

Universidade Estadual de Maringá



Ervin Kaminski Lenzi

Departamento de Física, Universidade Estadual de Maringá, Avenida Colombo, 5790, 87020-900, Maringá-PR, Brazil

Comb – Model and Extensions

Equation to be considered

$$\begin{split} \frac{\partial}{\partial t} \rho(x,y;t) &= \mathcal{D}_y \frac{\partial^2}{\partial y^2} \rho(x,y;t) + \mathcal{D}_x \delta(y) \bigg(\frac{\partial^2}{\partial x^2} - \overline{v}_x \frac{\partial}{\partial x} \bigg) \rho(x,y;t) \\ &- \nabla \cdot (\vec{v} \rho(x,y;t)) \, . \end{split}$$

It is subjected to the boundary and initial conditions

$$\rho(\pm\infty, y; t) = 0 \text{ and } \rho(x, \pm\infty; t) = 0$$

$$\rho(x, y; 0) = \widehat{\rho}(x, y)$$

First case

$$v_x \neq 0, v_y \neq 0, \text{ with } \overline{v}_x = 0$$

For this case, the solution is given by

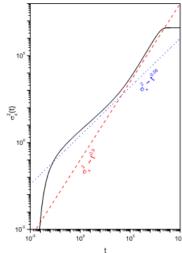
$$\begin{split} \rho(x,y;t) &= -\int_{-\infty}^{\infty} d\overline{y} \widehat{\rho}(x,y) \mathcal{G}(x,y,\overline{y};t) \\ \mathcal{G}(x,y,\overline{y};t) &= -e^{\frac{vy}{2\mathcal{D}y}(y-\overline{y})} e^{-\frac{v_y^2}{4\mathcal{D}y}t} \overline{\mathcal{G}}'(x,y,\overline{y};t) \\ \overline{\mathcal{G}}'(x,y,\overline{y};t) &= \frac{1}{\sqrt{4\pi\mathcal{D}_y t}} \delta(x-v_x t) \left(e^{-\frac{(y-\overline{y})^2}{4\mathcal{D}_y t}} - e^{-\frac{(|y|+|\overline{y}|)^2}{4\mathcal{D}_y t}} \right) + \frac{1}{\sqrt{8\mathcal{D}_x \sqrt{\mathcal{D}_y}}} \frac{|y|+|\overline{y}|}{\sqrt{4\pi\mathcal{D}_y}} \\ &\times \int_0^t d\overline{t} \frac{e^{-\frac{(|y|+|\overline{y}|)^2}{4\mathcal{D}_y (t-\overline{t})}}}{[(t-\overline{t})\overline{t}^{\frac{1}{2}}]^{\frac{3}{2}}} H_{1,1}^{1,0} \left[\sqrt{\frac{2}{\mathcal{D}_x} \sqrt{\frac{\mathcal{D}_y}{\overline{t}}}} |x-v_x \overline{t}| \left| \frac{\left(\frac{1}{4}, \frac{1}{4}\right)}{(0, 1)} \right| \right] \end{split}$$

Second Case

$$v_x \neq 0, v_y \neq 0, \text{ with } \overline{v}_x \neq 0$$

The Green function for this case is

$$\begin{split} \mathcal{G}(x,y,\overline{y};t) &= -e^{-\frac{v_y^{\overline{y}}}{4\mathcal{D}y}t}e^{\frac{vy}{2\mathcal{D}y}(y-\overline{y})}\widetilde{\mathcal{G}}(x,y,\overline{y};t) \\ \widetilde{\mathcal{G}}(x,y,\overline{y};t) &= \frac{1}{\sqrt{4\pi\mathcal{D}_yt}}\delta(x-v_xt)\left(e^{-\frac{(y-\overline{y})^2}{4\mathcal{D}yt}}-e^{-\frac{(|y|+|\overline{y}|)^2}{4\mathcal{D}yt}}\right) \\ &+ \frac{1}{t}\int_0^\infty du\,(|y|+|\overline{y}|+2\mathcal{D}_yu)\,\mathcal{G}_y\,(|y|,|\overline{y}|,2\mathcal{D}_yu;u)\,\mathcal{G}_x\,(x,-\overline{v}_xu,-v_xt;t) \\ \text{and } \mathcal{G}_\alpha(x,y,z;u) &= e^{-\frac{1}{4\mathcal{D}_\alpha u}(x+y+z)^2}/\sqrt{4\pi\mathcal{D}_\alpha u}\;. \end{split}$$



The time behavior of the mean square displacement for $v_y = 5 \ 10^{-4}, \ v_x = 0, \ \overline{v}_x = 1,$ $\mathcal{D}_y = 5, \ \mathcal{D}_x = 10 \ and \ \tilde{y} = 0.1$

Impedance, Fractional Diffusion Equation, Experimental Data

Equations

$$\begin{split} \mathcal{A} & \frac{\partial^{\gamma}}{\partial t^{\gamma}} n_{\alpha}(z,t) + \mathcal{B} \frac{\partial}{\partial t} n_{\alpha}(z,t) = -\frac{\partial}{\partial z} j_{\alpha}(z,t) - \int_{-\infty}^{t} \zeta(t-t') n_{\alpha}(z,t') dt' \;, \\ & j_{\alpha}(z,t) = -\mathcal{D}_{\alpha} \frac{\partial}{\partial z} n_{\alpha}(z,t) \mp \frac{q \mathcal{D}_{\alpha}}{k_{B} T} n_{\alpha}(z,t) \frac{\partial}{\partial z} V(z,t) \\ & \frac{\partial^{2}}{\partial z^{2}} V(z,t) = -\frac{q}{\varepsilon} \left[n_{+}(z,t) - n_{-}(z,t) \right] , \end{split}$$

These equations are subjected to the conditions

$$V\left(\pm \frac{d}{2}, t\right) = \pm \frac{V_0}{2} e^{i\omega t}$$

$$j_{\alpha}(z, t)|_{z=\pm \frac{d}{2}} = \pm k_{\alpha, e} E\left(z, t\right)|_{z=\pm \frac{d}{2}}$$

$$\pm \int_0^1 d\overline{\vartheta} \, \widetilde{\tau}(\overline{\vartheta}) \int_{-\infty}^t d\overline{t} \mathcal{K}_{\alpha}(t - \overline{t}, \overline{\vartheta}) \frac{\partial^{\overline{\vartheta}}}{\partial \overline{t}^{\overline{\vartheta}}} n_{\alpha}(z, \overline{t})\Big|_{z=\pm \frac{d}{2}}$$

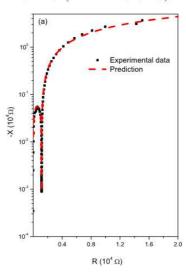
The impedance for this case is given by

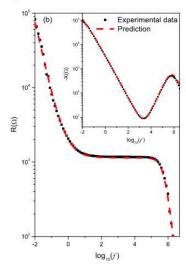
$$\mathcal{Z} = \frac{2}{\mathcal{S}\varepsilon\nu_{-}\Delta(i\omega)} \left(\frac{1}{\lambda^{2}\nu_{-}} \tanh\left(\nu_{-}d/2\right) + \frac{d}{2\mathcal{D}}\mathcal{E}(i\omega) \right)$$

$$\Delta(i\omega) = \left(1/\lambda^{2} + \chi\left(\Lambda(i\omega) + \omega_{e}\right)/\mathcal{D}\right) \left(i\omega + \omega_{e}\right)/\nu_{-}$$

$$+ \left(1/\lambda^{2} + \chi\left(i\omega + \omega_{e}\right)/\mathcal{D}\right) \Upsilon(i\omega) \tanh\left(\nu_{-}d/2\right)$$
with $\mathcal{E}(i\omega) = \chi\left(\Lambda(i\omega) + \omega_{e}\right)/\nu_{-}$ and $\mathcal{E}(i\omega) = \chi(\lambda(i\omega) + \omega_{e})/\nu_{-}$

with
$$\mathcal{E}(i\omega) = \chi \left(\Lambda(i\omega) + \omega_e\right) + \chi \nu_- \Upsilon(i\omega) \tanh \left(\nu_- d/2\right)$$
, $\chi = 1/\left(1 - \lambda^2 \overline{k}_e q/\left(\mathcal{D}\varepsilon\right)\right)$, $\omega_e = \overline{k}_e q/\varepsilon$ and \mathcal{S} is the electrode area.





These figure compare the experimental data with the prediction of the model presented here for the real, $R = Re \mathcal{Z}$, and imaginary, $X = Im \mathcal{Z}$, parts of the impedance. We have a good agreement between the experimental data and the predictions for the parameters values: $S = 3.14 \times 10^{-4} \ m^2$, $\varepsilon = 75\varepsilon_0$ ($\varepsilon_0 = 8.85 \times 10^{-12} \ C^2/(Nm^2)$), $\gamma = 0.98$, $\mathcal{D} = 2.9 \times 10^{-9} \ m^2/s$, $d = 10^{-3} m$, $\kappa_{a,1} = 8.6 \times 10^{-5} \ m/s$, $\kappa_{a,2} = 1.03 \times 10^{-7} \ m/s$, $\overline{k}_e = 0$, $\mathcal{A} = 0.99$, $\lambda = 2.734 \times 10^{-8} m$, $\tau = 1.5 \times 10^{-3} s$, $\eta_1 = 0.17$, and $\eta_2 = 0.825$.