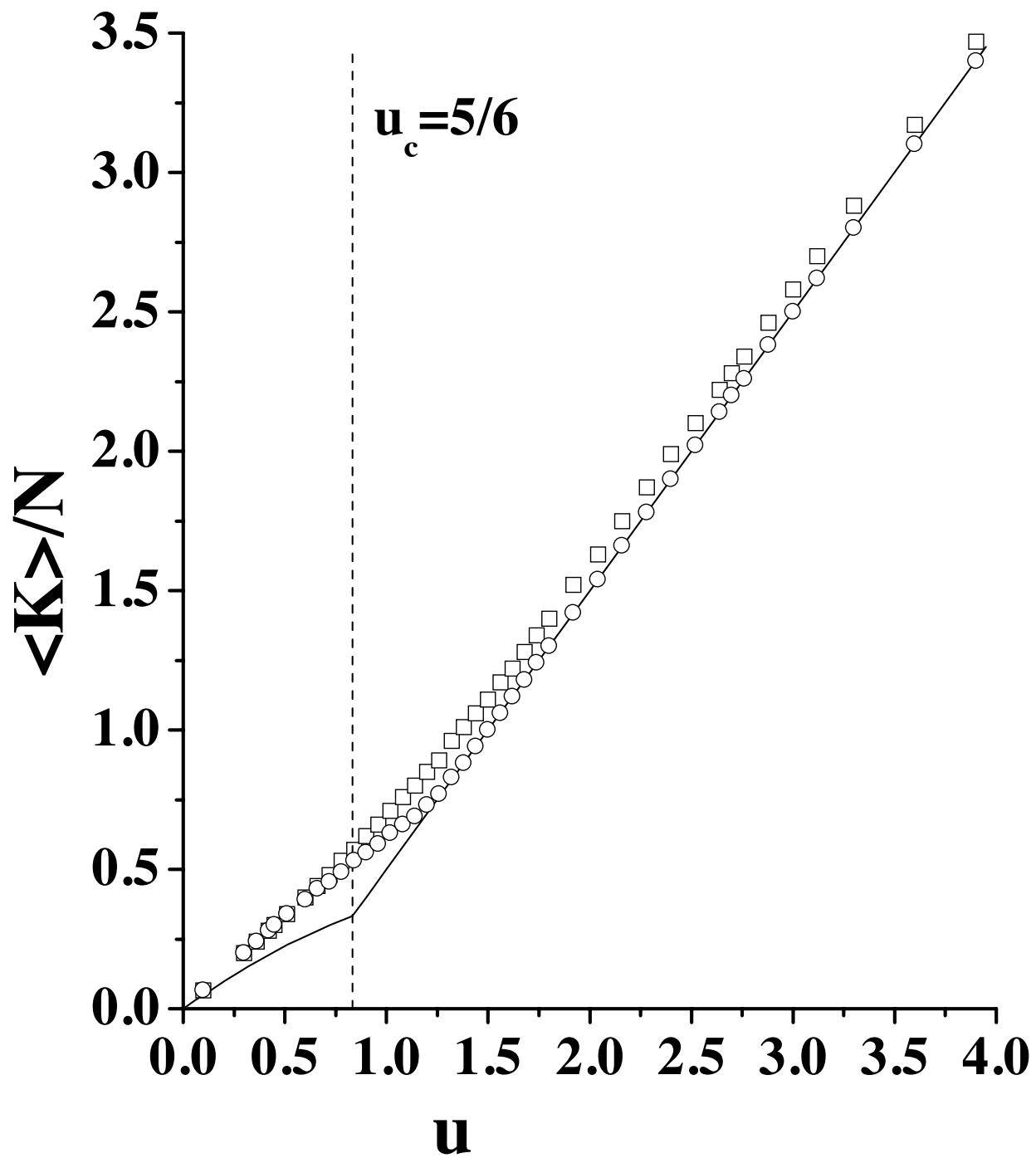
- 4th order Runge-Kutta-Merson Integrator
- Preserve: Total Energy and Spin Norm.

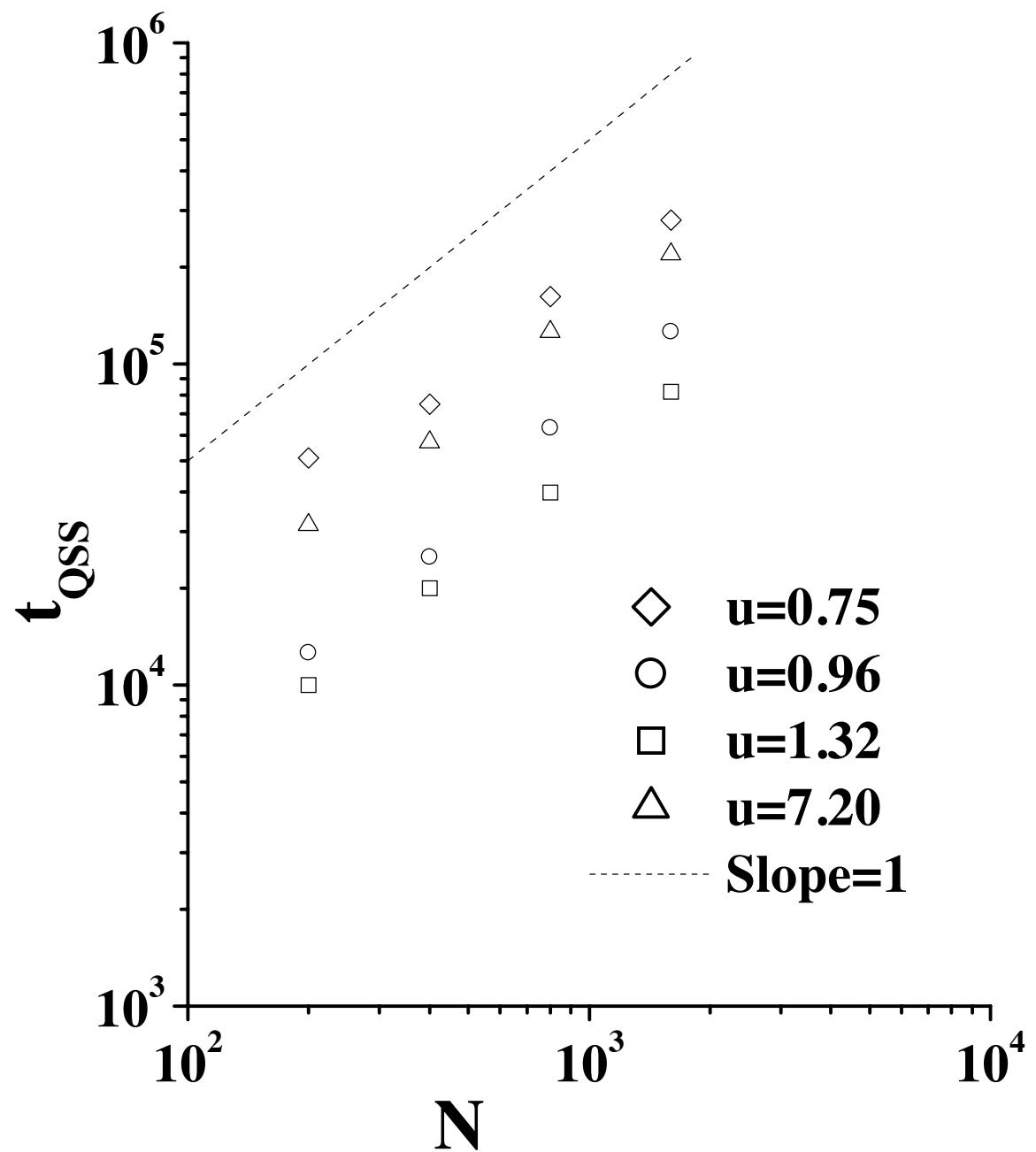
F. D. Nobre and C. Tsallis, Phys. Rev. E 68,
036115 (2003).

- Initial Magnetization: $m_z(0) = 1$
 - Initial Angular Momenta: Equal “Water Bags” for the Components $(\mu = x, y, z)$
- Similar MS for wider range of energies





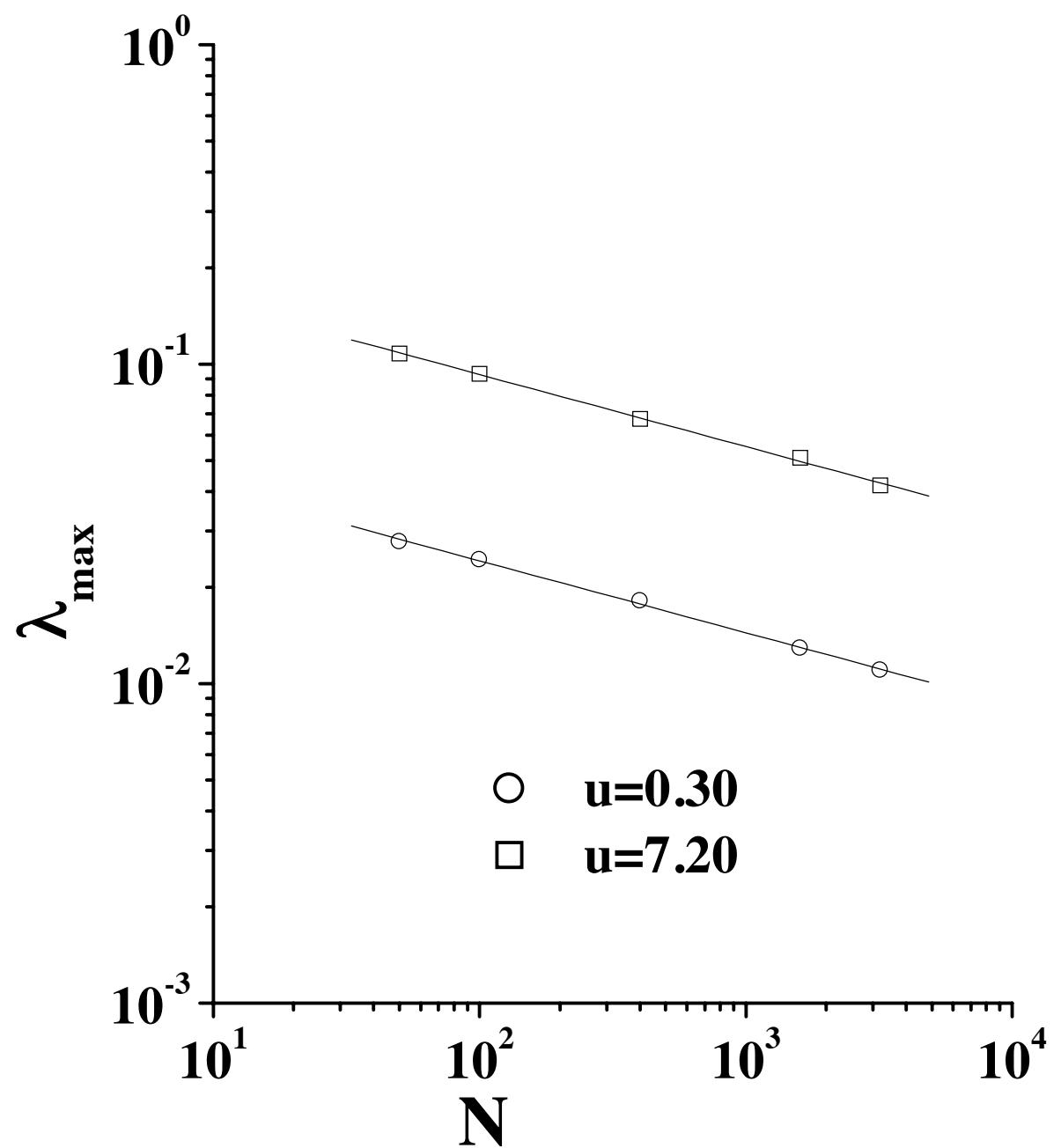




● Maximum Lyapunov Exponent

$$\lambda_{\max} = \lim_{t \rightarrow \infty} \frac{1}{t} \ln \left[\frac{d(t)}{d(0)} \right]$$

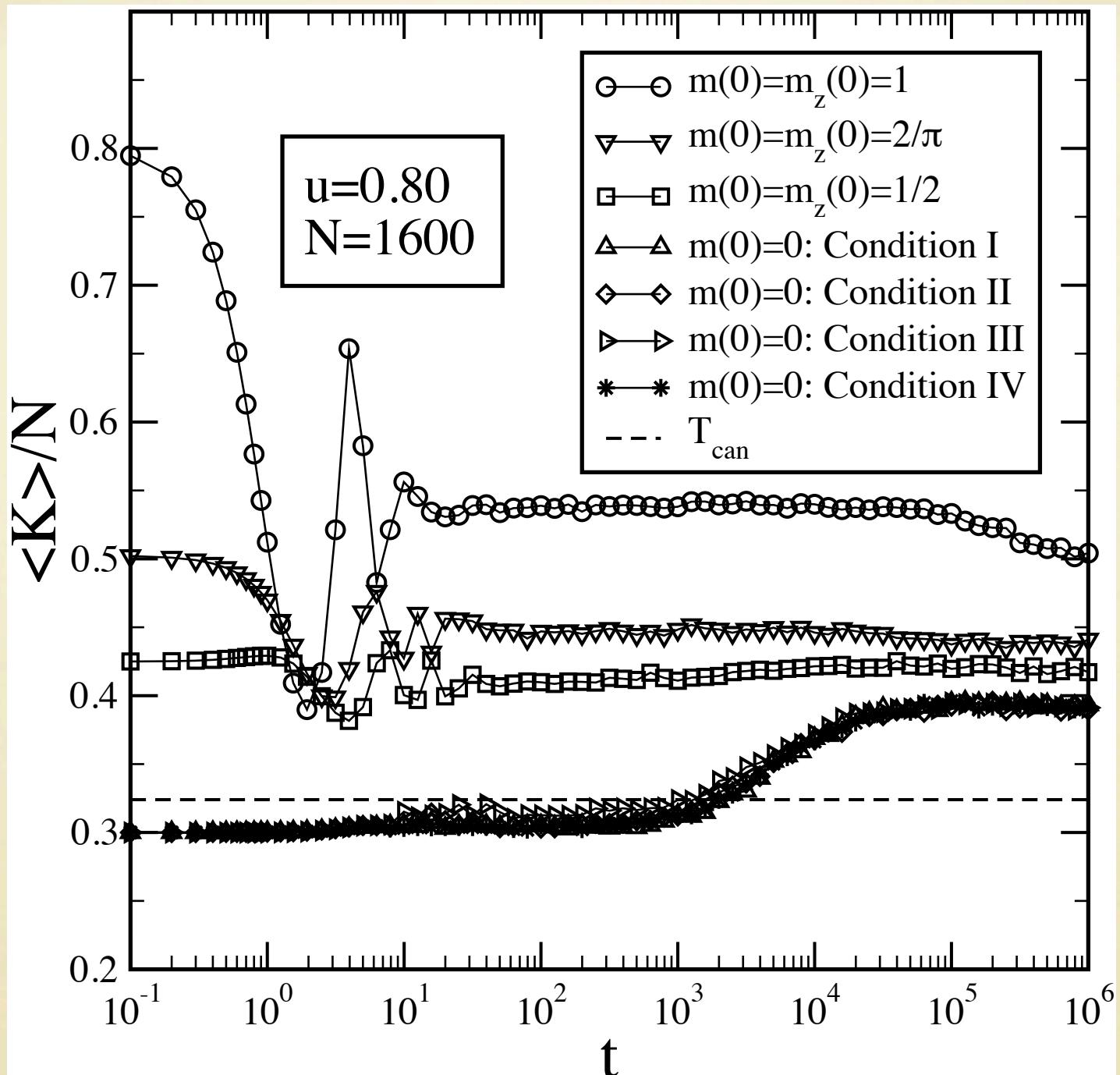
$$d(t) = \left\{ \sum_{i=1}^N \sum_{\mu=x,y,z} [(\delta L_{i\mu})^2 + (\delta S_{i\mu})^2] \right\}^{1/2}$$



F. D. Nobre and C. Tsallis, Physica A 344,
587 (2004).

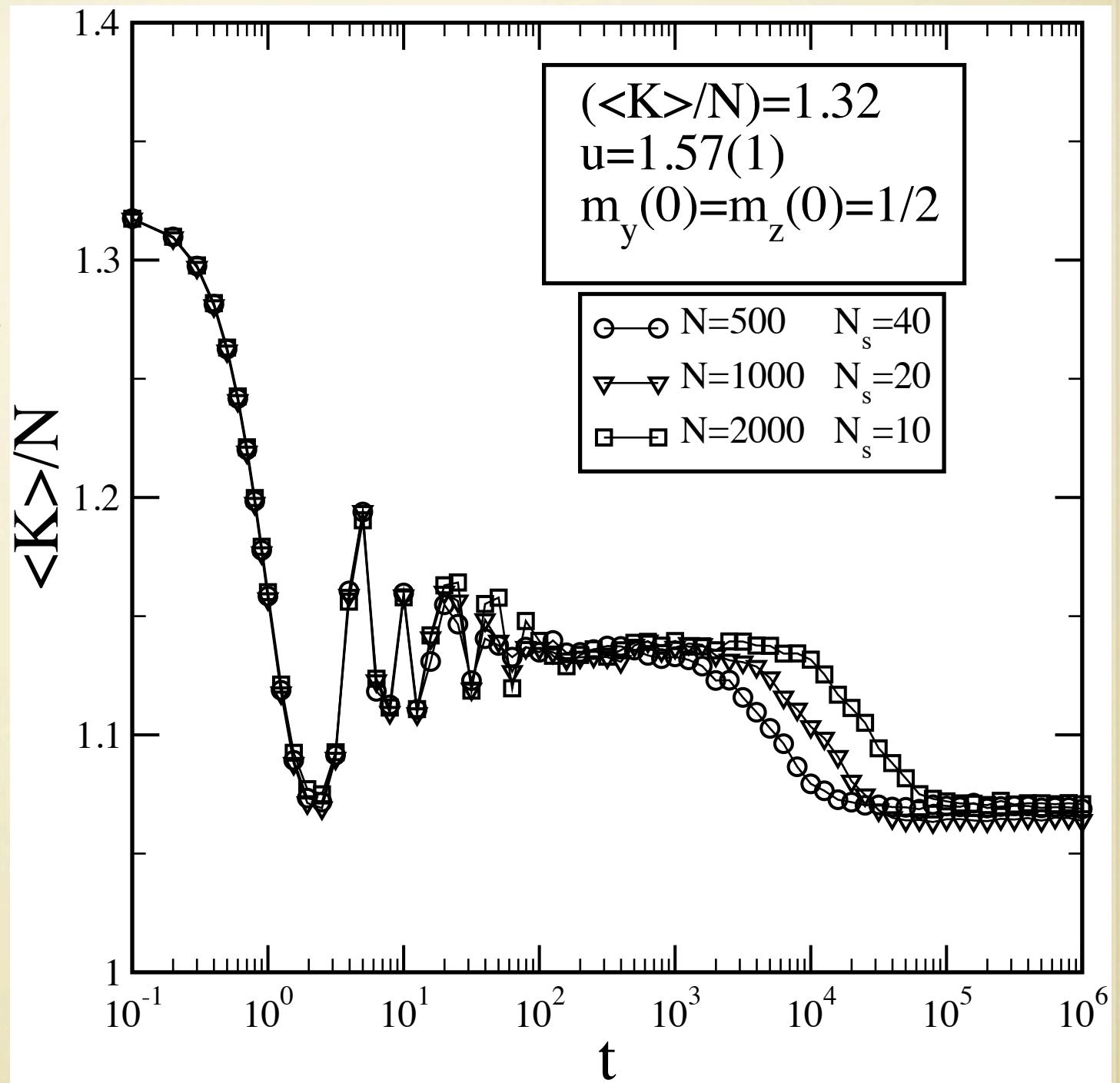
- Role of Initial Conditions on MS:
 - i) Different Initial Magnetizations; Equal “Water Bags” for the Angular-Momenta Components

→ Distinct MS



$S_y(0) = 1, \quad p = 1/2$

$S_z(0) = 1, \quad p = 1/2$



- PDFs of angular velocities:

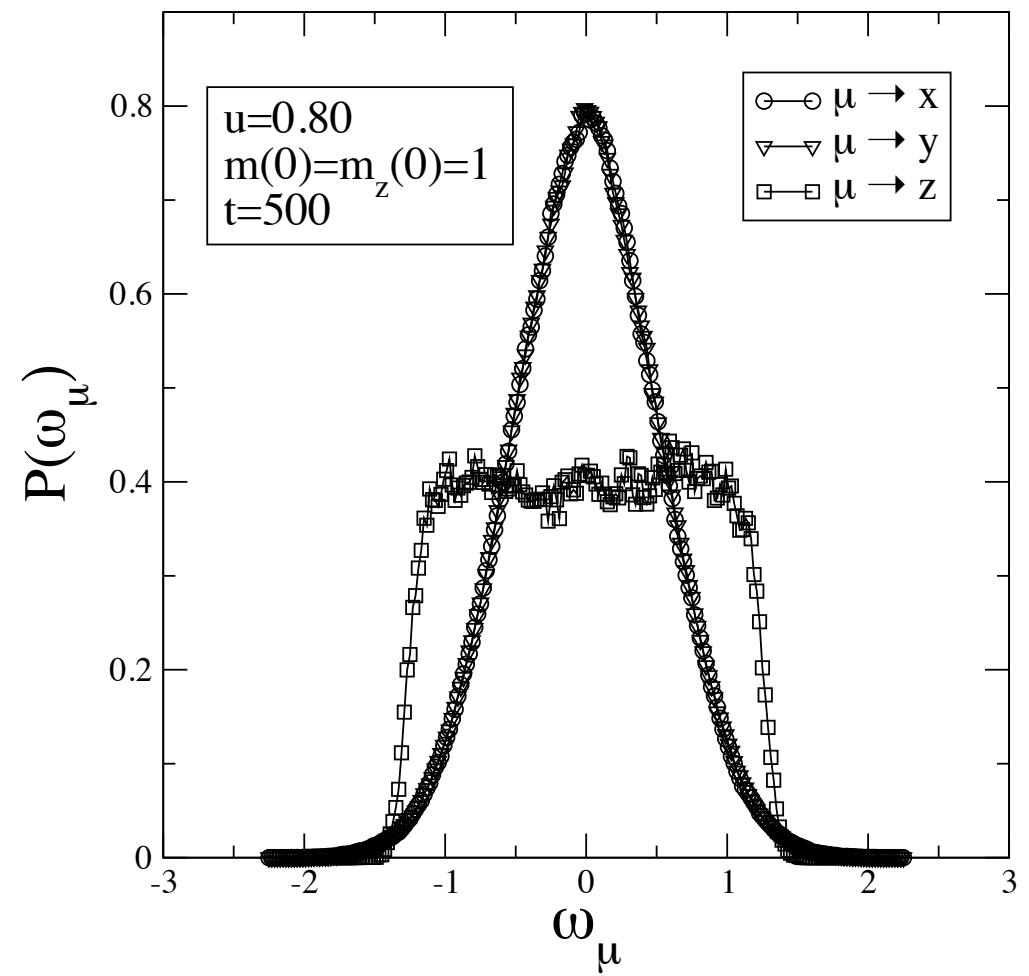
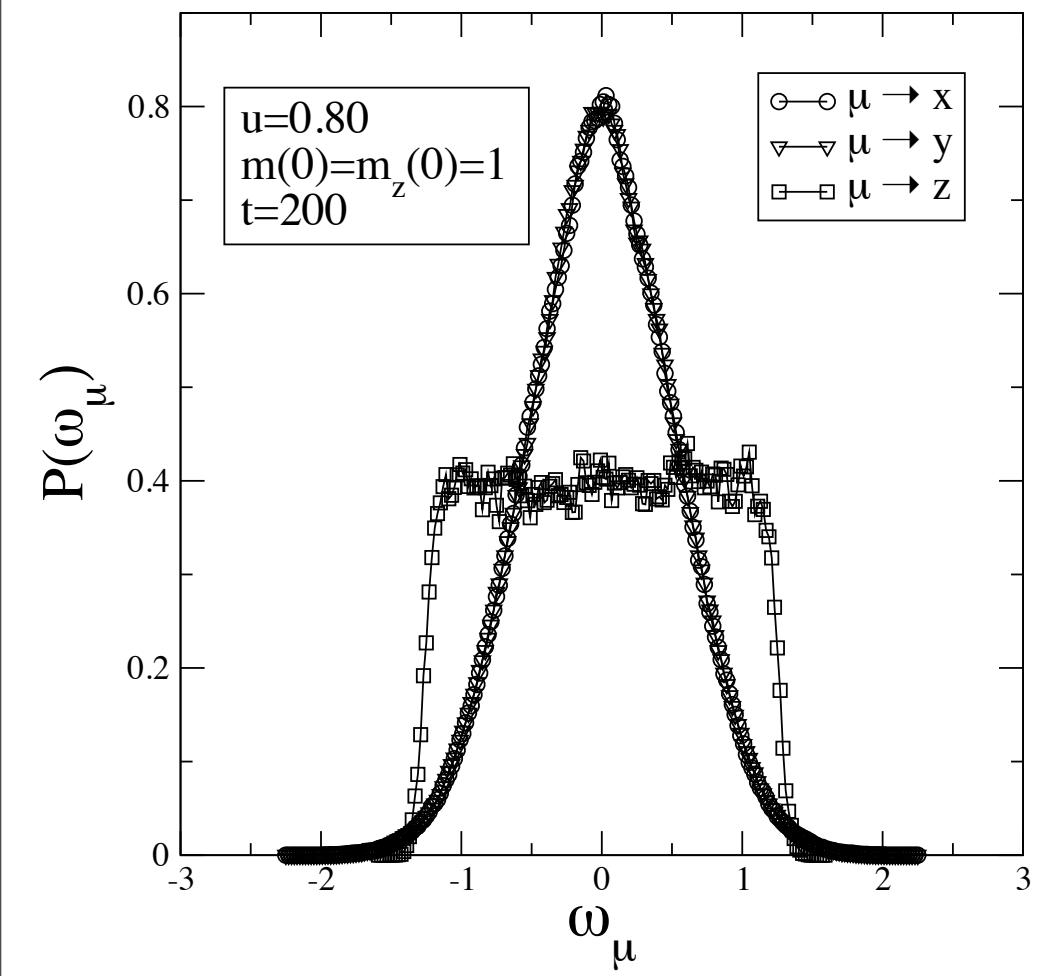
- i) Ensemble averages: $P(\omega_\mu)$

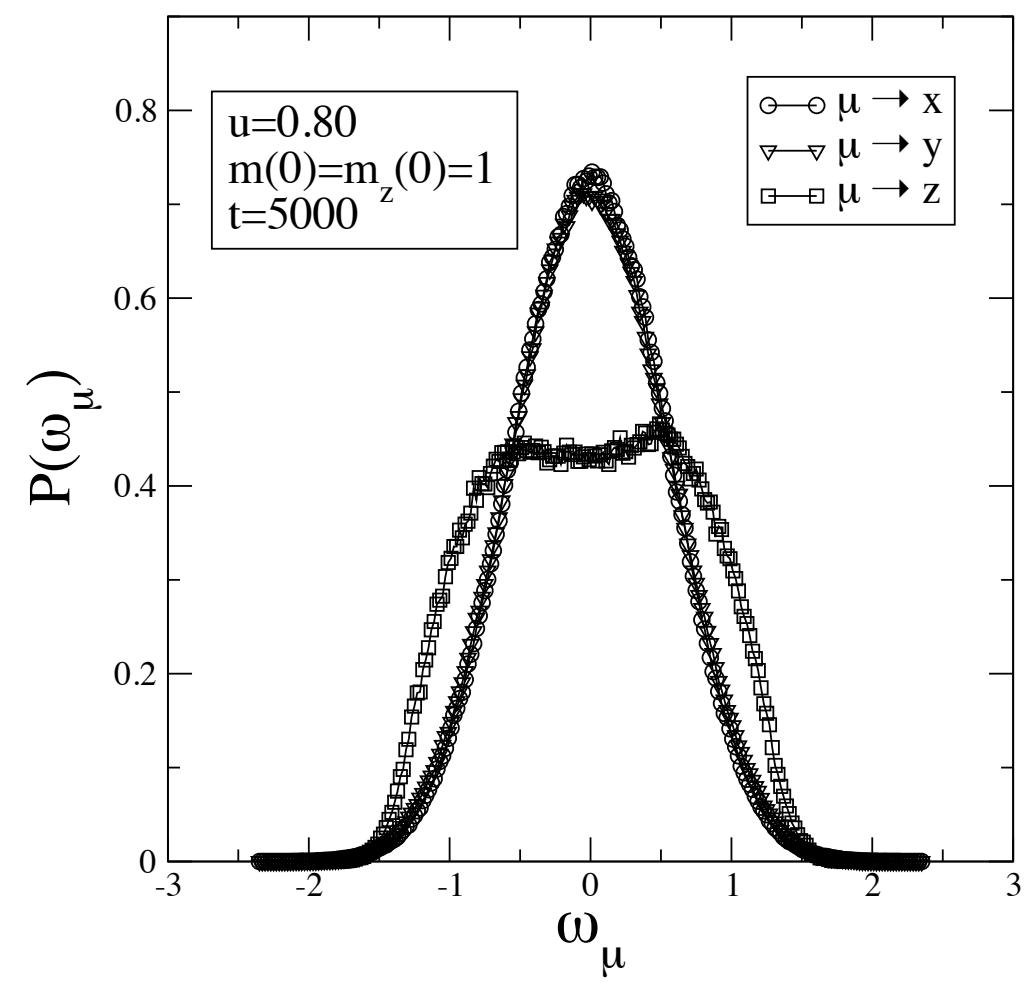
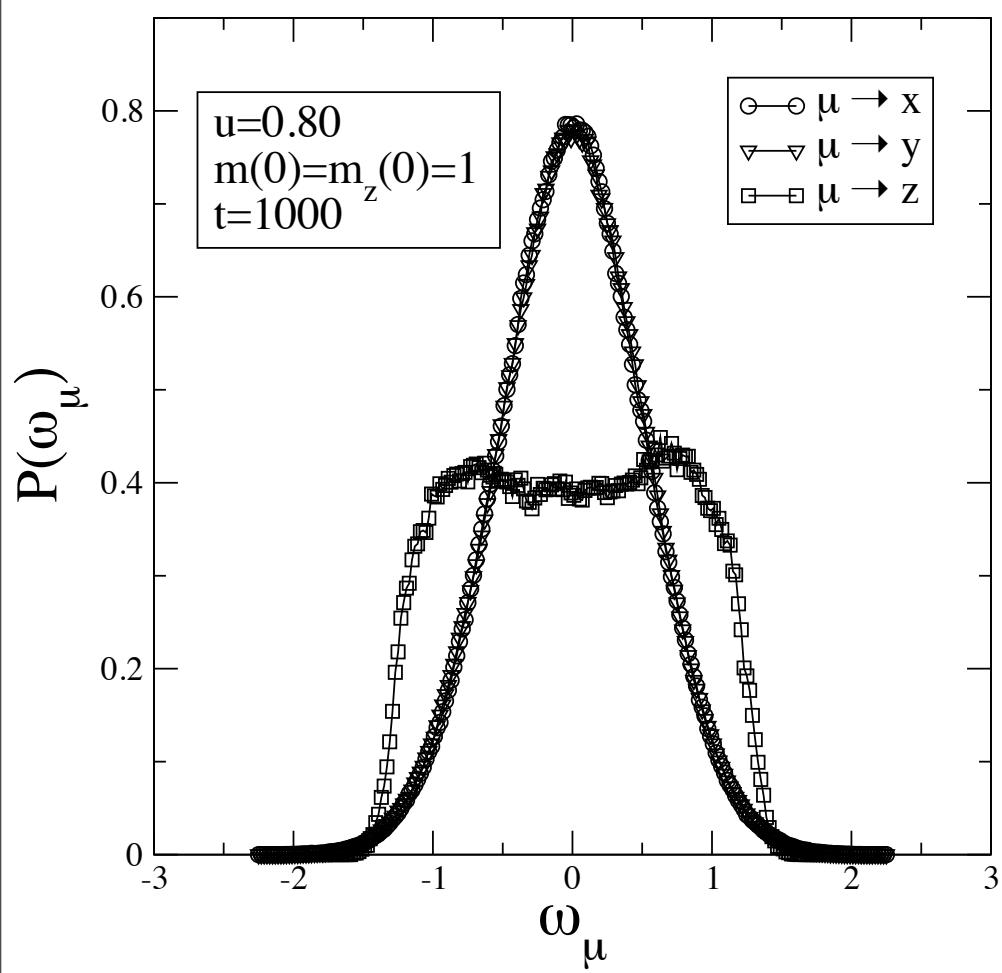
- ii) Time av. over single trajectories: $P(u_\mu)$

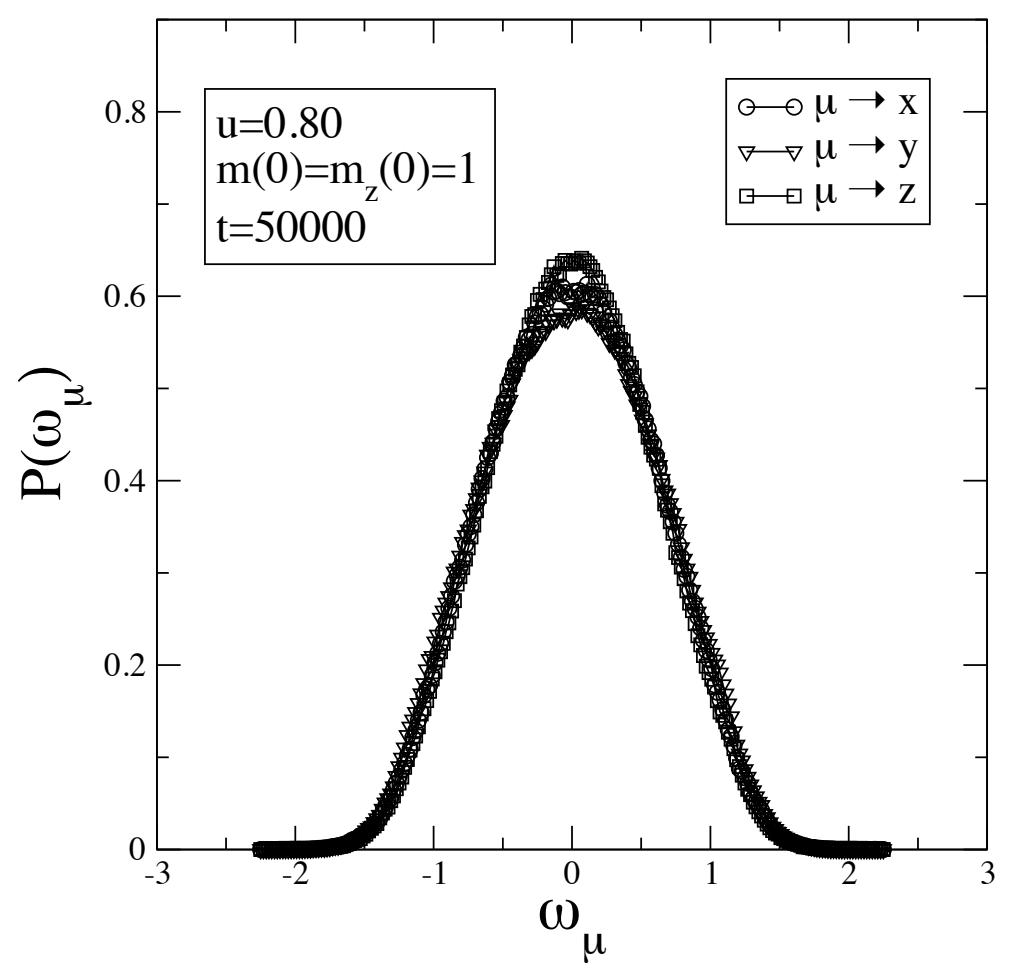
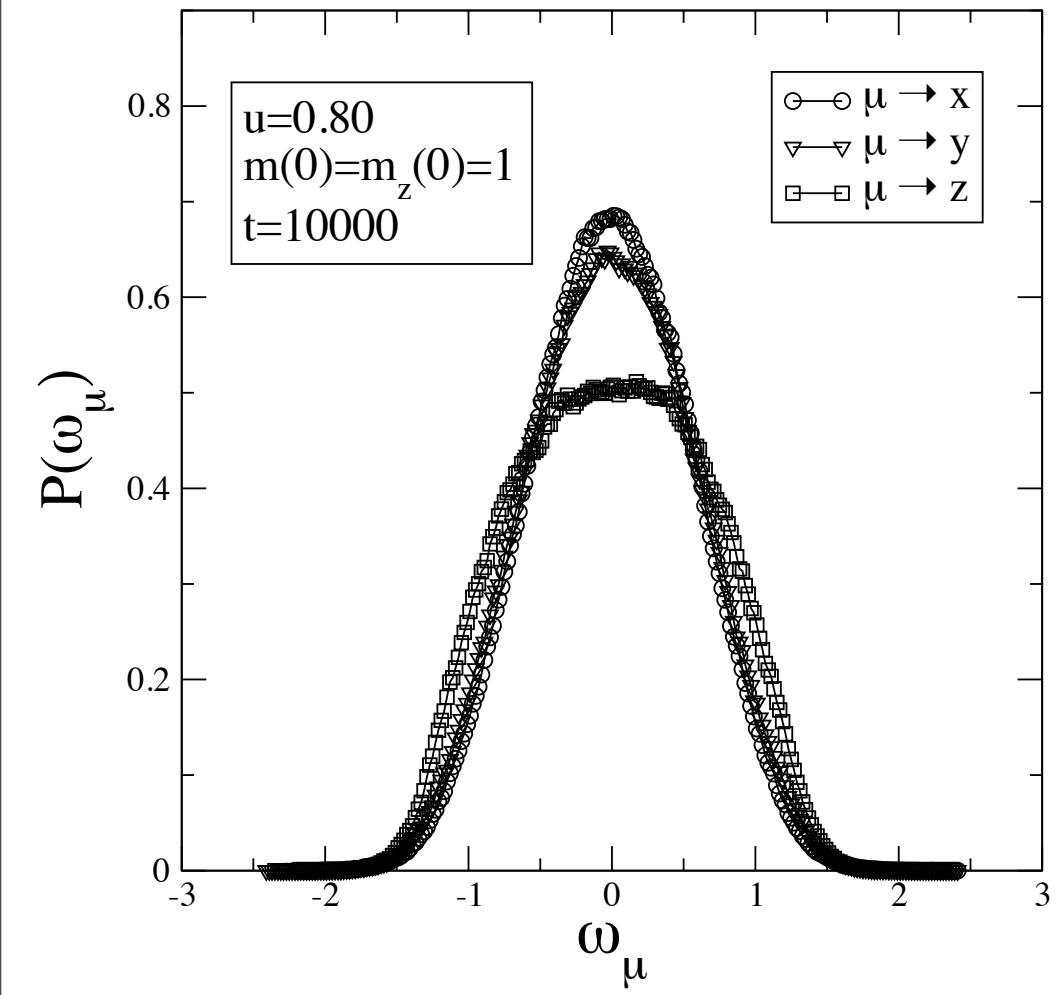
$$u_{i\mu} = \frac{1}{\sqrt{n}} \sum_{k=1}^n \omega_{i\mu}(k) \quad (i = 1, \dots, N; \mu = x, y, z)$$

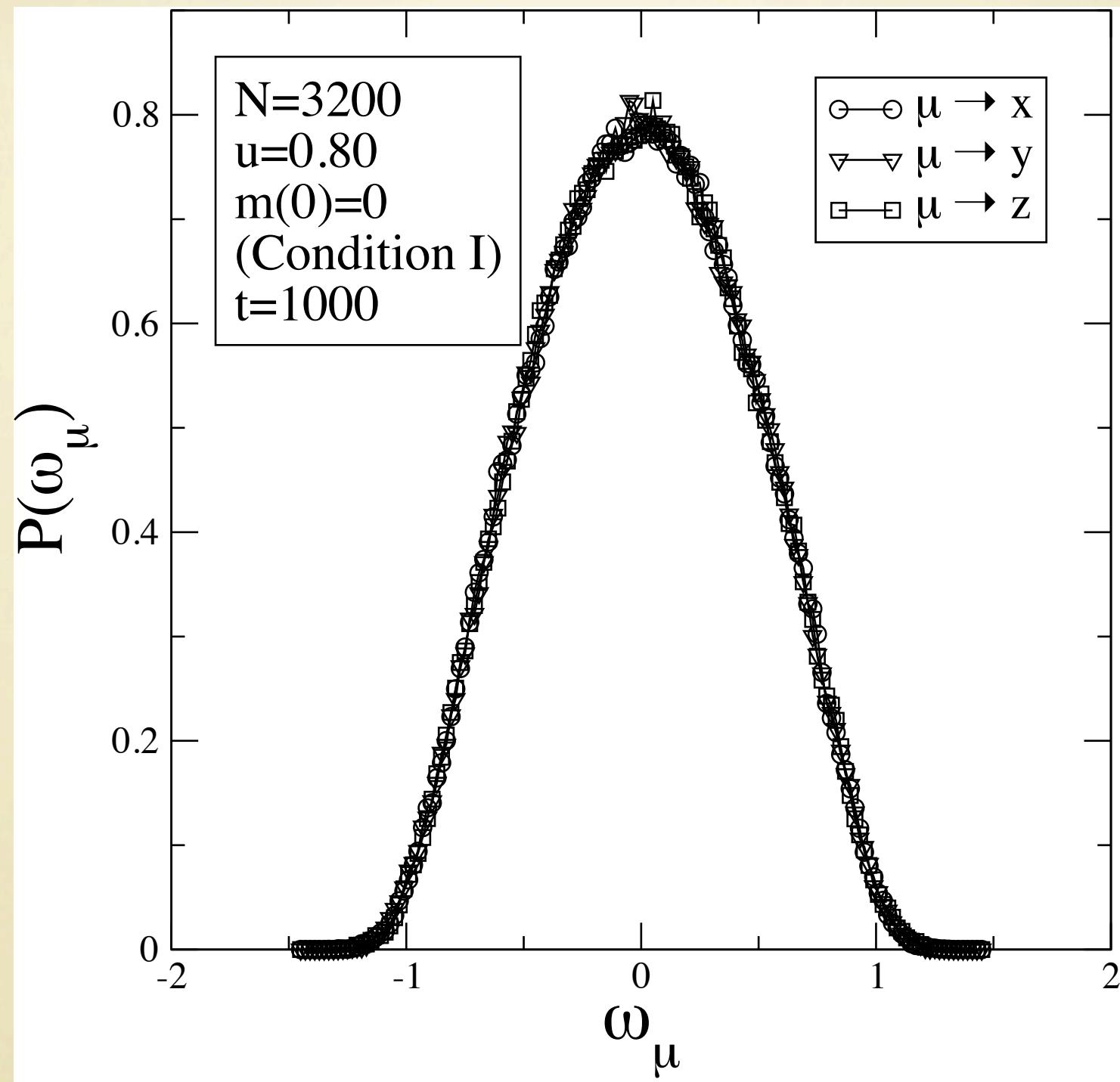
n : no. measures taken at fixed time interval δ

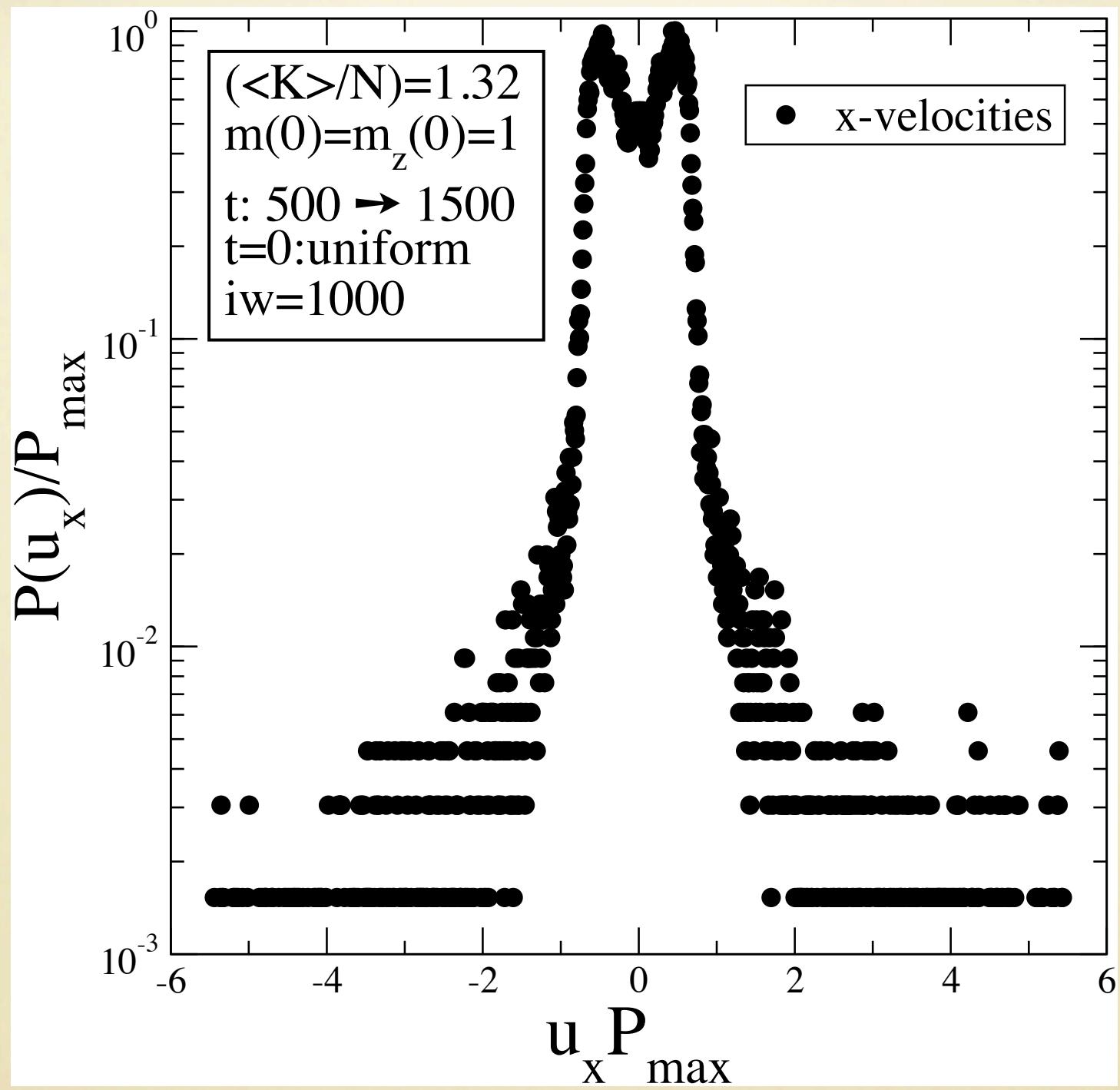
$n\delta$: total time for measurements

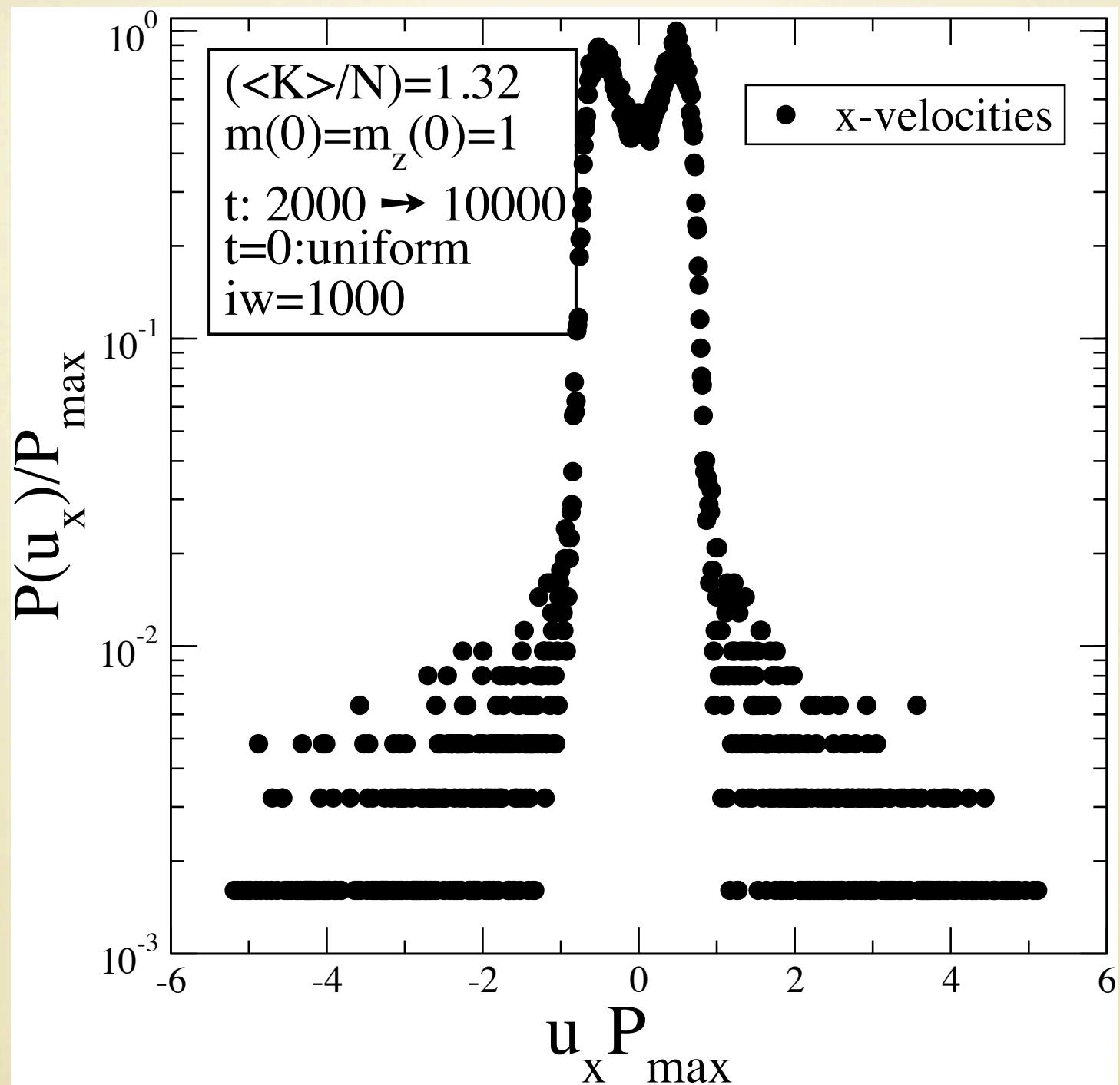


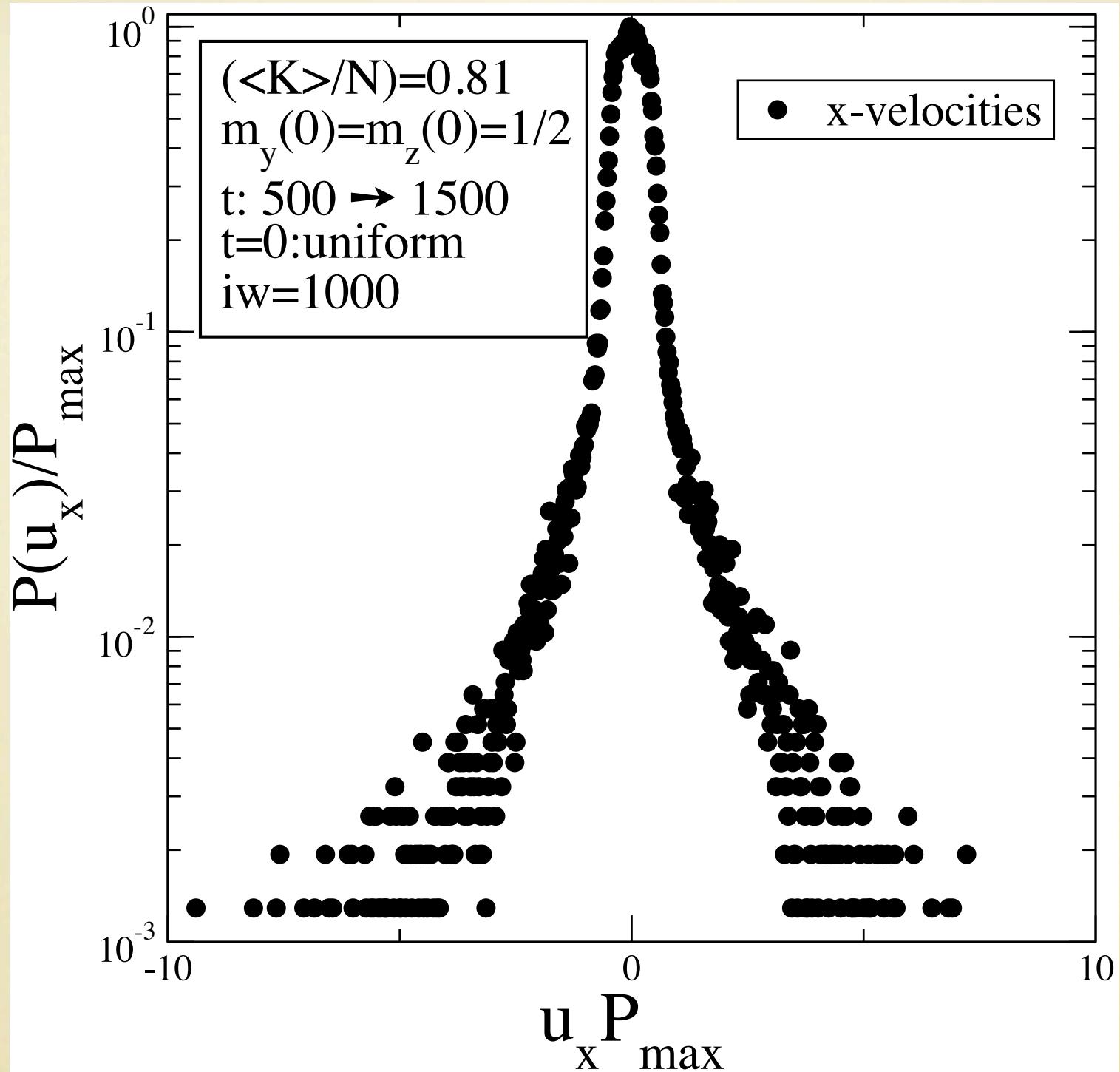


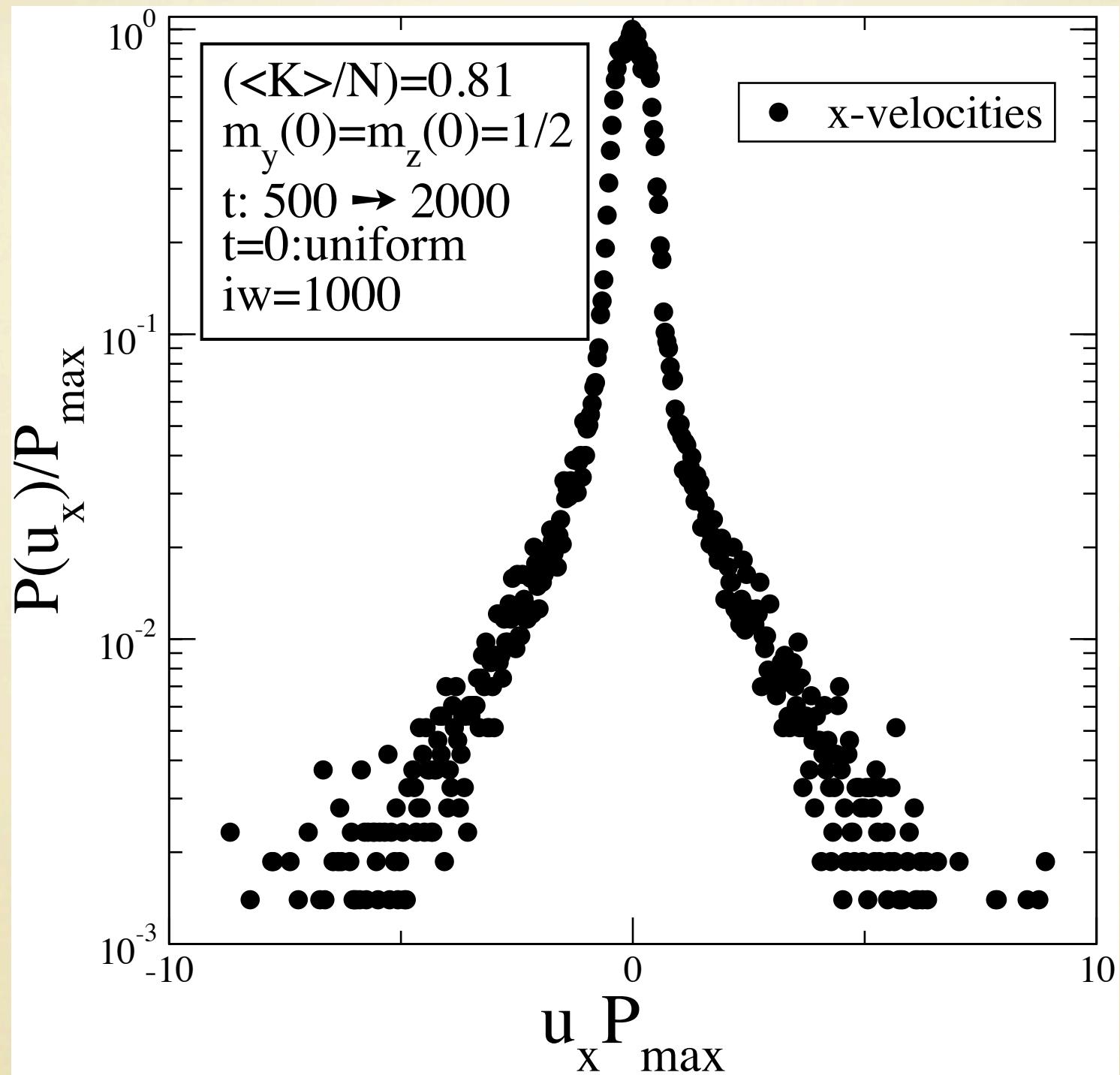


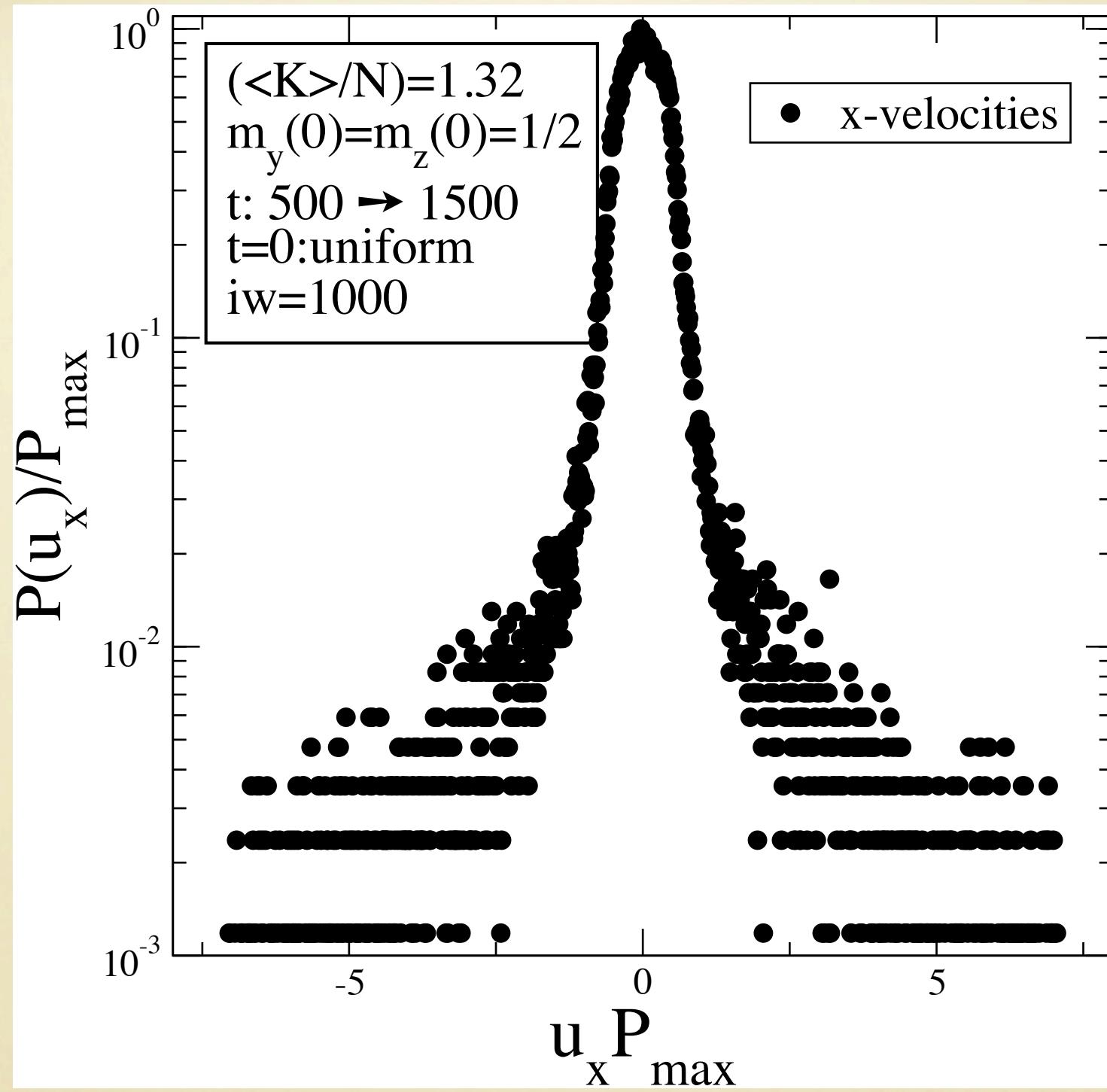


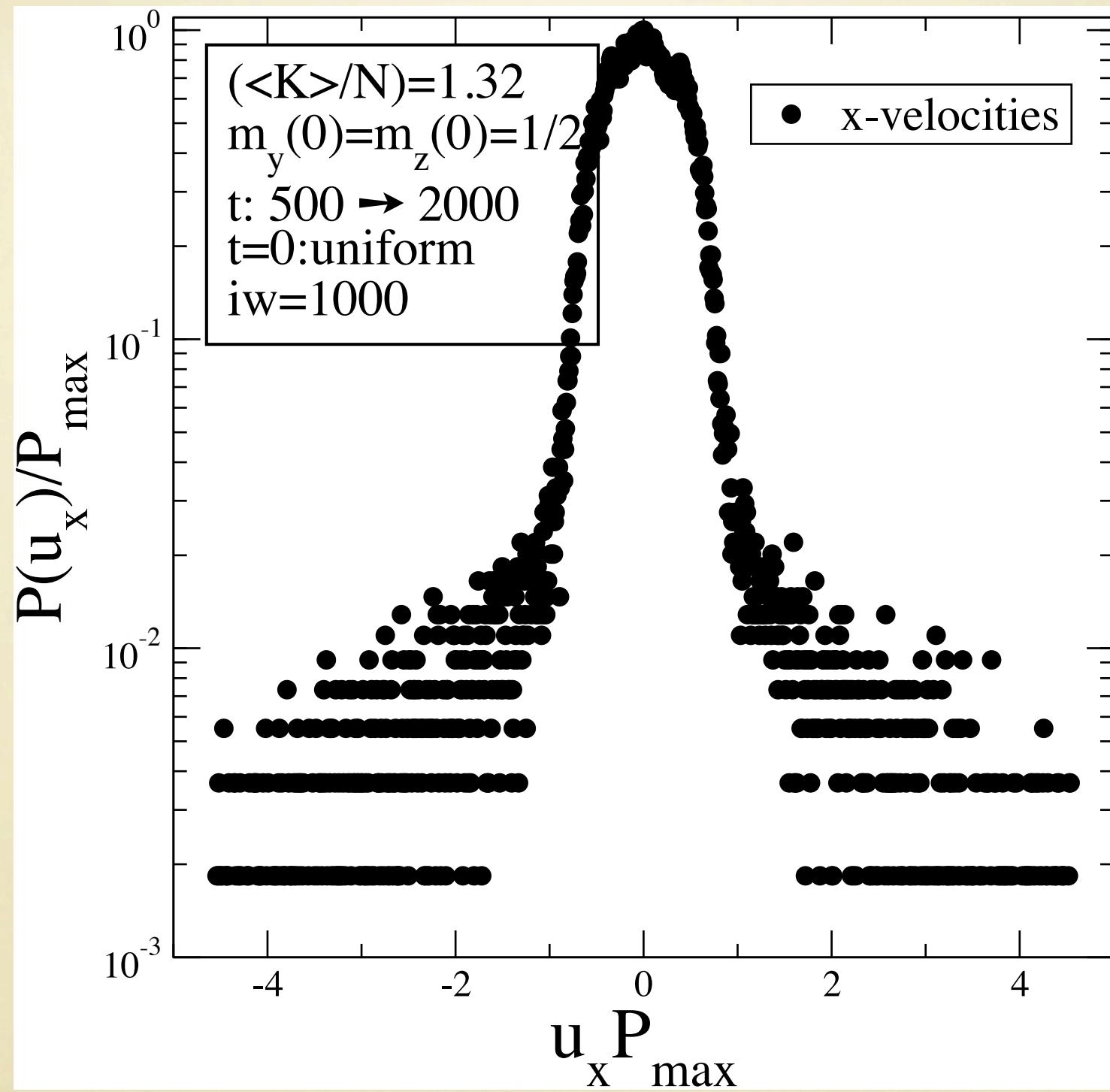












Conclusions

- Model is Non-Ergodic
 - Considering $N \rightarrow \infty$ and $t \rightarrow \infty$
- System remains on QSS forever
- CLT does NOT apply on QSS
 - Standard Statistical Mechanics does NOT hold on QSS

- Homogeneous Function (degree 1):

$$S(\lambda V, \lambda U, \lambda N) = \lambda S(V, U, N)$$

→ Extensivity

- Holds Only for Short-Range Interactions

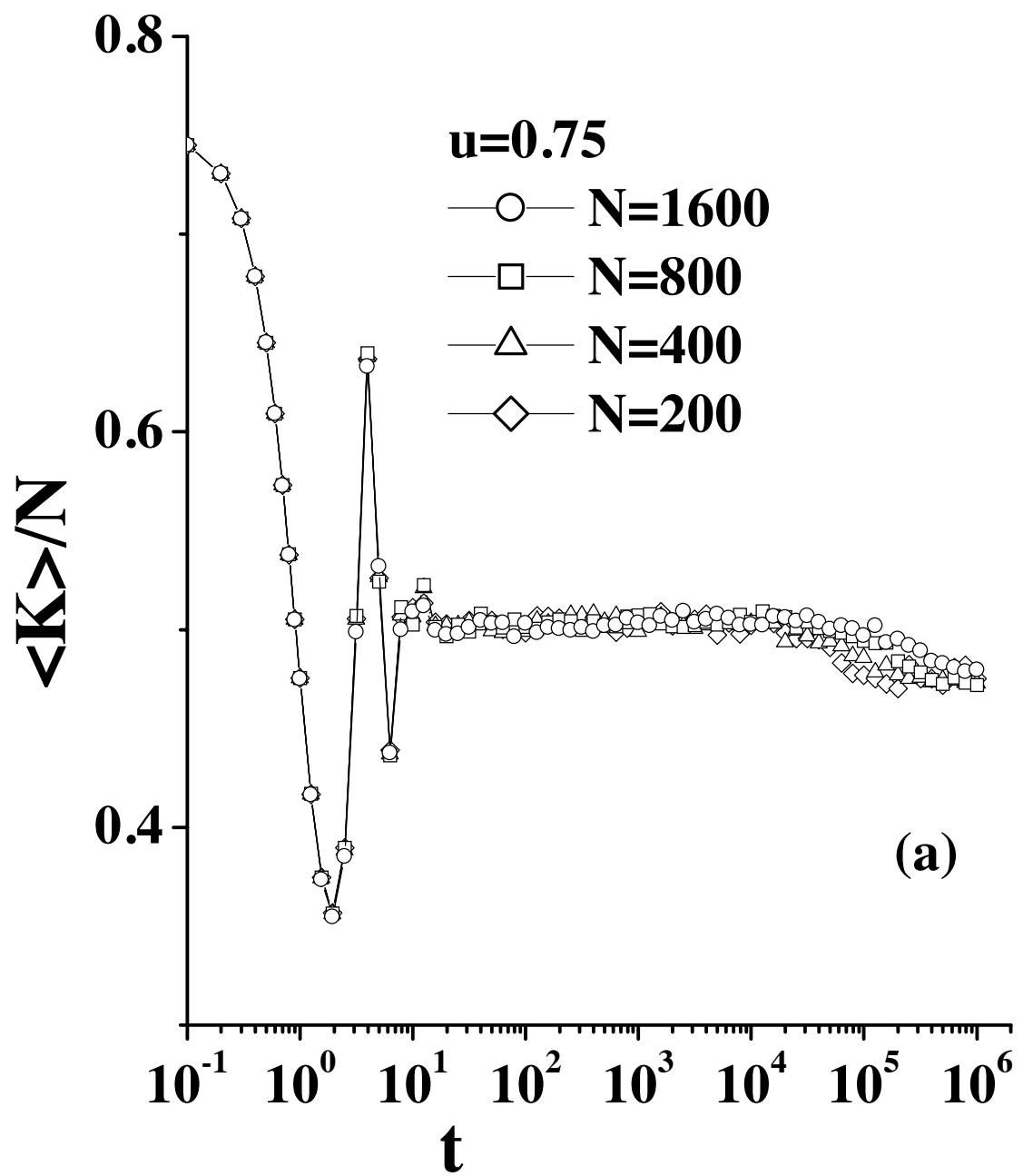
- Possible Alternative:
Nonextensive Statistical Mechanics

- Tsallis Entropy:

$$S_q = k \frac{1}{q-1} \left(1 - \sum_j p_j^q \right)$$

- Limit $q \rightarrow 1$:

$$S_1 = -k_B \sum_j p_j \ln p_j$$



ii) Different Initial “Water Bags”

$$\langle K_z \rangle = \eta \langle K \rangle \quad (0 \leq \eta \leq 1)$$

$$\langle K_x \rangle = \langle K_y \rangle = \frac{1}{2}(1 - \eta) \langle K \rangle$$

